## XX Congresso AIMETA

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## Giovanni Romano

DIST - Dipartimento di Ingegneria STrutturale Università di Napoli Federico II, Italia

Conferenza Generale
del 13 Settembre 2011

## On the Geometric Approach to Non-Linear Continuum Mechanics

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Linearized Continuum Mechanics (LCM) can be modeled by Linear Algebra (LA) and Calculus on Linear Spaces (CoLS).

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Linearized Continuum Mechanics (LCM) can be modeled by Linear Algebra (LA) and Calculus on Linear Spaces (CoLS).

Non-Linear Continuum Mechanics (NLCM) calls instead for Differential Geometry (DG) and Calculus on Manifolds (CoM) as natural tools to develop theoretical and computational models.


## Prolegomena

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## Prolegomena



## Prolegomena

In these days the angel of topology and the devil of abstract algebra fight for the soul of each individual mathematical domain.
H. Weyl, "Invariants", Duke Mathematical Journal 5 (3): (1939) 489-502

## A basic question

## Prolegomena



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## Adapted to NLCM

In these days the angel of differential geometry and the devil of algebra and calculus on linear spaces fight for the soul of each individual continuum mechanics domain.

## Prolegomena



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This lecture is in support of the angel.

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## Adapted to NLCM

In these days the angel of differential geometry and the devil of algebra and calculus on linear spaces fight for the soul of each individual continuum mechanics domain.
This lecture is in support of the angel.
Differential Geometry provides the tools to fly higher and see what before was shadowed or completely hidden.

## A basic question in NLCM

- How to compare material tensors at corresponding points in displaced configurations of a body?

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## A basic question in NLCM

- How to compare material tensors at corresponding points in displaced configurations of a body?
- Devil's temptation:

In 3D bodies it might seem as natural to compare by translation the involved material vectors.
This is tacitly done in literature, when evaluating the material time-derivative of the stress tensor $\mathbf{T}$ :

$$
\dot{\mathbf{T}}(\mathrm{p}, t):=\partial_{\tau=t} \mathbf{T}(\mathrm{p}, \tau)
$$

or the material time-derivative of the director $\mathbf{n}$ of a nematic liquid crystal:

$$
\dot{\mathbf{n}}(\mathrm{p}, t):=\partial_{\tau=t} \mathbf{n}(\mathrm{p}, \tau)
$$

These definitions are connection dependent and geometrically untenable when considering 1D and 2D models (wires and membranes).

## A basic question in NLCM

- How to compare material tensors at corresponding points in displaced configurations of a body?
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These definitions are connection dependent and geometrically untenable when considering 1D and 2D models (wires and membranes).

- Hint: Tangent vectors to a body placement are transformed into tangent vectors to another body placement by the tangent displacement map. This is the essence of the COVARIANCE PARADIGM.


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## Basic requirements

## DIMENSIONALITY INDEPENDENCE:

A geometrically consistent theoretical framework should be equally applicable to body models of any dimension.

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A geometrically consistent theoretical framework should be equally applicable to body models of any dimension.

COVARIANCE PARADIGM motivation ${ }^{1}$ :
${ }^{1}$ G. Romano, R. Barretta, Covariant hypo-elasticity.

## Basic requirements

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A geometrically consistent theoretical framework should be equally applicable to body models of any dimension.

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## Math1

Tangent vector to a manifold:

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## Tangent functor

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## Math1

Tangent vector to a manifold: velocity of a curve $\mathbf{c} \in \mathrm{C}^{1}([a, b] ; \mathbb{M}), \quad \lambda \in[a, b], \quad \mathbf{x}=\mathbf{c}(\lambda)$ base point

$$
\mathbf{v}:=\partial_{\mu=\lambda} \mathbf{c}(\mu) \in \mathbb{T}_{\mathbf{x}} \mathbb{M}
$$

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## Cotangent vector:

$$
\mathbf{v}^{*} \in L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M} ; \mathcal{R}\right) \in \mathbb{T}_{\mathbf{x}}^{*} \mathbb{M}
$$

## Math1

Tangent vector to a manifold: velocity of a curve $\mathbf{c} \in \mathrm{C}^{1}([a, b] ; \mathbb{M}), \quad \lambda \in[a, b], \quad \mathbf{x}=\mathbf{c}(\lambda)$ base point

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Tangent map:

## NLCM

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## Math1

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Tangent map:

- A map $\zeta \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{N})$ sends a curve $\mathbf{c} \in \mathrm{C}^{1}([a, b] ; \mathbb{M})$ into a curve $\boldsymbol{\zeta} \circ \mathbf{c} \in \mathrm{C}^{1}([a, b] ; \mathbb{N})$.


## Math1

Tangent vector to a manifold:
velocity of a curve $\mathbf{c} \in \mathbb{C}^{1}([a, b] ; \mathbb{M}), \quad \lambda \in[a, b], \quad \mathbf{x}=\mathbf{c}(\lambda)$ base point

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\mathbf{v}:=\partial_{\mu=\lambda} \mathbf{c}(\mu) \in \mathbb{T}_{\mathbf{x}} \mathbb{M}
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Cotangent vector:

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\mathbf{v}^{*} \in L\left(\mathbb{T}_{x} \mathbb{M} ; \mathcal{R}\right) \in \mathbb{T}_{x}^{*} \mathbb{M}
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Tangent map:

- A map $\boldsymbol{\zeta} \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{N})$ sends a curve $\mathbf{c} \in \mathrm{C}^{1}([a, b] ; \mathbb{M})$ into a curve $\boldsymbol{\zeta} \circ \mathbf{c} \in \mathrm{C}^{1}([a, b] ; \mathbb{N})$.
- The tangent map $T_{x} \zeta \in \mathrm{C}^{0}\left(\mathbb{T}_{x} \mathbb{M} ; \mathbb{T}_{\zeta(\mathrm{x})} \mathbb{N}\right)$ sends a tangent vector at $x \in \mathbb{M}$
$\mathbf{v} \in \mathbb{T}_{\mathbf{x}}(\mathbb{M}):=\partial_{\mu=\lambda} \mathbf{c}(\mu)$
into a tangent vector at $\zeta(x) \in \mathbb{N}$ $T_{\mathrm{x}} \boldsymbol{\zeta} \cdot \mathbf{v} \in \mathbb{T}_{\boldsymbol{\zeta}(\mathrm{x})}(\mathbb{N}):=\partial_{\mu=\lambda}(\boldsymbol{\zeta} \circ \mathbf{c})(\mu)$


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Tangent bundle

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## Math2

Tangent bundle

- disjoint union of tangent spaces:

$$
\mathbb{T M}:=\cup_{\mathbf{x} \in \mathbb{M}} \mathbb{T}_{\mathbf{x}} \mathbb{M}
$$



## Math2

Tangent bundle

- disjoint union of tangent spaces:

$$
\mathbb{T} \mathbb{M}:=\cup_{\mathbf{x} \in \mathbb{M}} \mathbb{T}_{\mathbf{x}} \mathbb{M}
$$



- Projection: $\boldsymbol{\tau}_{\mathbb{M}} \in \mathrm{C}^{1}(\mathbb{T M} ; \mathbb{M})$

$$
\mathbf{v} \in \mathbb{T}_{\mathbf{x}} \mathbb{M}, \quad \boldsymbol{\tau}_{\mathbb{M}}(\mathbf{v}):=\mathbf{x} \quad \text { base point }
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## Math2

Tangent bundle

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$$

- Surjective submersion:

$$
T_{\mathbf{v}} \tau_{\mathbb{M}} \in \mathrm{C}^{1}\left(\mathbb{T}_{\mathbf{v}} \mathbb{T M} ; \mathbb{T}_{\mathrm{x}} \mathbb{M}\right) \text { is surjective }
$$



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Tangent bundle

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- Tangent functor

$$
\zeta \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{N}) \quad \mapsto \quad T \zeta \in \mathrm{C}^{0}(\mathbb{T M} ; \mathbb{T})
$$

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- E, M manifolds


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Fiber bundles

- E,M manifolds
- Fiber bundle projection:

$\pi_{\mathbb{M}, \mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{M})$ surjective submersion


## Math3



- Fiber bundle projection: $\pi_{\mathrm{M}, \mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{M})$ surjective submersion
- Total space: E
- Base space: $\mathbb{M}$
- Fiber manifold: $\left(\boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}}(\mathbf{x})\right)^{-1}$ based at $\mathbf{x} \in \mathbb{M}$

Fiber bundles

- E, M manifolds


## Math3



- Fiber bundle projection: $\pi_{\mathbb{M}, \mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{M})$ surjective submersion
- Total space: E
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- Fiber manifold: $\left(\boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}}(\mathbf{x})\right)^{-1}$ based at $\mathbf{x} \in \mathbb{M}$
- Tangent bundle $T \pi_{\mathbb{M}, \mathrm{E}} \in \mathrm{C}^{0}(\mathbb{T E} ; \mathbb{T M})$

Fiber bundles

- E, M manifolds

Fiber bundles

## Math3



- Fiber bundle projection:

Fiber bundles

## Math3

Fiber bundles

- E, M manifolds
- Fiber bundle projection:
 $\pi_{\mathrm{M}, \mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{M})$ surjective submersion
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- Base space: $\mathbb{M}$
- Fiber manifold: $\left(\boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}}(\mathbf{x})\right)^{-1}$ based at $\mathbf{x} \in \mathbb{M}$
- Tangent bundle $\quad T \pi_{\mathrm{M}, \mathrm{E}} \in \mathrm{C}^{0}(\mathbb{T E} ; \mathbb{T M})$
- Vertical tangent subbundle $T \pi_{\mathbb{M}, \mathrm{E}} \in \mathrm{C}^{0}(\mathbb{V E} ; \mathbb{T M})$ with: $\delta \mathbf{e} \in \mathbb{V E} \subset \mathbb{T E} \quad \Longrightarrow \quad T_{\mathrm{e}} \boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}} \cdot \delta \mathbf{e}=0$

Fiber bundles

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Trivial and non-trivial fiber bundles


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Listing-Möbius strip


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## Sections of fiber bundles



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Sections of fiber bundles

- Fiber bundle $\quad \pi_{\mathbb{M}, \mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{M})$



## NLCM

## Math5

Sections of fiber bundles

- Fiber bundle $\quad \boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{M})$
- Sections

$\boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}} \circ \mathbf{S}_{\mathrm{E}, \mathbb{M}}=\mathrm{ID}_{\mathbb{M}}$

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Sections of fiber bundles

- Fiber bundle

$$
\boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{M})
$$

$$
\mathbf{s}_{\mathrm{E}, \mathbb{M}} \in \mathrm{C}^{1}(\mathbb{M} ; \mathrm{E}), \quad \boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}} \circ \mathbf{s}_{\mathrm{E}, \mathbb{M}}=\mathrm{ID}_{\mathbb{M}}
$$

- Tangent v.f.
- Sections

$$
\mathbf{v}_{\mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{T} \mathrm{E}), \quad \boldsymbol{\tau}_{\mathrm{E}} \circ \mathbf{v}_{\mathrm{E}}=\mathrm{ID}_{\mathrm{E}}
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## NLCM

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Sections of fiber bundles

- Fiber bundle

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$$

- Vertical tangent sections $\quad T \boldsymbol{\pi}_{\mathrm{M}, \mathrm{E}} \circ \mathbf{v}_{\mathrm{E}}=0$
- Sections

$$
\begin{aligned}
& \mathbf{s}_{\mathrm{E}, \mathbb{M}} \in \mathrm{C}^{1}(\mathbb{M} ; \mathrm{E}), \\
& \mathbf{v}_{\mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{T} \mathrm{E}),
\end{aligned}
$$



## Math5

Sections of fiber bundles

- Fiber bundle

$$
\boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{M})
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- Tangent v.f.

$$
\begin{array}{ll}
\mathbf{s}_{\mathrm{E}, \mathbb{M}} \in \mathrm{C}^{1}(\mathbb{M} ; \mathrm{E}), & \boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}} \circ \mathbf{s}_{\mathrm{E}, \mathbb{M}}=\mathrm{ID}_{\mathbb{M}} \\
\mathbf{v}_{\mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{T}), & \boldsymbol{\tau}_{\mathrm{E}} \circ \mathbf{v}_{\mathrm{E}}=\mathrm{ID}_{\mathrm{E}}
\end{array}
$$

- Sections

$$
T \boldsymbol{\pi}_{\mathrm{M}, \mathrm{E}} \circ \mathbf{v}_{\mathrm{E}}=0
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- Vertical tangent sections $\quad T \boldsymbol{\pi}_{\mathrm{M}, \mathrm{E}} \circ \mathbf{v}_{\mathrm{E}}=0$

Sections of tangent and bi-tangent bundles


## Math5

Sections of fiber bundles

- Fiber bundle

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Sections of tangent and bi-tangent bundles

- Tangent vector fields:

$$
\mathbf{v} \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{T} \mathbb{M}): \boldsymbol{\tau}_{\mathbb{M}} \circ \mathbf{v}=\operatorname{ID}_{\mathbb{M}}
$$

## Parallel transport

## Math5

Sections of fiber bundles

- Fiber bundle

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\boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{M})
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- Sections

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## Sections of tangent and bi-tangent bundles

- Tangent vector fields:

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\mathbf{v} \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{T} \mathbb{M}): \boldsymbol{\tau}_{\mathbb{M}} \circ \mathbf{v}=\operatorname{ID}_{\mathbb{M}}
$$

- Bi-tangent vector fields:

$$
\mathbf{X} \in \mathrm{C}^{1}(\mathbb{T M} ; \mathbb{T} \mathbb{T} \mathbb{M}): \boldsymbol{\tau}_{\mathbb{T} M} \circ \mathbf{X}=\mathrm{ID}_{\mathbb{T} M}
$$

## Math5

Sections of fiber bundles

- Fiber bundle

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\boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathbb{M})
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- Sections

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\mathbf{s}_{\mathrm{E}, \mathbb{M}} \in \mathrm{C}^{1}(\mathbb{M} ; \mathrm{E}), \quad \boldsymbol{\pi}_{\mathbb{M}, \mathrm{E}} \circ \mathbf{s}_{\mathrm{E}, \mathbb{M}}=\mathrm{ID}_{\mathbb{M}}
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## Sections of tangent and bi-tangent bundles

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- Bi-tangent vector fields:

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\mathbf{X} \in \mathrm{C}^{1}(\mathbb{T M} ; \mathbb{T} \mathbb{T} \mathbb{M}): \boldsymbol{\tau}_{\mathbb{T} M} \circ \mathbf{X}=\mathrm{ID}_{\mathbb{T} M}
$$

- Vertical bi-tangent vectors

$$
\mathbf{X} \in \operatorname{Ker} \boldsymbol{T}_{\mathbf{v}} \boldsymbol{\tau}_{\mathbb{M}}
$$

## Math6

## Tensor spaces

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Tensor spaces

- Covariant $\quad \mathbf{s}_{\mathbf{x}}^{\text {Cov }} \in \operatorname{Cov}_{\mathbf{x}}(\mathbb{T M})=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M}^{2} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M} ; \mathbb{T}_{\mathbf{x}}^{*} \mathbb{M}\right)$


## NLCM

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A basic question

## Basic

Tantent spaces
Tangent functor
Fiber bundles
Trivial and non-trivial
fiber bundles
Sections
Tensor bundle and sections

## Math6

Tensor spaces

- Covariant

$$
\mathbf{s}_{\mathbf{x}}^{\operatorname{Cov}} \in \operatorname{Cov}_{\mathbf{x}}(\mathbb{T} \mathbb{M})=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M}^{2} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M} ; \mathbb{T}_{\mathbf{x}}^{*} \mathbb{M}\right)
$$

- Contravariant $\mathbf{s}_{\mathbf{x}}^{\text {Con }} \in \operatorname{Con}_{\mathbf{x}}(\mathbb{T M})=L\left(\mathbb{T}_{\mathbf{x}}^{*} \mathbb{M}^{2} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathbf{x}}^{*} \mathbb{M} ; \mathbb{T}_{\mathbf{x}} \mathbb{M}\right)$


## NLCM

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## Basic

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Tensor spaces

- Covariant

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\mathbf{s}_{\mathbf{x}}^{\operatorname{Cov}} \in \operatorname{Cov}_{\mathbf{x}}(\mathbb{T} \mathbb{M})=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M}^{2} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M} ; \mathbb{T}_{\mathbf{x}}^{*} \mathbb{M}\right)
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- Mixed $\quad \mathbf{s}_{\mathbf{x}}^{\operatorname{Mix}} \in \operatorname{Mix}_{\mathbf{x}}(\mathbb{T M})=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M}, \mathbb{T}_{\mathbf{x}}^{*} \mathbb{M} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M} ; \mathbb{T}_{\mathbf{x}} \mathbb{M}\right)$


## Math6

## Tensor spaces

- Covariant $\quad \mathbf{s}_{\mathbf{x}}^{\text {Cov }} \in \operatorname{Cov}_{\mathbf{x}}(\mathbb{T M})=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M}^{2} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M} ; \mathbb{T}_{\mathbf{x}}^{*} \mathbb{M}\right)$
- Contravariant $\mathbf{s}_{\mathrm{x}}^{\text {Con }} \in \operatorname{Con}_{\mathrm{x}}(\mathbb{T M})=L\left(\mathbb{T}_{\mathrm{x}}^{*} \mathbb{M}^{2} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathrm{x}}^{*} \mathbb{M} ; \mathbb{T}_{\mathbf{x}} \mathbb{M}\right)$
- Mixed $\quad \mathbf{s}_{\mathbf{x}}^{\mathrm{Mix}} \in \operatorname{Mix}_{\mathbf{x}}(\mathbb{T M})=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M}, \mathbb{T}_{\mathbf{x}}^{*} \mathbb{M} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathrm{x}} \mathbb{M} ; \mathbb{T}_{\mathrm{x}} \mathbb{M}\right)$
- with the alteration rules:

$$
\mathbf{s}_{\mathrm{x}}^{\mathrm{Cov}}=\mathbf{g}_{\mathrm{x}} \circ \mathbf{s}_{\mathrm{x}}^{\mathrm{MIX}}, \quad \mathbf{s}_{\mathrm{x}}^{\mathrm{CoN}}=\mathbf{s}_{\mathrm{x}}^{\mathrm{Mix}} \circ \mathbf{g}_{\mathrm{x}}^{-1}
$$

## Math6

Tensor spaces

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## Tensor bundles and sections

## Math6

Tensor spaces

- Covariant $\quad \mathbf{s}_{\mathbf{x}}^{\text {Cov }} \in \operatorname{Cov}_{\mathbf{x}}(\mathbb{T} \mathbb{M})=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M}^{2} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M} ; \mathbb{T}_{\mathbf{x}}^{*} \mathbb{M}\right)$
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$$

## Tensor bundles and sections

- Tensor bundle

$$
\boldsymbol{\tau}_{\mathbb{M}}^{\mathrm{TEns}} \in \mathrm{C}^{1}(\operatorname{Tens}(\mathbb{T M}) ; \mathbb{M})
$$

## Math6

Tensor spaces

- Covariant $\quad \mathbf{s}_{\mathbf{x}}^{\text {Cov }} \in \operatorname{Cov}_{\mathbf{x}}(\mathbb{T M})=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M}^{2} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M} ; \mathbb{T}_{\mathbf{x}}^{*} \mathbb{M}\right)$
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\mathbf{s}_{\mathrm{x}}^{\mathrm{Cov}}=\mathbf{g}_{\mathrm{x}} \circ \mathbf{s}_{\mathrm{x}}^{\mathrm{MIX}}, \quad \mathbf{s}_{\mathrm{x}}^{\mathrm{CoN}}=\mathbf{s}_{\mathrm{x}}^{\mathrm{Mix}} \circ \mathbf{g}_{\mathrm{x}}^{-1}
$$

## Tensor bundles and sections

- Tensor bundle

$$
\boldsymbol{\tau}_{\mathbb{M}}^{\mathrm{TEns}} \in \mathrm{C}^{1}(\operatorname{Tens}(\mathbb{T M}) ; \mathbb{M})
$$

- Tensor field

$$
\mathbf{s}_{\mathbb{M}}^{\mathrm{TENS}} \in \mathrm{C}^{1}(\mathbb{M} ; \operatorname{Tens}(\mathbb{T M}))
$$

## Math6

## Tensor spaces

- Covariant $\quad \mathbf{s}_{\mathbf{x}}^{\operatorname{Cov}} \in \operatorname{Cov}_{\mathbf{x}}(\mathbb{T M})=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M}^{2} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathbf{x}} \mathbb{M} ; \mathbb{T}_{\mathbf{x}}^{*} \mathbb{M}\right)$
- Contravariant $\mathbf{s}_{\mathbf{x}}^{\text {Con }} \in \operatorname{Con}_{\mathbf{x}}(\mathbb{T M})=L\left(\mathbb{T}_{\mathbf{x}}^{*} \mathbb{M}^{2} ; \mathcal{R}\right)=L\left(\mathbb{T}_{\mathbf{x}}^{*} \mathbb{M} ; \mathbb{T}_{\mathbf{x}} \mathbb{M}\right)$
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$$

## Tensor bundles and sections

- Tensor bundle $\boldsymbol{\tau}_{\mathbb{M}}^{\text {TEns }} \in \mathrm{C}^{1}(\operatorname{Tens}(\mathbb{T M}) ; \mathbb{M})$
- Tensor field $\quad \mathbf{s}_{\mathbb{M}}^{\mathrm{TENS}} \in \mathrm{C}^{1}(\mathbb{M} ; \operatorname{Tens}(\mathbb{T M}))$
- with: $\boldsymbol{\tau}_{\mathbb{M}}^{\mathrm{Tens}} \circ \mathbf{s}_{\mathbb{M}}^{\mathrm{Tens}}=\mathrm{ID}_{\mathbb{M}}$


## Math7

## Push and pull

## NLCM

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Metric measurements

Metric theory
Events manifold fibrations

## Math7

Push and pull
Given a map $\zeta \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{N})$

- Pull-back of a scalar field

$$
f: \mathbb{N} \mapsto \operatorname{FuN}(\mathbb{N}) \quad \mapsto \quad \zeta \downarrow f: \mathbb{M} \mapsto \operatorname{FuN}(\mathbb{M})
$$

defined by:

$$
(\zeta \downarrow f)_{\mathrm{x}}:=\zeta \downarrow f_{\zeta(\mathrm{x})}:=f_{\zeta(\mathrm{x})} \in \operatorname{FuN}_{\mathrm{x}}(\mathbb{M}) .
$$

## NLCM

## Math7

Push and pull
Given a map $\zeta \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{N})$

- Pull-back of a scalar field

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$$

defined by:

$$
(\zeta \downarrow f)_{\mathrm{x}}:=\zeta \downarrow f_{\zeta(\mathrm{x})}:=f_{\zeta(\mathrm{x})} \in \operatorname{FUN}_{\mathrm{x}}(\mathbb{M})
$$

- Push-forward of a tangent vector field

$$
\mathbf{v} \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{T} \mathbb{M}) \quad \mapsto \quad \boldsymbol{\zeta} \uparrow \mathbf{v}: \mathbb{N} \mapsto \mathbb{T} \mathbb{N}
$$

defined by:

$$
(\zeta \uparrow \mathbf{v})_{\zeta(\mathrm{x})}:=\zeta \uparrow \mathbf{v}_{\mathrm{x}}=T_{\mathrm{x}} \zeta \cdot \mathbf{v}_{\mathrm{x}} \in \mathbb{T}_{\zeta(\mathrm{x})} \mathbb{N} .
$$

## Math8

## Push and pull of tensor fields

## NLCM

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## Math8

Push and pull of tensor fields

- Covectors

$$
\left\langle\boldsymbol{\zeta} \downarrow \mathbf{v}_{\zeta(\mathbf{x})}^{*}, \mathbf{v}_{\mathbf{x}}\right\rangle=\left\langle\mathbf{v}_{\boldsymbol{\zeta}(\mathbf{x})}^{*}, \boldsymbol{\zeta} \uparrow \mathbf{v}_{\mathbf{x}}\right\rangle=\left\langle T_{\zeta(\mathbf{x})}^{*} \boldsymbol{\zeta} \circ \mathbf{v}_{\zeta(\mathbf{x})}^{*}, \mathbf{v}_{\mathbf{x}}\right\rangle
$$

## NLCM

## Math8

Push and pull of tensor fields

- Covectors

$$
\left\langle\boldsymbol{\zeta} \backslash \mathbf{v}_{\zeta(\mathbf{x})}^{*}, \mathbf{v}_{\mathbf{x}}\right\rangle=\left\langle\mathbf{v}_{\boldsymbol{\zeta}(\mathrm{x})}^{*}, \boldsymbol{\zeta} \uparrow \mathbf{v}_{\mathbf{x}}\right\rangle=\left\langle T_{\zeta(\mathbf{x})}^{*} \boldsymbol{\zeta} \circ \mathbf{v}_{\zeta(\mathbf{x})}^{*}, \mathbf{v}_{\mathbf{x}}\right\rangle
$$

- Covariant tensors

$$
\zeta \downarrow \mathbf{s}_{\zeta(\mathrm{x})}^{\mathrm{Cov}}=T_{\zeta(\mathrm{x})}^{*} \boldsymbol{\zeta} \circ \mathbf{s}_{\zeta(\mathrm{x})}^{\mathrm{Cov}} \circ T_{\mathrm{x}} \boldsymbol{\zeta} \in \operatorname{Cov}(\mathbb{T M})_{\mathrm{x}}
$$

## Math8

Push and pull of tensor fields

- Covectors

$$
\left\langle\boldsymbol{\zeta} \backslash \mathbf{v}_{\zeta(\mathbf{x})}^{*}, \mathbf{v}_{\mathbf{x}}\right\rangle=\left\langle\mathbf{v}_{\boldsymbol{\zeta}(\mathrm{x})}^{*}, \boldsymbol{\zeta} \uparrow \mathbf{v}_{\mathbf{x}}\right\rangle=\left\langle T_{\zeta(\mathbf{x})}^{*} \boldsymbol{\zeta} \circ \mathbf{v}_{\zeta(\mathbf{x})}^{*}, \mathbf{v}_{\mathbf{x}}\right\rangle
$$

- Covariant tensors

$$
\zeta \backslash \mathbf{s}_{\zeta(\mathrm{x})}^{\mathrm{Cov}}=T_{\zeta(\mathrm{x})}^{*} \zeta \circ \mathbf{s}_{\zeta(\mathrm{x})}^{\mathrm{Cov}} \circ T_{\mathrm{x}} \zeta \in \operatorname{Cov}(\mathbb{T M})_{\mathrm{x}}
$$

- Contravariant tensors

$$
\boldsymbol{\zeta} \uparrow \mathbf{s}_{\mathrm{x}}^{\mathrm{CoN}}=T_{\mathrm{x}} \boldsymbol{\zeta} \circ \mathbf{s}_{\mathrm{x}}^{\mathrm{CoN}} \circ T_{\zeta(\mathrm{x})}^{*} \boldsymbol{\zeta} \in \operatorname{Con}(\mathbb{T} \mathbb{N})_{\zeta(\mathrm{x})}
$$

## Math8

Push and pull of tensor fields

- Covectors

$$
\left\langle\zeta \backslash \mathbf{v}_{\zeta(\mathbf{x})}^{*}, \mathbf{v}_{\mathbf{x}}\right\rangle=\left\langle\mathbf{v}_{\boldsymbol{\zeta}(\mathrm{x})}^{*}, \boldsymbol{\zeta} \uparrow \mathbf{v}_{\mathbf{x}}\right\rangle=\left\langle T_{\zeta(\mathbf{x})}^{*} \boldsymbol{\zeta} \circ \mathbf{v}_{\zeta(\mathrm{x})}^{*}, \mathbf{v}_{\mathbf{x}}\right\rangle
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- Covariant tensors

$$
\zeta \backslash \mathbf{s}_{\zeta(\mathrm{x})}^{\mathrm{Cov}}=T_{\zeta(\mathrm{x})}^{*} \zeta \circ \mathbf{s}_{\zeta(\mathrm{x})}^{\mathrm{Cov}} \circ T_{\mathrm{x}} \zeta \in \operatorname{Cov}(\mathbb{T M})_{\mathrm{x}}
$$

- Contravariant tensors

$$
\boldsymbol{\zeta} \uparrow \mathbf{s}_{\mathrm{x}}^{\mathrm{CoN}}=T_{\mathrm{x}} \boldsymbol{\zeta} \circ \mathbf{s}_{\mathrm{x}}^{\mathrm{CoN}} \circ T_{\zeta(\mathrm{x})}^{*} \boldsymbol{\zeta} \in \operatorname{Con}(\mathbb{T} \mathbb{N})_{\zeta(\mathrm{x})}
$$

- Mixed tensors

$$
\boldsymbol{\zeta} \uparrow \mathbf{s}_{\mathrm{x}}^{\mathrm{MIX}}=T_{\mathrm{x}} \zeta \circ \mathbf{s}_{\mathrm{x}}^{\mathrm{MIX}} \circ T_{\zeta(\mathrm{x})} \zeta^{-1} \in \operatorname{Mix}(\mathbb{T})_{\zeta(\mathrm{x})}
$$

## Math9

Parallel transport along a curve $\mathbf{c} \in \mathrm{C}^{1}([a, b] ; \mathbb{M})$

## NLCM

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Fiber buncles
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Parallel transport
Derivatives
Key contributions

## Math9

Parallel transport along a curve $\mathbf{c} \in \mathrm{C}^{1}([a, b] ; \mathbb{M})$

- Vector fields

$$
\begin{aligned}
& \mathbf{x}=\mathbf{c}(\mu), \quad \mathbf{v}_{\mathbf{x}} \in \mathbb{T}_{\mathbf{x}} \mathbb{M} \mapsto \mathbf{c}_{\lambda, \mu} \Uparrow \mathbf{v}_{\mathbf{x}} \in \mathbb{T}_{\mathbf{c}(\lambda)} \mathbb{M} \\
& \\
& \mathbf{c}_{\mu, \mu} \Uparrow \mathbf{v}_{\mathbf{x}}=\mathbf{v}_{\mathbf{x}} \\
& \\
& \mathbf{c}_{\lambda, \mu} \Uparrow \circ \mathbf{c}_{\mu, \nu} \Uparrow=\mathbf{c}_{\lambda, \nu} \Uparrow
\end{aligned}
$$

## Math9

Parallel transport along a curve $\mathbf{c} \in \mathrm{C}^{1}([a, b] ; \mathbb{M})$

- Vector fields

$$
\begin{array}{ll}
\mathbf{x}=\mathbf{c}(\mu), & \mathbf{v}_{\mathbf{x}} \in \mathbb{T}_{\mathbf{x}} \mathbb{M} \quad \mapsto \quad \mathbf{c}_{\lambda, \mu} \Uparrow \mathbf{v}_{\mathbf{x}} \in \mathbb{T}_{\mathbf{c}(\lambda)} \mathbb{M} \\
& \mathbf{c}_{\mu, \mu} \Uparrow \mathbf{v}_{\mathbf{x}}=\mathbf{v}_{\mathbf{x}} \\
& \mathbf{c}_{\lambda, \mu} \Uparrow \circ \mathbf{c}_{\mu, \nu} \Uparrow=\mathbf{c}_{\lambda, \nu} \Uparrow
\end{array}
$$

- Covector fields $\mathbf{v}_{\mathrm{x}}^{*} \in \mathbb{T}_{\mathrm{x}}^{*} \mathbb{M}$ (by naturality)

$$
\left\langle\mathbf{c}_{\lambda, \mu} \Uparrow \mathbf{v}_{\mathbf{x}}^{*}, \mathbf{c}_{\lambda, \mu} \Uparrow \mathbf{v}_{\mathbf{x}}\right\rangle=\mathbf{c}_{\lambda, \mu} \Uparrow\left\langle\mathbf{v}_{\mathbf{x}}^{*}, \mathbf{v}_{\mathbf{x}}\right\rangle
$$

- Tensor fields (by naturality)


## Math9

Parallel transport along a curve $\mathbf{c} \in \mathrm{C}^{1}([a, b] ; \mathbb{M})$

- Vector fields

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- Tensor fields (by naturality)



## Math9

Parallel transport along a curve $\mathbf{c} \in \mathrm{C}^{1}([a, b] ; \mathbb{M})$

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\mathbf{x}=\mathbf{c}(\mu), & \mathbf{v}_{\mathbf{x}} \in \mathbb{T}_{\mathbf{x}} \mathbb{M} \quad \mapsto \quad \mathbf{c}_{\lambda, \mu} \Uparrow \mathbf{v}_{\mathbf{x}} \in \mathbb{T}_{\mathbf{c}(\lambda)} \mathbb{M} \\
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\left\langle\mathbf{c}_{\lambda, \mu} \Uparrow \mathbf{v}_{\mathbf{x}}^{*}, \mathbf{c}_{\lambda, \mu} \Uparrow \mathbf{v}_{\mathbf{x}}\right\rangle=\mathbf{c}_{\lambda, \mu} \Uparrow\left\langle\mathbf{v}_{\mathbf{x}}^{*}, \mathbf{v}_{\mathbf{x}}\right\rangle
$$

- Tensor fields (by naturality)


Tullio Levi-Civita (1873-1941)

## Math10

Derivatives of a tensor field
$s \in \mathrm{C}^{1}(\mathbb{M} ; \operatorname{Tens}(\mathbb{T M}))$
along the flow of a tangent vector field

## Math10

Derivatives of a tensor field
$s \in \mathrm{C}^{1}(\mathbb{M} ; \operatorname{Tens}(\mathbb{T M}))$
along the flow of a tangent vector field

- Tangent vector fields and Flows

$$
\begin{aligned}
& \mathbf{v} \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{T M}) \\
& \mathbf{v}:=\partial_{\lambda=0} \mathbf{F I}_{\lambda}^{v}
\end{aligned}
$$

## Math10

Derivatives of a tensor field
$s \in \mathrm{C}^{1}(\mathbb{M} ; \boldsymbol{T e n s}(\mathbb{T M}))$
along the flow of a tangent vector field

- Tangent vector fields and Flows

$$
\begin{aligned}
& \mathbf{v} \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{T M}) \quad \mathbf{F I}_{\lambda}^{v} \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{M}) \\
& \mathbf{v}:=\partial_{\lambda=0} \mathrm{FI}_{\lambda}^{v}
\end{aligned}
$$

- Lie derivative - LD

$$
\mathcal{L}_{\mathbf{v}} \mathbf{s}:=\partial_{\lambda=0} \mathbf{F I}_{\lambda}^{\mathbf{v}} \downarrow\left(\mathbf{s} \circ \mathbf{F} \mathbf{I}_{\lambda}^{\mathbf{v}}\right)
$$

## Math10

Derivatives of a tensor field
$s \in \mathrm{C}^{1}(\mathbb{M} ; \boldsymbol{T e n s}(\mathbb{T M}))$
along the flow of a tangent vector field

- Tangent vector fields and Flows

$$
\begin{aligned}
& \mathbf{v} \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{T M}) \quad \mathbf{F I}_{\lambda}^{v} \in \mathrm{C}^{1}(\mathbb{M} ; \mathbb{M}) \\
& \mathbf{v}:=\partial_{\lambda=0} \mathrm{FI}_{\lambda}^{\mathrm{v}}
\end{aligned}
$$

- Lie derivative - LD

$$
\mathcal{L}_{\mathbf{v}} \mathbf{s}:=\partial_{\lambda=0} \mathbf{F l}_{\lambda}^{\mathrm{V}} \downarrow\left(\mathbf{s} \circ \mathbf{F} \mathbf{I}_{\lambda}^{\mathbf{v}}\right)
$$

- Parallel derivative - PD

$$
\nabla_{\mathbf{v}} \mathbf{s}:=\partial_{\lambda=0} \mathbf{F I}_{\lambda}^{\mathbf{v}} \Downarrow\left(\mathbf{s} \circ \mathbf{F l}_{\lambda}^{\mathbf{v}}\right)
$$

NLCM: Nonlinear Continuum Mechanics
Key contributions

## NLCM: Nonlinear Continuum Mechanics

## Key contributions

C. Truesdell \& W. Noll The non-linear field theories of mechanics Handbuch der Physik, Springer (1965)
C. Truesdell $A$ first Course in Rational Continuum Mechanics

Second Ed., Academic Press, New-York (1991). First Ed. (1977).

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J.E. Marsden Lectures on Geometric Methods in Mathematical Physics, SIAM, Philadelphia, PA (1981), on line version July 22 (2009)
J.E. Marsden \& T.J.R. Hughes Mathematical Foundations of Elasticity Prentice-Hall, Redwood City, Cal. (1983)
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G. Romano \& R. Barretta Covariant hypo-elasticity

Eur. J. Mech. A-Solids 30 (2011) 1012-1023
G. Romano, R. Barretta, M. Diaco Basic Geometric Issues in Non-Linear Continuum Mechanics, preprint (2011).

## NLCM: Nonlinear Continuum Mechanics

## NLCM: Nonlinear Continuum Mechanics

How to play the game according to a full geometric approach

## NLCM

Dralemomena
A basic question

## NLCM: Nonlinear Continuum Mechanics

How to play the game
according to a full geometric approach
Kinematics

## NLCM

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- Events manifold: E - four dimensional RiEmANn manifold


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- Observer split into space-time: $\gamma: \mathrm{E} \mapsto \mathcal{S} \times /$


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- Events manifold: E - four dimensional Riemann manifold
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- time is absolute (Classical Mechanics)


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- Events manifold: E - four dimensional Riemann manifold
- Observer split into space-time: $\gamma: \mathrm{E} \mapsto \mathcal{S} \times I$
- time is absolute (Classical Mechanics)
- distance between simultaneous events $\mapsto$ space-metric
- distance between localized events $\mapsto$ time-metric


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## lenght of symplex's edges

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## lenght of symplex's edges

- Norm axioms


$$
\|\mathbf{a}\| \geq 0, \quad\|\mathbf{a}\|=0 \quad \Longrightarrow \quad \mathbf{a}=0
$$

$$
\|\mathbf{a}\|+\|\mathbf{b}\| \geq\|\mathbf{c}\| \quad \text { triangle inequality }
$$

$$
\|\alpha \mathbf{a}\|=|\alpha|\|\mathbf{a}\|
$$

## Math11

## lenght of symplex's edges

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& \|\alpha \mathbf{a}\|=|\alpha|\|\mathbf{a}\|
\end{aligned}
$$

- Parallelogram rule


$$
\|\mathbf{a}+\mathbf{b}\|^{2}+\|\mathbf{a}-\mathbf{b}\|^{2}=2\left[\|\mathbf{a}\|^{2}+\|\mathbf{b}\|^{2}\right]
$$

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The metric tensor

- Theorem (Fréchet - von Neumann - Jordan)


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The metric tensor

- Theorem (Fréchet - von Neumann - Jordan)

$$
\mathbf{g}(\mathbf{a}, \mathbf{b}):=\frac{1}{4}\left[\|\mathbf{a}+\mathbf{b}\|^{2}-\|\mathbf{a}-\mathbf{b}\|^{2}\right]
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$$



Maurice René Fréchet (1878-1973)

$$
)^{2}=\operatorname{det}\left[\begin{array}{ccc}
\mathbf{g}\left(\mathbf{e}_{1}, \mathbf{e}_{1}\right) \cdots & \mathbf{g}\left(\mathbf{e}_{1}, \mathbf{e}_{3}\right) \\
\ldots & \cdots & \cdots \\
\mathbf{g}\left(\mathbf{e}_{3}, \mathbf{e}_{1}\right) & \cdots & \mathbf{g}\left(\mathbf{e}_{3}, \mathbf{e}_{3}\right)
\end{array}\right]
$$

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John von Neumann (1903-1957)

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Pascual Jordan (1902-1980)

## Math12

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$$



Kosaku Yosida (1909-1990)

## Events manifold fibrations

## Events manifold fibrations

- Time and space fibrations: $\gamma: \mathrm{E} \mapsto \mathcal{S} \times I$ (observer)

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\begin{aligned}
\pi_{I, \mathrm{E}} & =\pi_{I,(\mathcal{S} \times I)} \circ \gamma \\
\boldsymbol{\pi}_{\mathcal{S}, \mathrm{E}} & =\boldsymbol{\pi}_{\mathcal{S},(\mathcal{S} \times I)} \circ \gamma
\end{aligned}
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- Space-time metric: $\mathbf{g}_{\mathrm{E}}:=\pi_{\mathcal{S}, \mathrm{E}} \downarrow \mathbf{g}_{\mathcal{S}}+\pi_{l, \mathrm{E}} \downarrow \mathbf{g}_{/}$


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- Time-vertical subbundle: spatial vectors

$$
\mathbf{v} \in \mathbb{V}_{\mathbf{e}} \mathrm{E} \quad \Longleftrightarrow \quad T_{\mathbf{e}} \pi_{I, \mathrm{E}} \cdot \mathbf{v}=0
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- $\mathbf{v}_{\mathbf{e}} \in \mathbb{V}_{\mathrm{e}} \mathrm{E} \quad \Longleftrightarrow \quad \gamma \uparrow \mathbf{v}_{\mathbf{e}}=\left(v_{\mathrm{x}, t}, 0_{t}\right) \in \mathbb{T}_{\mathrm{x}} \mathcal{S} \times \mathbb{T}_{t} /$


## Trajectory



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## Trajectory



- Trajectory $\mapsto$ a manifold $\mathcal{T}_{\varphi}$ with injective immersion in the events time-bundle: $\mathbf{i}_{\mathrm{E}, \mathcal{I}_{\varphi}} \in \mathrm{C}^{1}\left(\mathcal{T}_{\varphi} ; \mathrm{E}\right)$


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- Trajectory time-fibration $\quad \pi_{I, \mathcal{T}_{\varphi}}:=\pi_{I, \mathrm{E}} \circ \mathbf{i}_{\mathrm{E}, \mathcal{T}_{\varphi}}$
- time bundle $\mapsto$ fibers: body placements $\Omega_{t}$

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\pi_{l, \mathcal{T}_{\varphi}}:=\pi_{l, \mathrm{E}} \circ \mathbf{i}_{\mathrm{E}, \mathcal{T}_{\varphi}}
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## Evolution

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## Evolution

- Evolution operator $\varphi^{\mathcal{T}} \varphi$

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- Evolution operator $\varphi^{\mathcal{T}_{\varphi}}$
- Displacements: diffeomorphisms between placements

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\boldsymbol{\varphi}_{\tau, t}^{\mathcal{T}_{\varphi}} \in \mathrm{C}^{1}\left(\boldsymbol{\Omega}_{t} ; \boldsymbol{\Omega}_{\tau}\right), \quad \tau, t \in I
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- Law of determinism (Chapman-Kolmogorov):

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- Trajectory speed:

$$
\mathbf{v}_{\mathcal{T}_{\varphi}}\left(\mathbf{e}_{t}\right):=\partial_{\tau=t} \varphi_{\tau, t}^{\mathcal{T}_{\varphi}}\left(\mathbf{e}_{t}\right) \quad \Longrightarrow \quad T_{\mathbf{e}} \pi_{I, \mathcal{T}_{\varphi}} \cdot \mathbf{v}_{\mathcal{T}_{\varphi}}\left(\mathbf{e}_{t}\right)=1_{t}
$$

## Body and particles

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## Body and particles

- Equivalence relation on the trajectory:

$$
\left(\mathbf{e}_{1}, \mathbf{e}_{2}\right) \in \mathcal{T}_{\varphi} \times \mathcal{T}_{\varphi}: \mathbf{e}_{2}=\varphi_{t_{2}, t_{1}}^{\mathcal{T}_{\varphi}}\left(\mathbf{e}_{1}\right)
$$

with $t_{i}=\pi_{l, \mathrm{E}}\left(\mathbf{e}_{i}\right), \quad i=1,2$.

## Body and particles

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& \qquad \quad\left(\mathbf{e}_{1}, \mathbf{e}_{2}\right) \in \mathcal{T}_{\boldsymbol{\varphi}} \times \mathcal{T}_{\boldsymbol{\varphi}}: \mathbf{e}_{2}=\varphi_{t_{2}, t_{1}}^{\mathcal{T}_{\boldsymbol{\varphi}}}\left(\mathbf{e}_{1}\right) \\
& \text { with } t_{i}=\boldsymbol{\pi}_{l, \mathrm{E}}\left(\mathbf{e}_{i}\right), \quad i=1,2 \\
& \text { Body }=\text { quotient manifold (foliation) }
\end{aligned}
$$

## Body and particles

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Body $=$ quotient manifold (foliation)
Particles $=$ equivalence classes (folia)

## Body and particles

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$$

with $t_{i}=\pi_{l, \mathrm{E}}\left(\mathbf{e}_{i}\right), \quad i=1,2$.
Body $=$ quotient manifold (foliation)
Particles $=$ equivalence classes (folia)

- mass conservation

$$
\begin{aligned}
& \quad \int_{\Omega_{t_{1}}} \mathbf{m}_{\tau_{\varphi}, t_{1}}=\int_{\Omega_{\mathrm{t}_{2}}} \mathbf{m}_{\tau_{\varphi}, t_{2}} \Longleftrightarrow \mathcal{L}_{\mathbf{v}_{\tau_{\varphi}}} \mathbf{m}_{\tau_{\varphi}}=0 \\
& \mathbf{m}_{\tau_{\varphi}} \in \mathrm{C}^{1}\left(\mathcal{T}_{\varphi} ; \operatorname{VoL}\left(\mathbb{T} \mathcal{T}_{\varphi}\right)\right) \text { mass form }
\end{aligned}
$$

## Tensor fields in NLCM

## Tensor fields in NLCM

| Space-time fields | $\mathbf{s}_{\mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \operatorname{TENS}(\mathbb{T} \mathrm{E}))$ | Space-time metric tensor | NLCM <br> Prolegon |
| :---: | :---: | :---: | :---: |
| Spatial fields | $\mathbf{s}_{\mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \operatorname{TENS}(\mathbb{V} \mathrm{E}))$ | Spatial metric tensor | Basic <br> Tangent <br> Tangent |

## Tensor fields in NLCM

| Space-time fields | $\mathbf{s}_{\mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \operatorname{TENS}(\mathbb{T} \mathrm{E}))$ | Space-time metric tensor |
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| Spatial fields | $\mathbf{s}_{\mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \operatorname{TENS}(\mathbb{V} \mathrm{E}))$ | Spatial metric tensor |
| Trajectory fields | $\mathbf{s}_{\mathcal{T}_{\varphi}} \in \mathrm{C}^{1}\left(\mathcal{T}_{\boldsymbol{\varphi}} ; \operatorname{TENS}\left(\mathbb{T} \mathcal{T}_{\varphi}\right)\right)$ | Trajectory metric, trajectory speed |
| Material fields | $\mathbf{s}_{\mathcal{T}_{\varphi}} \in \mathrm{C}^{1}\left(\mathcal{T}_{\varphi} ; \operatorname{TENS}\left(\mathbb{V} \mathcal{T}_{\varphi}\right)\right)$ | Stress, stressing, material metric, stretching. |

## Tensor fields in NLCM

| Space-time fields | $\mathbf{s}_{\mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \operatorname{TENS}(\mathbb{T} \mathrm{E}))$ | Space-time metric tensor |
| :---: | :---: | :---: |
| Spatial fields | $\mathbf{s}_{\mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \operatorname{TENs}(\mathbb{V} \mathrm{E}))$ | Spatial metric tensor |
| Trajectory fields | $\mathbf{s}_{\mathcal{I}_{\varphi}} \in \mathrm{C}^{1}\left(\mathcal{T}_{\varphi} ; \operatorname{TENS}\left(\mathbb{T} \mathcal{I}_{\varphi}\right)\right)$ | Trajectory metric, trajectory speed |
| Material fields | $\mathbf{s}_{\mathcal{T}_{\varphi}} \in \mathrm{C}^{1}\left(\mathcal{T}_{\boldsymbol{\varphi}} ; \operatorname{TENS}\left(\mathbb{V} \mathcal{I}_{\varphi}\right)\right)$ | Stress, stressing, material metric, stretching. |
| Trajectory-based space-time fields | $\mathbf{s}_{\mathrm{E}, \mathcal{T}_{\varphi}} \in \mathrm{C}^{1}\left(\mathcal{T}_{\varphi} ; \operatorname{TENS}(\mathbb{T E})\right)$ | Trajectory speed (immersed) |
| Trajectory-based spatial fields | $\mathbf{s}_{\mathrm{E}, \mathcal{I}_{\varphi}} \in \mathrm{C}^{1}\left(\mathcal{T}_{\varphi} ; \operatorname{TENS}(\mathbb{V E})\right)$ | Virtual velocity, acceleration, momentum, force |

## Covariance Paradigm

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## Covariance Paradigm

Material fields at different times along the trajectory must be compared by push along the material displacement. Material fields on push-related trajectories must be compared by push along the relative motion.

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Push and parallel transport along the motion

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## Time derivatives $=$ <br> derivatives along the flow of the trajectory speed

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derivatives along the flow of the trajectory speed

Lie time derivative - LTD

- Trajectory and material tensor field

$$
\dot{\mathbf{s}}_{\mathcal{T}_{\varphi}}:=\mathcal{L}_{\mathbf{v}_{\mathcal{T}_{\varphi}}} \mathbf{s}_{\mathcal{T}_{\varphi}}=\partial_{\lambda=0} \mathbf{F l}_{\lambda}^{\mathbf{v}_{\mathcal{T}_{\varphi}}} \downarrow\left(\mathbf{s}_{\mathcal{T}_{\varphi}} \circ \mathbf{F I}_{\lambda}^{\mathbf{v}_{\mathcal{T}_{\varphi}}}\right),
$$

## Time derivatives $=$

Lie time derivative - LTD

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\dot{\mathbf{s}}_{T_{\varphi}}:=\mathcal{L}_{\mathbf{v}_{T_{\varphi}}} \mathbf{s}_{\mathcal{T}_{\varphi}}=\partial_{\lambda=0} \mathbf{F I}_{\lambda}^{\mathbf{v}_{\varphi}} \downarrow\left(\mathbf{s}_{\tau_{\varphi}} \circ \mathbf{F l}_{\lambda}^{\mathbf{v}_{\boldsymbol{\tau}}}\right),
$$

## Material time-derivative - MTD

- Trajectory-based space-time and spatial fields

$$
\dot{\mathbf{s}}_{\mathrm{E}, \mathcal{T}_{\varphi}}:=\nabla_{\mathbf{v}_{\varphi}}^{\mathrm{E}} \mathbf{s}_{\mathrm{E}, \mathcal{T}_{\varphi}}=\partial_{\lambda=0} \mathbf{F}_{\lambda}^{\mathbf{v}_{\mathrm{E}, \mathcal{T}_{\varphi}}} \Downarrow^{\mathrm{E}}\left(\mathbf{s}_{\mathrm{E}, \mathcal{I}_{\varphi}} \circ \mathbf{F}_{\lambda}^{\mathbf{v}_{\tau_{\varphi}}}\right),
$$

with $\mathbf{v}_{\mathrm{E}, \tau_{\varphi}}:=\mathbf{i}_{\mathrm{E}, \mathcal{T}_{\varphi}} \uparrow \mathbf{v}_{\tau_{\varphi}}$.

## Rivers and Cogwheels

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## Rivers and Cogwheels

$$
\left(\mathcal{L}_{\mathbf{v}_{\mathcal{T}_{\varphi}}} \mathbf{s}_{\mathcal{T}_{\varphi}}\right)_{t}:=\partial_{\tau=t} \boldsymbol{\varphi}_{\tau, t} \downarrow\left(\mathbf{s}_{\mathcal{T}_{\varphi}, \tau} \circ \boldsymbol{\varphi}_{\tau, t}\right)=\partial_{\tau=t} \mathbf{s}_{\mathcal{T}_{\varphi}, \tau}+\mathcal{L}_{\boldsymbol{\pi}_{\mathcal{S}, \mathcal{T}_{\varphi}} \downarrow \mathbf{v}_{\mathcal{T}_{\varphi}}} \mathbf{s}_{\mathcal{T}_{\varphi}, t}
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$$

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Rivers and Cogwheels
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## Rivers and Cogwheels

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Gottfried Wilhelm von Leibniz (1646-1716)

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Both conditions are not fulfilled in solid mechanics, in general.

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## Acceleration

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MTD of the velocity field

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\end{aligned}
$$

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This is the celebrated Euler split formula, applicable only in special problems of hydrodynamics, where it was originally conceived.
This eventually led to the Navier-Stokes-St.Venant differential equation of motion in fluid-dynamics.

Trivial and non-trivial

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\dot{\varepsilon}_{\tau_{\varphi}, t}:=\frac{1}{2}\left(\mathcal{L}_{\mathbf{v}_{\tau_{\varphi}}} \mathbf{g}_{\tau_{\varphi}}\right)_{t}=\frac{1}{2} \partial_{\tau=t}\left(\varphi_{\tau, t} \backslash \mathbf{g}_{\tau_{\varphi}, \tau}\right)
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Leonhard Euler (1707-1783)


- Euler's formula (generalized)

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\frac{1}{2} \mathcal{L}_{\mathbf{v}_{\tau_{\varphi}}} \mathbf{g}_{\tau_{\varphi}}=\frac{1}{2} \nabla_{v_{\tau_{\varphi}}}^{\tau_{\varphi}} \mathbf{g}_{\tau_{\varphi}}+\operatorname{sym}\left(\mathbf{g}_{\tau_{\varphi}} \circ\left(\operatorname{ToRS}^{\tau_{\varphi}}+\nabla^{\mathcal{T}_{\varphi}}\right) \mathbf{v}_{\tau_{\varphi}}\right)
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$$

- Mixed form of the stretching tensor (standard):

$$
\mathbf{D}_{\mathcal{T}_{\varphi}}:=\mathbf{g}_{\mathcal{T}_{\varphi}}^{-1} \circ \frac{1}{2} \mathcal{L}_{\mathbf{v}_{\mathcal{T}_{\varphi}}} \mathbf{g}_{\mathcal{T}_{\varphi}}=\operatorname{sym}\left(\nabla^{\mathcal{T}_{\varphi}} \mathbf{v}_{\mathcal{T}_{\varphi}}\right)
$$

$$
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## Stress and stressing

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## Stress and stressing

- Stress: $\boldsymbol{\sigma}_{\mathcal{T}_{\boldsymbol{\varphi}}} \in \mathrm{C}^{1}\left(\mathcal{T}_{\boldsymbol{\varphi}} ; \operatorname{CoN}\left(\mathbb{V} \mathcal{T}_{\boldsymbol{\varphi}}\right)\right)$ in duality with the
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## Dralexamena

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The expression in terms of parallel derivative:

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- Treatments which do not adopt a full geometric approach do not even perceive the difficulties revealed by the previous investigation.


## Objective stress rate tensors

A sample of objective stress rate tensors

Co-rotational stress rate tensor, Zaremba (1903), Jaumann $(1906,1911)$, Prager (1960):

$$
\stackrel{\circ}{\mathbf{T}}=\dot{\mathbf{T}}-\mathbf{W} \mathbf{T}+\mathbf{T} \mathbf{W}
$$

with $\dot{\mathbf{T}}$ material time derivative.
Convective stress tensor rate, Zaremba (1903), Oldroyd (1950), Truesdell (1955), Sedov (1960), Truesdell \& Noll (1965):

$$
\stackrel{\Delta}{\mathbf{T}}=\dot{\mathbf{T}}+\mathbf{L}^{T} \mathbf{T}+\mathbf{T L}
$$

## Objective stress rate tensors

A sample of objective stress rate tensors

Co-rotational stress rate tensor, Zaremba (1903), Jaumann $(1906,1911)$, Prager (1960):

$$
\stackrel{\circ}{\mathbf{T}}=\dot{\mathbf{T}}-\mathbf{W} \mathbf{T}+\mathbf{T} \mathbf{W}
$$

with $\dot{\mathbf{T}}$ material time derivative.
Convective stress tensor rate, Zaremba (1903), Oldroyd (1950), Truesdell (1955), Sedov (1960), Truesdell \& Noll (1965):

$$
\stackrel{\Delta}{\mathbf{T}}=\dot{\mathbf{T}}+\mathbf{L}^{T} \mathbf{T}+\mathbf{T L}
$$

These formulas, and similar ones in literature, rely on the application of Leibniz rule and on taking the parallel derivative of the material stress tensor field according to the trajectory connection.

The lack of regularity that may prevent to take partial time derivatives and the lack of conservation of time-verticality by parallel transport, are not taken into account.

## Deformation gradient

The equivalence class of all material displacements whose tangent map have the common value:

$$
T_{\mathbf{x}} \boldsymbol{\varphi}_{\tau, t} \in L\left(\mathbb{T}_{\mathbf{x}} \boldsymbol{\Omega}_{t} ; \mathbb{T}_{\boldsymbol{\varphi}_{\tau, t}(\mathrm{x})} \boldsymbol{\Omega}_{\tau}\right)
$$

- is called the first jet of $\boldsymbol{\varphi}_{\tau, t}$ at $\mathbf{x} \in \boldsymbol{\Omega}_{t}$ in differential geometry
- and the relative deformation gradient in continuum mechanics.

The chain rule between tangent maps:

$$
\boldsymbol{T}_{\boldsymbol{\varphi}_{\tau, s}(\mathrm{x})} \boldsymbol{\varphi}_{\tau, s}=\boldsymbol{T}_{\varphi_{t, s}(\mathrm{x})} \boldsymbol{\varphi}_{\tau, t} \circ T_{\mathrm{x}} \boldsymbol{\varphi}_{t, s}
$$

implies the corresponding one between material deformation gradients:

$$
\mathbf{F}_{\tau, s}=\mathbf{F}_{\tau, t} \circ \mathbf{F}_{t, s}
$$

Time rate of deformation gradient, Truesdell \& Noll (1965)

$$
\dot{\mathbf{F}}_{t, s}=\mathbf{L}_{t} \mathbf{F}_{t, s}
$$

with $\dot{\mathbf{F}}_{t, s}:=\partial_{\tau=t} \mathbf{F}_{\tau, s}$ and $\mathbf{L}_{t}:=\partial_{\tau=t} \mathbf{F}_{\tau, t}$ time derivatives.

$$
\mathbf{L}_{t}(\mathbf{x}) \cdot \mathbf{h}_{\mathbf{x}}:=\partial_{\tau=t} \mathbf{F}_{\tau, t}(\mathbf{x}) \cdot \mathbf{h}_{\mathbf{x}} \in \mathbb{T}_{\mathbf{x}} \boldsymbol{\Omega}_{t}, \quad \forall \mathbf{h}_{\mathbf{x}} \in \mathbb{T}_{\mathbf{x}} \boldsymbol{\Omega}_{t}
$$

with $\mathbf{F}_{\tau, t}(\mathbf{x}) \cdot \mathbf{h}_{\mathbf{x}} \in \mathbb{T}_{\mathbf{x}} \boldsymbol{\Omega}_{\tau}$. The LIE time derivative gives:

$$
\partial_{\tau=t}\left(T_{\mathrm{x}} \boldsymbol{\varphi}_{\tau, t}\right)^{-1} \cdot\left(T_{\mathrm{x}} \boldsymbol{\varphi}_{\tau, t} \cdot \mathbf{h}_{\mathrm{x}}\right)=\partial_{\tau=t} \mathbf{h}_{\mathrm{x}}=0
$$

## Change of observer

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A basic question
Rasir
Tangent spaces
Tangent functor
Fiber bundles
Trivial and non-trivial fiber bundles

Sections
Tensor bundle and sections

Push and pull
Push and pull of tensor fields

Parallel transport
Derivatives
Key contributions

## Kinematics

Metric measurements
Metric theory
Events manifold fibrations

## Change of observer

- Change of observer $\zeta_{\mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E} ; \mathrm{E})$, time-bundle automorphism


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- Relative motion $\quad \boldsymbol{\zeta} \in \mathrm{C}^{1}\left(\mathcal{T}_{\varphi} ; \mathcal{T}_{\boldsymbol{\zeta} \uparrow \varphi}\right)$, time-bundle diffeomorphism

- Pushed motion

$$
\begin{aligned}
& \zeta_{t}\left(\Omega_{t}\right) \xrightarrow{(\zeta \uparrow \varphi)_{\tau, t}} \zeta_{\tau}\left(\Omega_{\tau}\right) \\
& \begin{array}{rc}
\zeta_{t} \uparrow & \zeta_{\tau}^{\uparrow} \\
\Omega_{t} \xrightarrow{\varphi_{\tau, t}} & \Longleftrightarrow \Omega_{\tau}
\end{array}
\end{aligned}
$$

## Consequences of the Covariance Paradigm

Time Invariance and Frame Invariance of material fields

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## Consequences of the Covariance Paradigm

Time Invariance and Frame Invariance of material fields

- Time Invariance $\quad \mathbf{s}_{\mathcal{T}_{\boldsymbol{\varphi}}, \tau}=\boldsymbol{\varphi}_{\tau, t} \uparrow \mathbf{s}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}$


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## Consequences of the Covariance Paradigm

Time Invariance and Frame Invariance of material fields

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- Frame Invariance

$$
\mathbf{s}_{\mathcal{T}_{\zeta \uparrow \varphi}}=\zeta \uparrow \mathbf{s}_{\mathcal{T}_{\varphi}}
$$

with: $\boldsymbol{\zeta} \in \mathrm{C}^{1}\left(\mathcal{T}_{\varphi} ; \mathcal{T}_{\boldsymbol{\zeta} \uparrow \varphi}\right)$ relative motion

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## Properties of Lie derivative

- Push of Lie time derivative to a fixed configuration

$$
\boldsymbol{\varphi}_{t, \mathrm{FIX}} \downarrow\left(\mathcal{L}_{\mathbf{v}_{\tau_{\varphi}}} \mathbf{s}_{\mathcal{T}_{\varphi}}\right)_{t}=\partial_{\tau=t} \boldsymbol{\varphi}_{\tau, \mathrm{FIX}} \downarrow \mathbf{s}_{\mathcal{T}_{\boldsymbol{\varphi}}, \tau}
$$

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$$

- Lie time derivative along pushed motions

$$
\mathcal{L}_{\mathbf{v}_{\mathcal{T}_{\boldsymbol{\zeta}} \boldsymbol{\varphi}}}\left(\boldsymbol{\zeta} \uparrow \mathbf{s}_{\varphi}\right)=\boldsymbol{\zeta} \uparrow\left(\mathcal{L}_{\mathbf{v}_{\tau_{\varphi}}} \mathbf{s}_{\mathcal{T}_{\varphi}}\right)
$$

## Constitutive laws

## Constitutive laws

- Constitutive operator $\mathbf{H}_{\mathcal{T}_{\varphi}}$

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A basic question

## Constitutive laws

- Constitutive operator $\mathbf{H}_{\mathcal{T}_{\varphi}}$

A material bundle morphism whose domain and codomain are Whitney products of material tensor bundles

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A basic question

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A material bundle morphism whose domain and codomain are Whitney products of material tensor bundles

- Constitutive time invariance

$$
\begin{gathered}
\mathbf{H}_{\mathcal{T}_{\varphi}, \tau}=\boldsymbol{\varphi}_{\tau, t} \uparrow \mathbf{H}_{\mathcal{T}_{\varphi}, t} \\
\left(\boldsymbol{\varphi}_{\tau, t} \uparrow \mathbf{H}_{\mathcal{T}_{\varphi}, t}\right)\left(\boldsymbol{\varphi}_{\tau, t} \uparrow \mathbf{s}_{\mathcal{T}_{\varphi}, t}\right)=\boldsymbol{\varphi}_{\tau, t} \uparrow\left(\mathbf{H}_{\mathcal{T}_{\varphi}, t}\left(\mathbf{s}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}\right)\right)
\end{gathered}
$$

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\left(\boldsymbol{\varphi}_{\tau, t} \uparrow \mathbf{H}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}\right)\left(\boldsymbol{\varphi}_{\tau, t} \uparrow \mathbf{s}_{\tau_{\varphi}, t}\right)=\boldsymbol{\varphi}_{\tau, t} \uparrow\left(\mathbf{H}_{\tau_{\varphi}, t}\left(\mathbf{s}_{\tau_{\varphi}, t}\right)\right)
\end{gathered}
$$

- Constitutive invariance under relative motions

$$
\begin{gathered}
\mathbf{H}_{\tau_{\zeta \uparrow \varphi}}=\boldsymbol{\zeta} \uparrow \mathbf{H}_{\tau_{\varphi}} \\
\left(\boldsymbol{\zeta} \uparrow \mathbf{H}_{\tau_{\varphi}}\right)\left(\boldsymbol{\zeta} \uparrow \mathbf{s}_{\tau_{\varphi}}\right)=\boldsymbol{\zeta} \uparrow\left(\mathbf{H}_{\tau_{\varphi}}\left(\mathbf{s}_{\tau_{\varphi}}\right)\right)
\end{gathered}
$$

## Hypo-elasticity

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Key contributions

## Hypo-elasticity

- Constitutive hypo-elastic law $\mathbf{e l}_{\mathcal{T}_{\varphi}}$ elastic stretching

$$
\left\{\begin{aligned}
\dot{\varepsilon}_{\mathcal{T}_{\varphi}} & =\mathbf{e l}_{\mathcal{T}_{\varphi}} \\
\mathbf{e l}_{\mathcal{T}_{\varphi}} & =\mathbf{H}_{\mathcal{T}_{\varphi}}^{\mathrm{HYPO}}
\end{aligned} \boldsymbol{\sigma}_{\mathcal{T}_{\varphi}}\right) \cdot \dot{\boldsymbol{\sigma}}_{\mathcal{T}_{\varphi}}
$$

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\end{aligned}\right.
$$

- Cauchy integrability

$$
\begin{aligned}
&\left\langle d_{F} \mathbf{H}_{\mathcal{T}_{\varphi}}^{\mathrm{HYPO}}\left(\sigma_{\mathcal{T}_{\varphi}}\right) \cdot\right.\left.\delta \boldsymbol{\sigma}_{\mathcal{T}_{\varphi}} \cdot \delta_{1} \sigma_{\mathcal{T}_{\varphi}}, \delta_{2} \sigma_{\mathcal{T}_{\varphi}}\right\rangle=\text { symmetric } \\
& \Longrightarrow \quad \mathbf{H}_{\mathcal{T}_{\varphi}}^{\mathrm{HYPO}}\left(\sigma_{\mathcal{T}_{\varphi}}\right)=d_{F} \boldsymbol{\Phi}_{\mathcal{T}_{\varphi}}\left(\sigma_{\mathcal{T}_{\varphi}}\right)
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$$

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\end{aligned}\left(\boldsymbol{\sigma}_{\mathcal{T}_{\varphi}}\right) \cdot \dot{\boldsymbol{\sigma}}_{\mathcal{T}_{\varphi}}\right.
$$

- Cauchy integrability

$$
\left\langle d_{F} \mathbf{H}_{\mathcal{T}_{\varphi}}^{\mathrm{HYPO}}\left(\boldsymbol{\sigma}_{\mathcal{T}_{\varphi}}\right) \cdot \delta \boldsymbol{\sigma}_{\mathcal{T}_{\varphi}} \cdot \delta_{1} \boldsymbol{\sigma}_{\mathcal{T}_{\varphi}}, \delta_{2} \boldsymbol{\sigma}_{\mathcal{T}_{\varphi}}\right\rangle=\text { symmetric }
$$

$$
\Longrightarrow \quad \mathbf{H}_{\mathcal{T}_{\varphi}}^{\mathrm{HYPO}}\left(\sigma_{\mathcal{T}_{\varphi}}\right)=d_{F} \boldsymbol{\Phi}_{\mathcal{T}_{\varphi}}\left(\sigma_{\mathcal{T}_{\varphi}}\right)
$$

- Green integrability

$$
\begin{aligned}
&\left\langle\mathbf{H}_{\mathcal{T}_{\varphi}}^{\mathrm{HYPO}}\left(\boldsymbol{\sigma}_{\mathcal{T}_{\varphi}}\right) \cdot \delta_{1} \boldsymbol{\sigma}_{\mathcal{T}_{\varphi}}, \delta_{2} \sigma_{\mathcal{T}_{\varphi}}\right\rangle=\text { symmetric } \\
& \Longrightarrow \quad \Phi_{\mathcal{T}_{\varphi}}\left(\sigma_{\mathcal{T}_{\varphi}}\right)=d_{F} E_{\mathcal{T}_{\varphi}}^{*}\left(\sigma_{\mathcal{T}_{\varphi}}\right)
\end{aligned}
$$

## Elasticity

- Elastic constitutive operator: hypo-elastic constitutive operator which is integrable and time invariant


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$$
\left\{\begin{aligned}
\dot{\varepsilon}_{\mathcal{T}_{\varphi}} & =\mathbf{e} \mathbf{I}_{\mathcal{T}_{\varphi}} \\
\mathbf{e l}_{\mathcal{T}_{\varphi}} & =d_{F}^{2} E_{\mathcal{T}_{\varphi}}^{*}\left(\boldsymbol{\sigma}_{\mathcal{T}_{\varphi}}\right) \cdot \dot{\boldsymbol{\sigma}}_{\mathcal{T}_{\varphi}}
\end{aligned}\right.
$$

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\end{aligned}\right.
$$

- pull-back to reference:

$$
\begin{aligned}
\boldsymbol{\varphi}_{t, \mathrm{FIX}} \downarrow \mathbf{e l}_{\mathcal{T}_{\boldsymbol{\varphi}}, t} & =d_{F}^{2} E_{\mathrm{FIX}}^{*}\left(\boldsymbol{\varphi}_{t, \mathrm{FIX}} \downarrow \boldsymbol{\sigma}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}\right) \cdot \partial_{\tau=t} \boldsymbol{\varphi}_{\tau, \mathrm{FIX}} \downarrow \boldsymbol{\sigma}_{\boldsymbol{\varphi}, \tau} \\
& =\partial_{\tau=t} d_{F} E_{\mathrm{FIX}}^{*}\left(\boldsymbol{\varphi}_{\tau, \mathrm{FIX}} \downarrow \boldsymbol{\sigma}_{\boldsymbol{\varphi}, \tau}\right)
\end{aligned}
$$

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\mathbf{e l}_{\mathcal{T}_{\varphi}} & =d_{F}^{2} E_{\mathcal{T}_{\varphi}}^{*}\left(\boldsymbol{\sigma}_{\mathcal{T}_{\varphi}}\right) \cdot \dot{\boldsymbol{\sigma}}_{\mathcal{T}_{\varphi}}
\end{aligned}\right.
$$

- pull-back to reference:

$$
\begin{aligned}
\boldsymbol{\varphi}_{t, \mathrm{FIX}} \downarrow \mathbf{e l}_{\mathcal{T}_{\boldsymbol{\varphi}}, t} & =d_{F}^{2}{E_{\mathrm{FIX}}^{*}}^{*}\left(\boldsymbol{\varphi}_{t, \mathrm{FIX}} \downarrow \boldsymbol{\sigma}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}\right) \cdot \partial_{\tau=t} \boldsymbol{\varphi}_{\tau, \mathrm{FIX}} \downarrow \boldsymbol{\sigma}_{\boldsymbol{\varphi}, \tau} \\
& =\partial_{\tau=t} d_{F} E_{\mathrm{FIX}}^{*}\left(\boldsymbol{\varphi}_{\tau, \mathrm{FIX}} \downarrow \boldsymbol{\sigma}_{\boldsymbol{\varphi}, \tau}\right) \\
\boldsymbol{\varphi}_{\tau, \mathrm{FIX}} & :=\boldsymbol{\varphi}_{\tau, t} \circ \boldsymbol{\varphi}_{t, \mathrm{FIX}} \\
E_{\mathrm{FIX}}^{*} & :=\boldsymbol{\varphi}_{t, \mathrm{FIX}} \downarrow E_{\mathcal{T}_{\boldsymbol{\varphi}}, t}^{*} \quad \text { time invariant }
\end{aligned}
$$

## Conservativeness of hyper-elasticity

## Conservativeness of hyper-elasticity

Green integrability of the elastic operator $\mathbf{H}_{\mathcal{T}_{\varphi}}$ as a function of the Kirchнoff stress tensor field implies conservativeness:

$$
\oint_{I} \int_{\Omega_{t}}\left\langle\boldsymbol{\sigma}_{\mathcal{T}_{\varphi}, t}, \mathbf{e l}_{\mathcal{T}_{\varphi}, t}\right\rangle \mathbf{m}_{\mathcal{T}_{\varphi}, t} d t=0
$$

for any cycle in the stress time-bundle,
i.e. for any stress path $\boldsymbol{\sigma}_{\mathcal{T}_{\varphi}} \in \mathrm{C}^{1}\left(I ; \operatorname{CoN}\left(\mathbb{V} \mathcal{T}_{\varphi}\right)\right)$
such that:

$$
\boldsymbol{\sigma}_{\mathcal{T}_{\varphi}, t_{2}}=\boldsymbol{\varphi}_{t_{2}, t_{1}} \uparrow \boldsymbol{\sigma}_{\mathcal{T}_{\varphi}, t_{1}}, \quad I=\left[t_{1}, t_{2}\right]
$$

## Elasto-visco-plasticity

## Elasto-visco-plasticity

- Constitutive law
$\mathbf{e l}_{\mathcal{T}_{\varphi}}$ elastic stretching
$\mathbf{p l}_{\mathcal{T}_{\varphi}}$ visco-plastic stretching

$$
\left\{\begin{aligned}
\dot{\varepsilon}_{\mathcal{T}_{\varphi}} & =\mathbf{e l}_{\mathcal{T}_{\varphi}}+\mathbf{p} \mathbf{I}_{\mathcal{T}_{\varphi}} \\
\mathbf{e l}_{\mathcal{T}_{\varphi}} & =d_{F}^{2} E_{\mathcal{T}_{\varphi}}^{*}\left(\sigma_{\mathcal{T}_{\varphi}}\right) \cdot \dot{\sigma}_{\mathcal{T}_{\varphi}} \\
\mathbf{p l}_{\mathcal{T}_{\varphi}} & \in \partial_{F} \mathcal{F}_{\mathcal{T}_{\varphi}}\left(\boldsymbol{\sigma}_{\varphi}\right)
\end{aligned}\right.
$$

stretching additivity
hyper-elastic law
visco-plastic flow rule

## Reference strains

## Reference strains

- total strain in the time interval $I=[s, t]$ :

$$
\varepsilon_{\mathcal{T}_{\varphi}, t, s}:=\varphi_{t, s} \downarrow \mathbf{g}_{\mathcal{T}_{\varphi}, t}-\mathbf{g}_{\mathcal{T}_{\varphi}, s}
$$

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$$
\varepsilon_{\mathcal{T}_{\varphi}, t, s}:=\varphi_{t, s} \downarrow \mathbf{g}_{\mathcal{T}_{\varphi}, t}-\mathbf{g}_{\mathcal{T}_{\varphi}, s}
$$

- reference total strain:

$$
\begin{aligned}
\varepsilon_{\mathcal{T}_{\boldsymbol{\varphi}}, l}^{\mathrm{FIX}} & :=\frac{1}{2} \int_{I} \partial_{\tau=t} \boldsymbol{\varphi}_{\tau, \mathrm{FIX}} \downarrow \mathbf{g}_{\mathcal{T}_{\boldsymbol{\varphi}}, \tau} d t \\
& =\frac{1}{2} \boldsymbol{\varphi}_{t, \mathrm{FIX}} \downarrow \mathbf{g}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}-\frac{1}{2} \boldsymbol{\varphi}_{s, \mathrm{FIX}} \downarrow \mathbf{g}_{\mathcal{T}_{\boldsymbol{\varphi}}, s} \\
& =\frac{1}{2} \boldsymbol{\varphi}_{s, \mathrm{FIX}} \downarrow\left(\boldsymbol{\varphi}_{t, s} \downarrow \mathbf{g}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}-\mathbf{g}_{\mathcal{T}_{\boldsymbol{\varphi}}, s}\right)=\frac{1}{2} \boldsymbol{\varphi}_{s, \mathrm{FIX}} \downarrow \varepsilon_{\mathcal{T}_{\boldsymbol{\varphi}}, t, s}
\end{aligned}
$$

## Reference strains

- total strain in the time interval $I=[s, t]$ :

$$
\varepsilon_{\mathcal{T}_{\varphi}, t, s}:=\varphi_{t, s} \downarrow \mathbf{g}_{\mathcal{T}_{\varphi}, t}-\mathbf{g}_{\mathcal{T}_{\varphi}, s}
$$

- reference total strain:

$$
\begin{aligned}
\varepsilon_{\mathcal{T}_{\varphi}, l}^{\mathrm{FIX}} & :=\frac{1}{2} \int_{I} \partial_{\tau=t} \varphi_{\tau, \mathrm{FIX}} \downarrow \mathbf{g}_{\mathcal{T}_{\varphi}, \tau} d t \\
& =\frac{1}{2} \boldsymbol{\varphi}_{t, \mathrm{FIX}} \downarrow \mathbf{g}_{\tau_{\varphi}, t}-\frac{1}{2} \boldsymbol{\varphi}_{s, \mathrm{FIX}} \downarrow \mathbf{g}_{\tau_{\varphi}, s} \\
& =\frac{1}{2} \boldsymbol{\varphi}_{s, \mathrm{FIX}} \downarrow\left(\boldsymbol{\varphi}_{t, s} \downarrow \mathbf{g}_{\tau_{\varphi}, t}-\mathbf{g}_{\tau_{\varphi}, s}\right)=\frac{1}{2} \boldsymbol{\varphi}_{s, \mathrm{FIX}} \downarrow \varepsilon_{\mathcal{T}_{\varphi}, t, s}
\end{aligned}
$$

- reference elastic and visco-plastic strain:

$$
\mathbf{e l}_{\mathcal{T}_{\varphi}, l}^{\mathrm{FIX}}:=\int_{I} \boldsymbol{\varphi}_{t, \mathrm{FIX}} \backslash \mathbf{e}_{\tau_{\varphi}, t} d t, \quad \mathbf{p}_{\mathcal{T}_{\varphi}, l}^{\mathrm{FIX}}:=\int_{I} \boldsymbol{\varphi}_{t, \mathrm{FIX}} \downarrow \mathbf{p}_{\mathcal{T}_{\varphi}, t} d t
$$

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$$

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$$
\begin{aligned}
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& =\frac{1}{2} \boldsymbol{\varphi}_{t, \mathrm{FIX}} \downarrow \mathbf{g}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}-\frac{1}{2} \varphi_{s, \mathrm{FIX}} \downarrow \mathbf{g}_{\mathcal{T}_{\boldsymbol{\varphi}}, s} \\
& =\frac{1}{2} \boldsymbol{\varphi}_{s, \mathrm{FIX}} \downarrow\left(\boldsymbol{\varphi}_{t, s} \downarrow \mathbf{g}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}-\mathbf{g}_{\mathcal{T}_{\boldsymbol{\varphi}}, s}\right)=\frac{1}{2} \boldsymbol{\varphi}_{s, \mathrm{FIX}} \downarrow \varepsilon_{\mathcal{T}_{\boldsymbol{\varphi}}, t, s}
\end{aligned}
$$

- reference elastic and visco-plastic strain:

$$
\mathbf{e l}_{\mathcal{T}_{\varphi}, l}^{\mathrm{FIX}}:=\int_{I} \varphi_{t, \mathrm{FIX}} \backslash \mathbf{e}_{\mathcal{T}_{\varphi}, t} d t, \quad \mathbf{p}_{\mathcal{T}_{\varphi}, l}^{\mathrm{FIX}}:=\int_{I} \boldsymbol{\varphi}_{t, \mathrm{FIX}} \downarrow \mathbf{p}_{\mathcal{T}_{\varphi}, t} d t
$$

- additivity of reference strains:

$$
\varepsilon_{\mathcal{T}_{\boldsymbol{\varphi}}, I}^{\mathrm{FIX}}=\mathbf{e l}_{\mathcal{T}_{\boldsymbol{\varphi}}, I}^{\mathrm{FIX}}+\mathbf{p} \mathbf{I}_{\mathcal{T}_{\boldsymbol{\varphi}}, l}^{\mathrm{FIX}}
$$

## Material Frame Indifference (MFI)

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Ansatz

Draleriomena

## A basic question

Basic
Tansent spaces
Tangent functor
ciber bundloe
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Push and pull of tensor fields

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## Derivatives

Key contributions

## Material Frame Indifference (MFI)

Ansatz

- Material fields are frame invariant


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## A basic question

## Material Frame Indifference (MFI)

Ansatz

- Material fields are frame invariant


## Principle of MFI

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A basic question

## Material Frame Indifference (MFI)

Ansatz

- Material fields are frame invariant


## Principle of MFI

- Any constitutive law must conform to the principle of MFI which requires that material fields, fulfilling the law, will still fulfill it when evaluated by another Euclid observer


## Material Frame Indifference (MFI)

Ansatz

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## Principle of MFI

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$$
\mathbf{H}_{\mathcal{T}_{\zeta}{ }^{\mathrm{iso} \uparrow \varphi}}\left(\zeta^{\mathrm{iso}} \uparrow \mathbf{s}_{\mathcal{T}_{\varphi}}\right)=\zeta^{\mathrm{iso}} \uparrow \mathbf{H}_{\mathcal{T}_{\varphi}}\left(\mathbf{s}_{\mathcal{T}_{\varphi}}\right)
$$

for any isometric relative motion $\zeta^{\text {iso }} \in \mathrm{C}^{1}\left(\mathcal{T}_{\varphi} ; \mathcal{T}_{\zeta^{\text {iso }} \uparrow \varphi}\right)$ induced by a change of Euclid observer $\zeta_{E}^{\text {iso }} \in \mathrm{C}^{1}(\mathbf{E} ; \mathbf{E})$.

## Material Frame Indifference (MFI)

Ansatz

- Material fields are frame invariant


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$$
\mathbf{H}_{\mathcal{T}_{\zeta^{\mathrm{iso}} \uparrow \varphi}}\left(\zeta^{\mathrm{iso}} \uparrow \mathbf{s}_{\mathcal{T}_{\varphi}}\right)=\zeta^{\mathrm{iso}} \uparrow \mathbf{H}_{\mathcal{T}_{\varphi}}\left(\mathbf{s}_{\mathcal{T}_{\varphi}}\right)
$$

for any isometric relative motion $\zeta^{\text {iso }} \in \mathrm{C}^{1}\left(\mathcal{T}_{\varphi} ; \mathcal{T}_{\zeta^{\text {iso }} \uparrow \varphi}\right)$ induced by a change of Euclid observer $\zeta_{E}^{\text {iso }} \in \mathrm{C}^{1}(\mathbf{E} ; \mathbf{E})$.

## Equivalent condition

- Constitutive operators must be frame invariant


## MFI in elasto-visco-plasticity

## MFI in elasto-visco-plasticity

- Frame invariance of the hypo-elastic operator

$$
\mathbf{H}_{\mathcal{T}_{\boldsymbol{\zeta}^{\mathrm{ISO}} \uparrow \boldsymbol{\varphi}}^{\mathrm{HYPO}}}=\zeta^{\mathrm{ISO}} \uparrow \mathbf{H}_{\mathcal{T}_{\boldsymbol{\varphi}}}^{\mathrm{HYPO}}
$$

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A basic question

## MFI in elasto-visco-plasticity

- Frame invariance of the hypo-elastic operator

$$
\mathbf{H}_{\mathcal{T}_{\zeta^{\mathrm{SOO}} \uparrow \varphi}^{\mathrm{HYPO}}}=\zeta^{\mathrm{ISO}} \uparrow \mathbf{H}_{\mathcal{T}_{\varphi}}^{\mathrm{HYPO}}
$$

Pushed operator

$$
\left(\boldsymbol{\zeta}^{\mathrm{ISO}} \uparrow \mathbf{H}_{\mathcal{T}_{\varphi}}^{\mathrm{HYPO}}\right)\left(\boldsymbol{\zeta}^{\mathrm{ISO}} \uparrow \boldsymbol{\sigma}_{\mathcal{T}_{\varphi}}\right) \cdot \boldsymbol{\zeta}^{\mathrm{ISO}} \uparrow \dot{\boldsymbol{\sigma}}_{\mathcal{T}_{\varphi}}=\boldsymbol{\zeta}^{\mathrm{ISO}} \uparrow\left(\mathbf{H}_{\mathcal{T}_{\varphi}}^{\mathrm{HYPO}}\left(\boldsymbol{\sigma}_{\mathcal{T}_{\varphi}}\right) \cdot \dot{\boldsymbol{\sigma}}_{\mathcal{T}_{\varphi}}\right)
$$

## MFI in elasto-visco-plasticity

- Frame invariance of the hypo-elastic operator

$$
\mathbf{H}_{\mathcal{T}_{\zeta^{\mathrm{ISO}} \uparrow \varphi}^{\mathrm{HYPO}}}^{\mathrm{YP}}=\boldsymbol{\zeta}^{\mathrm{ISO}} \uparrow \mathbf{H}_{\mathcal{T}_{\varphi}}^{\mathrm{HYPO}}
$$

Pushed operator

$$
\left(\boldsymbol{\zeta}^{\mathrm{ISO}} \uparrow \mathbf{H}_{\mathcal{T}_{\boldsymbol{\varphi}}}^{\mathrm{HYPO}}\right)\left(\boldsymbol{\zeta}^{\mathrm{ISO}} \uparrow \boldsymbol{\sigma}_{\mathcal{T}_{\boldsymbol{\varphi}}}\right) \cdot \boldsymbol{\zeta}^{\mathrm{ISO}} \uparrow \dot{\boldsymbol{\sigma}}_{\mathcal{T}_{\boldsymbol{\varphi}}}=\boldsymbol{\zeta}^{\mathrm{ISO}} \uparrow\left(\mathbf{H}_{\mathcal{T}_{\boldsymbol{\varphi}}}^{\mathrm{HYPO}}\left(\boldsymbol{\sigma}_{\mathcal{T}_{\boldsymbol{\varphi}}}\right) \cdot \dot{\boldsymbol{\sigma}}_{\mathcal{T}_{\boldsymbol{\varphi}}}\right)
$$

Examples:

- the simplest hypo-elastic operator is Green integrable and frame invariant:

$$
\mathbf{H}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}^{\mathrm{HYP}}\left(\mathbf{T}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}\right):=\frac{1}{2 \mu} \mathbb{I}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}-\frac{\nu}{E} \mathbf{I}_{\mathcal{T}_{\boldsymbol{\varphi}}, t} \otimes \mathbf{I}_{\mathcal{T}_{\boldsymbol{\varphi}}, t}
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These results provide answers to unsolved questions posed in:
J.C. Simó \& K.S. Pister, Remarks on rate constitutive equations for finite deformation problems: computational implications, Comp. Meth. Appl. Mech. Eng. 46 (1984) 201-215.
J. C. Simó \& M. Ortiz, A unified approach to finite deformation elastoplastic analysis based on the use of hyperelastic constitutive equations, Comp. Meth. Appl. Mech. Eng. 49 (1985) 221-245.

## Achievements

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NLCM
Dralenamena
A basic question
Racie
Tangent spaces
Tangent functor
Fiben bundles
Trivial and non-trivial fiber bundles

Sections
Tensor bundle and sections

Push and pull
Push and pull of tensor fields

Parallel transport

## Derivatives

Key contributions

## Kinematics

Metric measurements
Metric theory
Events manifold fibrations

## Achievements

- Notion of spatial and material fields

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Drelemame
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- Notions and treatments of constitutive models in the nonlinear range should be revised and reformulated


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- Notions and treatments of constitutive models in the nonlinear range should be revised and reformulated
- Algorithms for numerical computations must be modified to comply with the covariant theory; multiplicative decomposition of the deformation gradient should be deemed as geometrically inconsistent


[^0]:    ${ }^{2}$ See e.g.

    1) C. Truesdell, A first Course in Rational Continuum Mechanics Second Ed. Academic Press, New-York (1991). First Ed. 1977
    2) M.E. Gurtin, An Introduction to Continuum Mechanics Academic Press, San Diego (1981)
    3) J.E. Marsden \& T.J.R. Hughes, Mathematical Foundations of Elasticity Prentice-Hall, Redwood City, Cal. (1983)
