

Università degli Studi di Napoli Federico II
 Corsi di
 Dinamica e Simulazione del Volo (prof. D. Coiro)
 Aeroelasticità (prof. F. Marulo)

DYNAMICS OF FLEXIBLE AIRCRAFT
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Outline

Part 1: Introduction to the dynamics of flexible aircraft

- Motivations for the study of flight dynamics of flexible aircraft
 - faster, lighter, more deformable
 - active control of deformations for improved riding qualities
- Qualitative analysis of flexibility static effects on A/C dynamics
 - fuselage deformation
 - aileron reversal

Part 2: Writing the equations of motion of flexible aircraft

- An historical perspective (with some technical considerations):
 - transport variables (ODE's) and deformation variables (PDE's)
 - the problem of the “body frame”
- A mixed Newtonian-Lagrangian approach
 - generalized Euler equation
 - assumed modes method

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Introduction to the dynamics of flexible aircraft

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Why a flexible aircraft model?

Trend in aircraft technologies:

- faster aircraft → higher dynamic pressure → higher loads;
- lighter aircraft → less structural weight;
- more efficient aircraft → slender structures;

} → more deformation.

Result:

- wider deformations significantly alter aircraft shape, depending on flight and manoeuvre condition;
- lower structural frequencies close to flight dynamic frequencies and/or control system bandwidth.

Consequence:

- coupling between structural dynamics and piloting tasks needs to be taken into account at all levels;
- rigid body dynamics no longer sufficient for a satisfactory description of aircraft behaviour within its operational flight envelope

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What happens?

Static effects

- variation of stability margins, stability derivatives and control power (always in a worsening direction!).

Dynamic effects:

- **aeroelastic response:** coupling between aerodynamics and structural dynamics (includes significant unsteady aerodynamic effects) at relatively high frequency → usually not relevant for flight dynamics and control within pilot tasks bandwidth;
- **flexible aircraft response: coupling between pilot/SCAS input to control surfaces, aircraft response and deformation**

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What happens?

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Need for active structural control

Performance degradation in terms of

- reduced stability margins and control power
- lighter damping

requires active control systems to compensate for the effects of structural deformations on aircraft response to controls.

First application of active structural control on flying vehicles:

- rocket launchers (tested; nowadays current technology);
- transatmospheric vehicles (only theory!).

First application to fixed wing aircraft: Rockwell B1-B (Structural Mode Control System).

Many modern transport jet aircraft now features...

- active flutter suppression system;
- active structure control for improved ride qualities in turbulent air;
- and more...

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Need for active structural control

B1-B with small fully-rotating canard for active deformation control

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Modelling issues

- Flexible aircraft dynamics described by means of a **hybrid system of differential equations** (mixed and coupled ODE's and PDE's);
- Some form of discretisation for elastic DoF's is needed for direct simulation;
- Need for a simple yet reliable aerodynamic model, that allows for direct numerical simulation within a reasonable CPU time;
- Elastic DoF's frequencies may induce (usually unwanted) couplings with
 - pilot commands and response to aircraft acceleration (degradation of handling qualities);
 - actuator dynamics (aero-servo-elastic problem);
 - external disturbances (e.g. turbulence).

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Engineering issues

- When an automatic control system is required, a state observer (e.g. a Kalman filter) is needed in order to provide reasonable estimate of "elastic states" from available measurements;
- The problem of estimation is made more demanding by the high level of "noise" (e.g. vibrations, turbulence, etc.);
- High performance actuators are necessary in order to control the system up to relatively high bandwidths;
- An engineering choice needs to be made between
 - **all purpose configurations** (a single set of aerodynamic surfaces used for full aircraft control, including flexible DoF's);
 - **dedicated surfaces** for deformation control (conventional aerodynamic surfaces used for aircraft control).

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Static vs dynamic flexible A/C models

- Method of **quasi-static deformations**:
 - deformations are sufficiently large to affect aircraft stability derivatives and/or control power;
 - structural mode frequencies are sufficiently high, such that only steady-state deformations affect aircraft dynamics.
- Lagrangian or Hamiltonian approaches to derive **fully dynamic models** of flexible aircraft that couples structural and transport dynamic variables:
 - based on FEM (high order, highly accurate, complex, may require unsteady aerodynamics);
 - based on low-order approximations of deformation (e.g. method of assumed modes);

Both techniques require choice of a suitable set of “body” axes and the definition of transport variables.
- **Note:** Dynamic models discussed in Part 2.

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Effects of fuselage deformation (static case)

Primary effect of fuselage deformation in the longitudinal plane is a reduction of tail incidence

Tail angle of attack:

rigid aircraft $\alpha_t^* = \alpha_{wb} + I_t - \epsilon$

with

α_{wb}	aircraft angle of attack
I_t	tail incidence setting
ϵ	downwash angle

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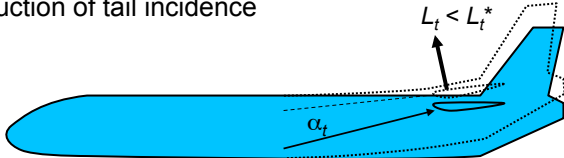
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Effects of fuselage deformation

Primary effect of fuselage deformation in the longitudinal plane is a reduction of tail incidence



Tail angle of attack:
 rigid aircraft $\alpha_t^* = \alpha_{wb} + i_t - \varepsilon$
 elastic aircraft $\alpha_t = \alpha_{wb} + i_t - \varepsilon - k L_t$

Tail incidence is reduced by fuselage flexibility!

Tail lift coefficient: $C_{L_t} = C_{L_{\alpha t}} [\alpha_{wb} + i_t - \varepsilon - k (\frac{1}{2} \rho V_t^2 S_t C_{L_t})]$
 Solving for C_{L_t} : $C_{L_t} = \frac{C_{L_{\alpha t}} (\alpha_{wb} + i_t - \varepsilon)}{1 + k \frac{1}{2} \rho V_t^2 S_t C_{L_{\alpha t}}}$

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Effects of fuselage deformation

It is possible to define an effective tail lift gradient:

$$C_{L_{\alpha t}}^{eff} = \frac{C_{L_{\alpha t}}}{1 + k \frac{1}{2} \rho V_t^2 S_t C_{L_{\alpha t}}} < C_{L_{\alpha t}}$$

Effects of fuselage deformation on longitudinal static stability

Position of neutral point: $\frac{x_N}{\bar{c}} = \frac{x_{AC_{wb}}}{\bar{c}} + \frac{C_{L_{\alpha t}}}{C_{L_{\alpha_{wb}}}} \bar{V}_H \left(1 - \frac{\partial \varepsilon}{\partial \alpha} \right)$

Static stability derivative: $C_{m_{\alpha}} = C_{L_{\alpha}} \left(\frac{x_{CG}}{\bar{c}} - \frac{x_N}{\bar{c}} \right) < 0$ if $x_{CG} < x_N$

Variation of x_N : $\frac{\Delta x_N}{\bar{c}} = \frac{\Delta C_{L_{\alpha t}}}{C_{L_{\alpha_{wb}}}} \bar{V}_H \left(1 - \frac{\partial \varepsilon}{\partial \alpha} \right)$

where $\Delta C_{L_{\alpha t}} = C_{L_{\alpha t}}^{eff} - C_{L_{\alpha t}} = - \frac{\frac{1}{2} k \rho V_t^2 S_t C_{L_{\alpha t}}^2}{1 + \frac{1}{2} k \rho V_t^2 S_t C_{L_{\alpha t}}} < 0$

that is, **static margin is reduced!**

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Effects of fuselage deformation

Effects of fuselage deformation on longitudinal control power

Tail lift coefficient: $C_{L_t} = C_{L_{\alpha_t}} [\alpha_{wb} + i_t - \varepsilon - k L_t] + C_{L_t \delta_E} \delta_E$

Solving for C_{L_t} :
$$C_{L_t} = \frac{C_{L_{\alpha_t}} (\alpha_{wb} + i_t - \varepsilon)}{1 + \frac{1}{2} k \rho V_t^2 S_t C_{L_{\alpha_t}}} + \left[\frac{C_{L_t \delta_E}}{1 + \frac{1}{2} k \rho V_t^2 S_t C_{L_{\alpha_t}}} \right] \delta_E$$

Effective control power:
$$C_{L_t \delta_E}^{eff} = \frac{C_{L_t \delta_E}}{1 + \frac{1}{2} k \rho V_t^2 S_t C_{L_{\alpha_t}}} < C_{L_t \delta_E}$$

Again, **control power is reduced!!**

Effects of fuselage deformation on pitch damping

Tail incidence induced by pitch motion: $\Delta \alpha_t = q l_t / V_t$

Tail contribution to pitch damping coefficient:
$$C_{m_q} = -2 C_{L_{\alpha_t}} \bar{V}_H \frac{l_t}{\bar{c}}$$

The reduction of $C_{L_{\alpha_t}}$ **decreases pitch damping coefficient too!!!**

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Effects of fuselage deformation

Effects of fuselage deformation on longitudinal modes

- Static margin reduction lowers short-period frequency
- Reduced horizontal tail lift gradient lowers damping

Consequence: worse HQ expected

- A contribution to pitch moment vs speed derivative shows up

In general $|M_u| \ll 1$ is neglected. It becomes non-negligible at high speed because of

- aerodynamic effects due to transonic aerodynamics;
- fuselage flexibility

Remembering that
$$C_{m_i} = -\bar{V}_H C_{L_i} = -\bar{V}_H \frac{C_{L_{\alpha_t}} (\alpha_{wb} + i_t - \varepsilon)}{1 + \frac{1}{2} k \rho V_t^2 S_t C_{L_{\alpha_t}}}$$

one gets
$$C_{m_u} = \frac{\partial C_{m_i}}{\partial (\Delta u / V)} = -\bar{V}_H V \frac{\partial C_{L_i}}{\partial u}$$

$$\approx \bar{V}_H \frac{C_{L_{\alpha_t}} (\alpha_{wb} + i_t - \varepsilon)}{\left(1 + \frac{1}{2} k \rho V_t^2 S_t C_{L_{\alpha_t}}\right)^2} k \rho V_t^2 S_t C_{L_{\alpha_t}} = -\frac{2k Q_t S_t C_{L_{\alpha_t}}}{1 + k Q_t S_t C_{L_{\alpha_t}}} C_{m_i}$$

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Wing torsion induced by aileron deflection

- Aileron deflection generates a rolling moment by means of an anti-symmetric deflection of aerodynamic surfaces on the wing (assuming $C_{l\delta_A} < 0$)

- A significant torsion moment is also developed, that reduces wing section incidence and, as a consequence, the lift increment and the resulting control moment

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Wing torsion induced by aileron deflection

- Torsion moment is roughly proportional to dynamic pressure and aileron deflection: $M_T \propto \frac{1}{2} \rho V^2 \delta_A$
- Torsional deformation ϑ is proportional to torsion moment M_T , and it corresponds to a variation of airfoil incidence in the opposite direction, that is: $\vartheta(y) = -\Delta\alpha(y) \propto M_T$
- As a consequence, the variation of roll control moment coefficient ΔC_l due to torsion effects can be written in the form

$$\Delta C_l = k_T \frac{1}{2} \rho V^2 \delta_A$$
- The resulting total roll control moment is thus given by

$$\Delta C_l^{tot} = \Delta C_l^{rig} + \Delta C_l^{flex} = C_{l\delta_A} \delta_A + \frac{1}{2} k_T \rho V^2 \delta_A$$
 and the effective control moment gradient becomes

$$C_{l\delta_A}^{eff} = C_{l\delta_A} + \frac{1}{2} k_T \rho V^2$$
- Aileron reversal speed V_R is defined as the velocity such that the control gradient vanishes $C_{l\delta_A} + \frac{1}{2} k_T \rho V_R^2 = 0$
- The effective roll control gradient becomes $C_{l\delta_A}^{eff} = C_{l\delta_A} \left[1 - \frac{V^2}{V_R^2} \right]$

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Equations of motion of flexible aircraft

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Why is it so difficult?

Transport vs deformation variables

- Transport variables (position vector, speed, Euler angles and angular velocity)
 - depend on time t only;
 - describe “global” properties of aircraft motion;
 - their evolution can be described in terms of a set of ODE’s.
- Deformation variables
 - depend on time t and position in an aircraft-“fixed” frame;
 - describe “local” properties of structure motion;
 - their evolution is described by means of a set of PDE’s;
- A hybrid set of highly coupled highly nonlinear ordinary and partial differential equations is obtained.
- Some form of discretization of the latter is required.
- Coupling extends to inertial terms (highest differential order), so that the system of O+PDE’s is in non-normal form.

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Why is it so difficult?

**Is there something like an aircraft-fixed frame?
Unfortunately the answer is NO!!!**

- Body-fixed frame centred in aircraft CoG defined for the rigid aircraft case only.
- Position of CoG and other inertial properties (e.g. moments of inertia) depend on deformation state.
- At the same time a frame representative of aircraft position and attitude is necessary!
- Two possible choices:
 1. Mean axes (Milne, 1964)
 2. Pseudo-body axes (Tuzcu and Meirovitch, 2003)

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Linear and angular momentum balance

Momentum balance $\frac{d}{dt}(mv_C) = \Sigma F$

Absolute angular momentum balance $\frac{d}{dt}(h_O) = \Sigma M_O$
(pole O fixed in the inertial frame)

Relative angular momentum balance $\frac{d}{dt}(h_C) = \Sigma M_C$
(only when pole C is the centre of mass)

For a deformable body it is

$$h_C = \int_{\mathcal{E}} (r \times v) \delta m = \int_{\mathcal{E}} [r \times (v_C + \omega \times r + \dot{r})] \delta m =$$

$$= \underbrace{\int_{\mathcal{E}} \delta m}_{0} \times v_C + \underbrace{\int_{\mathcal{E}} [r \times (\omega \times r)] \delta m}_{I \omega} + \underbrace{\int_{\mathcal{E}} (r \times \dot{r}) \delta m}_{\text{distortional component of angular momentum}}$$

by definition of CoM "rigid-body" term

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Definition of the “mean axes” frame

Statement: whichever the deformation state, there always exists a frame centred in the centre of mass such that

$$\int_{\mathcal{E}} (r \times \dot{r}) \delta m = 0$$

This frame is known as the **mean axes frame**.

Proof: assume two different body-frames, F_{B1} and F_{B2} . It is

$$h_{C1} = I_{B1} \omega_{B1} + \int_{\mathcal{E}_1} (r_{B1} \times \dot{r}_{B1}) \delta m$$

$$h_{C2} = I_{B2} (\omega + \Delta\omega)_{B2} + \int_{\mathcal{E}_2} (r_{B2} \times \dot{r}_{B2}) \delta m$$

$$\int_{\mathcal{E}_2} (r_{B2} \times \dot{r}_{B2}) \delta m = \int_{\mathcal{E}_2} [T_{B2B1} \tilde{R}_{B1} T_{B1B2}] [T_{B2B1} (\dot{r}_{B1} + \Delta\omega_{B1} \times r_{B1})] \delta m$$

$$= T_{B2B1} \int_{\mathcal{E}_2} (\tilde{R}_{B1} \dot{r}_{B1}) \delta m + T_{B2B1} \int_{\mathcal{E}_2} [\tilde{R}_{B1} (\Delta\omega_{B1} \times r_{B1})] \delta m$$

$$= T_{B2B1} \left[\int_{\mathcal{E}_2} (r_{B1} \times \dot{r}_{B1}) \delta m + I_{B1} \Delta\omega_{B1} \right]$$

The distortional term thus disappears if $\Delta\omega_{B1} = -I_{B1}^{-1} \int_{\mathcal{E}_2} (r_{B1} \times \dot{r}_{B1}) \delta m$ that is, the coordinate transformation matrix evolves according to the equation $T_{B1B2} = \Delta\tilde{\Omega}_{B1} T_{B1B2}$.

In such a case the angular momentum has the form $h_C = I\omega$

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Pros & cons of the “mean axes” frame

Statement: whichever the deformation state, there always exists a frame centred in the centre of mass such that

$$\int_{\mathcal{E}} (r \times \dot{r}) \delta m = 0$$

This frame is known as the **mean axes frame**.

Pros:

- frame centred in the vehicle centre of mass → relative angular momentum balance equation holds;
- major inertial coupling term between transport and deformation DoF’s is removed from E.o.M. → if variations of inertia tensor are negligible, no coupling on the LHS of E.o.M.

Cons:

- significant coupling between transport and deformation DoF’s is always present in the aerodynamic terms;
- centre of mass depends on deformation state, so that its velocity depends on deformation rate too;
- most important: mean axes always exist, but the identification of their actual position is far from trivial (i.e. truly difficult!).

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Pros & cons of the “mean axes” frame

Definition: A set of **pseudo-body axes** is given by a reference frame attached to the centre of mass of the undeformed aircraft structure and fixed with respect to it.

Pros:

- simple definition;
- deformation state naturally described with respect to the undeformed condition.

Cons:

- the actual centre of mass moves with respect to the origin of the frame → usual angular momentum balance equations (w.r.t. the CoM) can no longer be used;
- fully coupled equations are obtained (local inertial acceleration depends on transport acceleration, which in turn depends on time derivative of deformation rates).

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Why is it so difficult? Last issue

Derivation of a finite order model (ODE system) may follow different paths:

- Direct derivation from a global Hamiltonian discretized by means of FEM in a set of mean axes (Cavin III & Dusto, 1977);
- Inertially decoupled models with aerodynamic coupling only in a set of approximate mean axes (Waszak et al., 1987);
- Direct derivation from a global Lagrangian function in terms of quasi-velocity variables expressed in a set of pseudo-body axes; discretization performed on the full set of hybrid O+PDE's (Tuzcu and Meirovitch, 2003);
- Mixed Newtonian-Lagrangian approach in a set of pseudo-body axes (Avanzini, Capello, Piacenza, 2014; from approach by Junkins et al. for space structures); discretization of deformation variables performed on the Lagrangian → **amplitudes of assumed modes** used as **generalized variables**; transport dynamics derived by **generalized Euler equation**.

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Generalized Euler equation

Momentum balance (for constant mass): $\frac{d}{dt}(mv_C) = F$

Relative angular momentum balance
(referred to body centre of mass) $\frac{d}{dt}(h_C) = M_C$

For a different pole $A \neq C$ it is $M_A = M_C + r_{AC} \times F$ and $a_C = a_A + \frac{d^2 r_{AC}}{dt^2}$

$$h_A = \int_{\mathcal{B}} \left(r_{AP} \times \frac{dr_{AP}}{dt} \right) \delta m = \int_{\mathcal{B}} \left[(r_{AC} + r_{CP}) \times \frac{d}{dt} (r_{AC} + r_{CP}) \right] \delta m$$

$$= \int_{\mathcal{B}} \left(r_{AC} \times \frac{dr_{AC}}{dt} \right) \delta m + \int_{\mathcal{B}} \left(r_{AC} \times \frac{dr_{CP}}{dt} \right) \delta m + \int_{\mathcal{B}} \left(r_{CP} \times \frac{dr_{AC}}{dt} \right) \delta m + \int_{\mathcal{B}} \left(r_{CP} \times \frac{dr_{CP}}{dt} \right) \delta m$$

0
from the definition of CoM

$$= m \left(r_{AC} \times \frac{dr_{AC}}{dt} \right) + h_C$$

$$\frac{d}{dt} \left(h_A - m r_{AC} \times \frac{dr_{AC}}{dt} \right) = M_A - m r_{AC} \times \left(a_A + \frac{d^2 r_{AC}}{dt^2} \right)$$

$$\frac{dh_A}{dt} - m r_{AC} \times \frac{d^2 r_{AC}}{dt^2} = M_A - m r_{AC} \times a_A - m r_{AC} \times \frac{d^2 r_{AC}}{dt^2}$$

Generalized Euler equation: $\frac{dh_A}{dt} + s_A \times a_A = M_A$
where $s_A = m r_{AC}$ is the static moment

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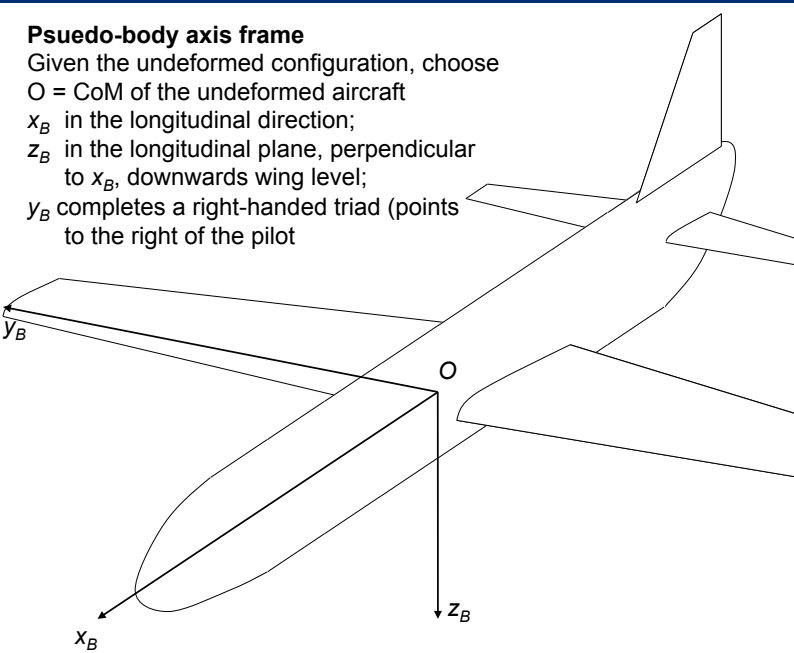
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Description of flexible aircraft

Pseudo-body axis frame
Given the undeformed configuration, choose
 $O = \text{CoM}$ of the undeformed aircraft
 x_B in the longitudinal direction;
 z_B in the longitudinal plane, perpendicular to x_B , downwards wing level;
 y_B completes a right-handed triad (points to the right of the pilot)



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Description of flexible aircraft

Transport variables

$\mathbf{v}_B = (u, v, w)^T$ velocity components

$\boldsymbol{\omega}_B = (p, q, r)^T$ angular velocity components

ϕ, θ, ψ roll, pitch and yaw angles

$\mathbf{r}_1 = (x, y, -h)^T$ position vector

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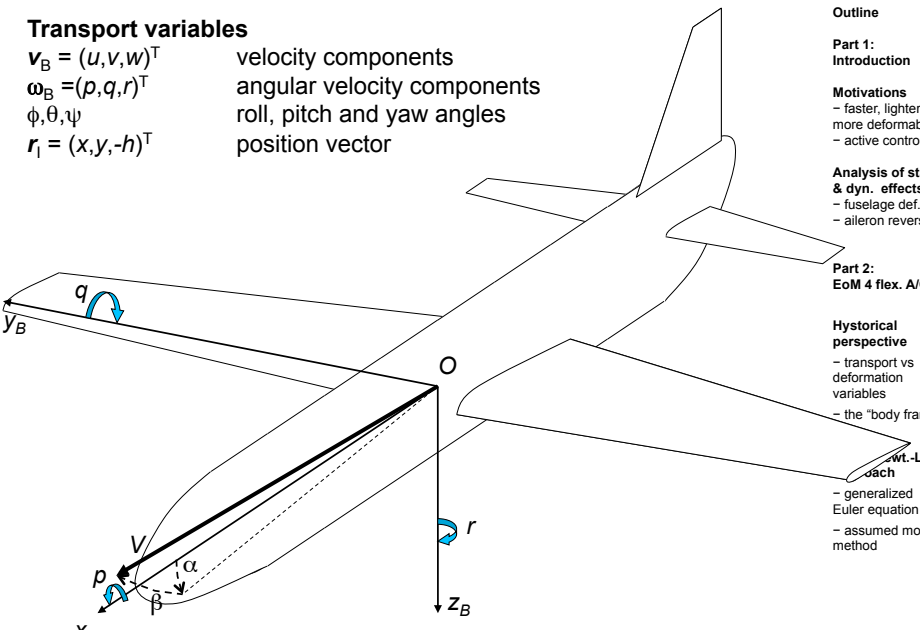
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Description of flexible aircraft

Deformable elements

aft. portion of the fuselage (with tip mass)

wing

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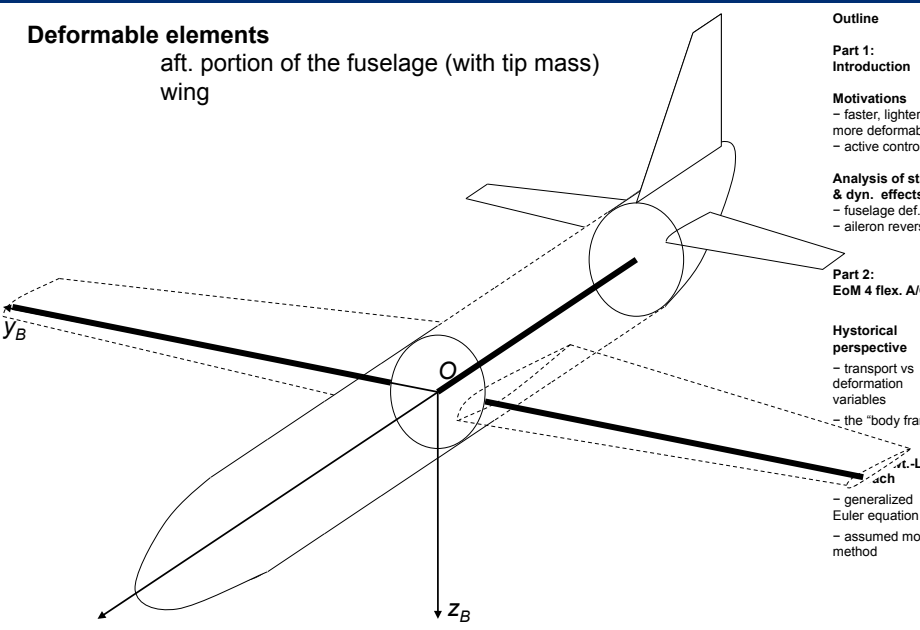
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Description of flexible aircraft

Deformable elements
aft. portion of the fuselage
wings

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State variables and discretization

State vector $\mathbf{x} = (\mathbf{x}_T^T, \mathbf{x}_D^T, \dot{\mathbf{x}}_D^T)^T$

Transport variables: $\mathbf{x}_T = (u, v, w, p, q, r, \phi, \theta, \psi; \Delta x, \Delta y, -h)^T$

Deformation variables:

1. Discretization

Fuselage ($i=z, y$) flexural deformation
(z =longitudinal; y =lateral)

Fuselage torsional deformation

Wing ($i=r, l$) flexural deformations
(r =right; l =left)

Wing ($i=r, l$) torsional deformation
(r =right; l =left)

2. Resulting state vector for deformation variables:

$\mathbf{x}_D = (\eta_1^{(f,y)}, \dots, \eta_N^{(f,y)}; \eta_1^{(f,z)}, \dots, \eta_N^{(f,z)}; \sigma_1^{(f)}, \dots, \sigma_N^{(f)};$
 $\eta_1^{(w,l)}, \dots, \eta_N^{(w,l)}; \eta_1^{(w,r)}, \dots, \eta_N^{(w,r)}; \sigma_1^{(w,l)}, \dots, \sigma_N^{(w,l)}; \sigma_1^{(w,r)}, \dots,$
 $\sigma_N^{(w,r)})^T$

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Lagrangian derivation of flexible dynamics

Kinetic energy:
 Definition $T = \frac{1}{2} \int_{\mathcal{E}} (\mathbf{v}_P \cdot \mathbf{v}_P) \delta m$
 $= \int_{\mathcal{E}} [\mathbf{v}_O + \omega \times (\mathbf{r}_P + \rho_P) + \rho_P] \cdot [\mathbf{v}_O + \omega \times (\mathbf{r}_P + \rho_P) + \rho_P] \delta m$

Terms $\mathbf{T} = \mathbf{T}(\mathbf{v}_B, \omega_B, \dot{\mathbf{x}}_D)$
 $= \mathbf{T}_{RIG}(\mathbf{v}_B, \omega_B) + \mathbf{T}_{FLEX}(\mathbf{x}_D) + \mathbf{T}_{COUP}(\mathbf{v}_B, \omega_B, \dot{\mathbf{x}}_D)$

Potential energy:

$$V = \frac{1}{2} \int_0^l \left[EI_{f,y}(x_f) \left(\frac{\partial^2 \xi_{f,z}}{\partial x^2} \right)^2 + EI_{f,z}(x_f) \left(\frac{\partial^2 \xi_{f,y}}{\partial x^2} \right)^2 + GJ_{f,x}(x_f) \left(\frac{\partial \theta_f}{\partial x} \right)^2 \right] dx_f$$

$$+ \frac{1}{2} \int_0^w \left[EI_{w,l}(x_w) \left(\frac{\partial^2 \xi_{w,l}}{\partial x^2} \right)^2 + GJ_{w,l}(x_w) \left(\frac{\partial \theta_{w,l}}{\partial x} \right)^2 \right] dx_w$$

$$+ \frac{1}{2} \int_0^w \left[EI_{w,r}(x_w) \left(\frac{\partial^2 \xi_{w,r}}{\partial x^2} \right)^2 + GJ_{w,r}(x_w) \left(\frac{\partial \theta_{w,r}}{\partial x} \right)^2 \right] dx_w$$

Terms $\mathbf{V} = \mathbf{V}(\mathbf{x}_D) = \frac{1}{2} \mathbf{x}_D^T \mathbf{K} \mathbf{x}_D$ with \mathbf{K} block diagonal

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Lagrangian derivation of flexible dynamics

Lagrange equations

$$\frac{d}{dt} \left(\frac{dT}{dx_D} \right) - \left(\frac{dT}{dx_D} \right) + \left(\frac{dV}{dx_D} \right) = Q$$

with Q a vector of generalized non-conservative forces (including aerodynamic ones)

$$\delta W = \sum_{k=1}^{N_{tot}} Q_k \delta x_{D_k} = \int_B \mathbf{f}_P^{(NC)} \cdot \delta \xi$$

Comments:

- Aerodynamics forces can be expressed in the form of a Raileigh function

$$\mathbf{F} = \mathbf{F}(\dot{\mathbf{x}}_D) = \frac{1}{2} \dot{\mathbf{x}}_D^T \mathbf{F} \dot{\mathbf{x}}_D$$
 such that $Q_A = \mathbf{F} \dot{\mathbf{x}}_D = \partial \mathbf{F} / \partial \dot{\mathbf{x}}_D$ (omitted for sake of brevity).
- The coupling term $\mathbf{T}_{COUP}(\mathbf{v}_B, \omega_B, \dot{\mathbf{x}}_D)$ in the kinetic energy causes inertial coupling between flexible variable dynamics and transport acceleration.
- The additional term in the generalized Euler equation ($m r_{AC} x_{A_A}$) and the variation of mass properties with deformation couple transport variable dynamics with deformation rate derivatives.

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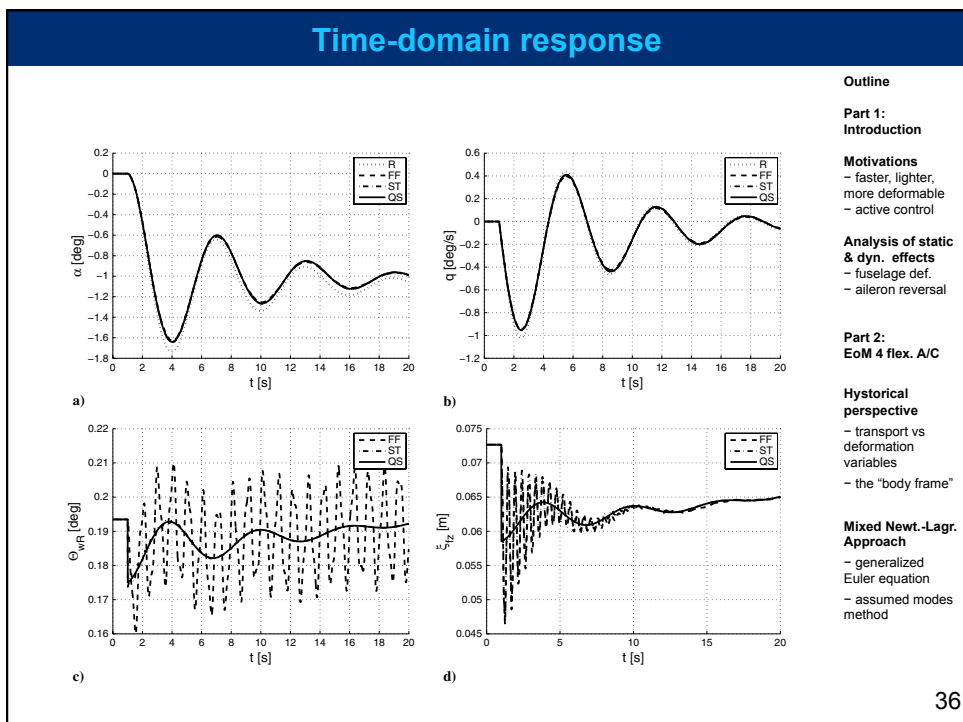
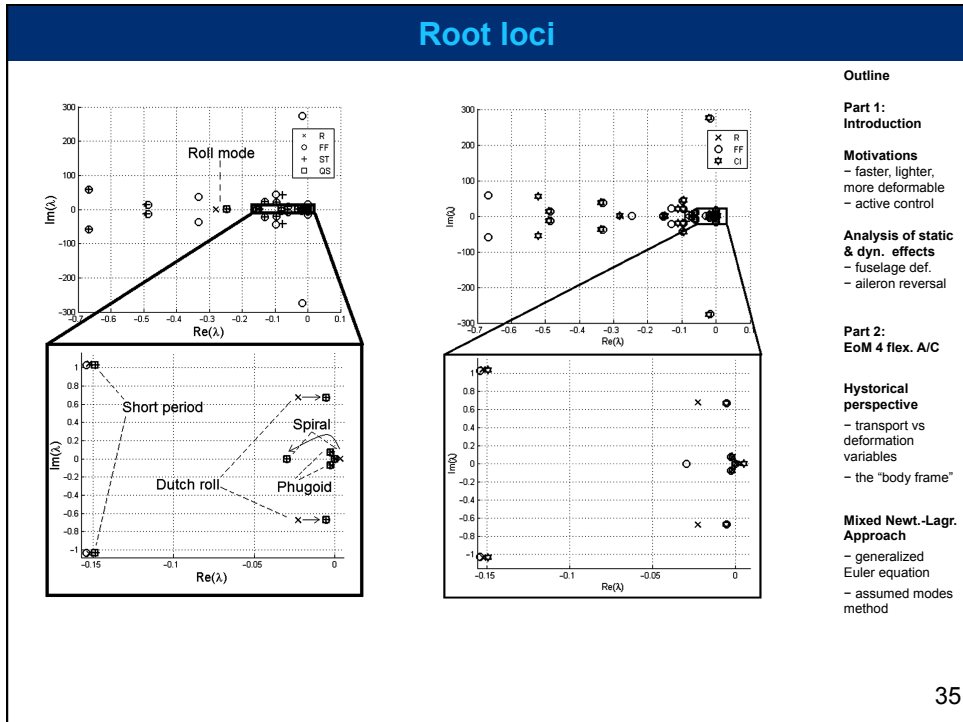
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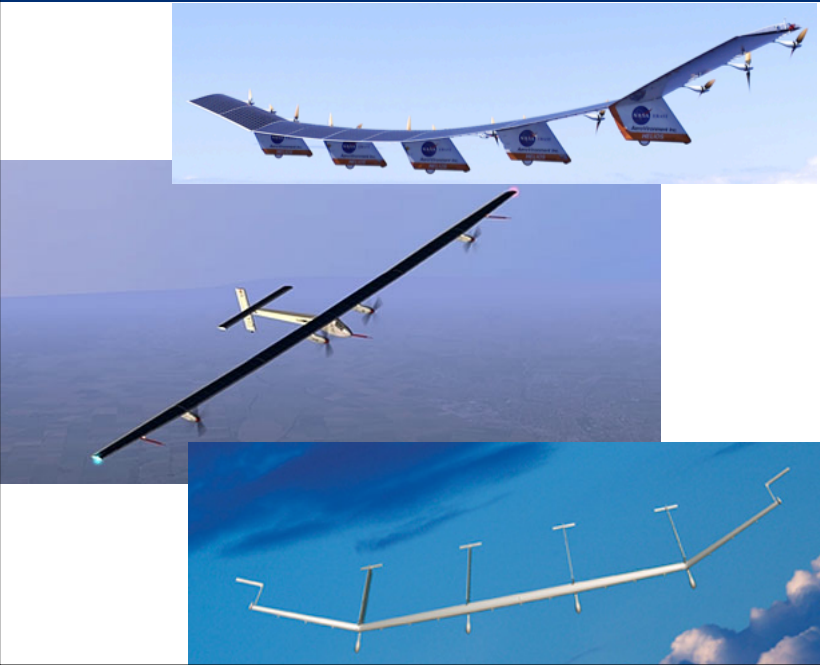
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Open issues...



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... that someone should solve!



- the "body frame"

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Questions?



If not too bored...

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