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28 June 2011



NLCM = Non-Linear Continuum Mechanicsand DG = Differential Geometry

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Covariance Paradigm

NLCM = Non-Linear Continuum Mechanicsand DG = Differential Geometry

NLCM is an important source of inspiration for DG and DG is the natural tool to develop a mathematical modeling of NLCM



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Hermann Weyl (1885-1955)

In these days the angel of topology and the devil of abstract algebra fight for the soul of each individual mathematical domain.

H. Weyl, "Invariants", Duke Mathematical Journal 5 (3): (1939) 489-502

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A basic question in NLCM

How to compare the metric and stress tensors at corresponding points in displaced placements of a body? The G-Factor Impact in NLCM

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Covariance Paradigm

A basic question in NLCM

- How to compare the metric and stress tensors at corresponding points in displaced placements of a body?
- Devil's temptation:

In 3D bodies it might seem as natural to compare by translation the traction vectors corresponding to translated normals to cutting surfaces. This is tacitly done when writing the stress time-rate as \dot{T} but is geometrically untenable as may be more clearly seen by considering 1D and 2D models (wires and membranes).

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A basic question in NLCM

- How to compare the metric and stress tensors at corresponding points in displaced placements of a body?
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In 3D bodies it might seem as natural to compare by translation the traction vectors corresponding to translated normals to cutting surfaces. This is tacitly done when writing the stress time-rate as \dot{T} but is geometrically untenable as may be more clearly seen by considering 1D and 2D models (wires and membranes).

Hint:

Tangent vectors to a body placement may be transformed into tangent vectors to another body placement only by means of the differential of the displacement map. This is the essence of the COVARIANCE PARADIGM. The G-Factor Impact in NLCM

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DIMENSIONALITY INDEPENDENCE: A geometrically consistent framework should be equally applicable to body models of any dimension.

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DIMENSIONALITY INDEPENDENCE: A geometrically consistent framework should be equally applicable to body models of any dimension.

Motivation for the COVARIANCE PARADIGM¹

¹G. Romano, R. Barretta, 2011. Covariant hypo-elasticity. Eur. J. Mech. A-Solids. DOI: 10.1016/j.euromechsol.2011.05.005 The G-Factor Impact in NLCM

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Tangent vector to a manifold: velocity of a curve

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Tangent vector to a manifold: velocity of a curve $\mathbf{c} \in C^1([a, b]; \mathbb{M}), \quad \lambda \in [a, b], \quad \mathbf{x} = \mathbf{c}(\lambda) \quad \text{base point}$ $\mathbf{v} := \partial_{\mu=\lambda} \mathbf{c}(\mu) \in \mathbb{T}_{\mathbf{x}}\mathbb{M}$ The G-Factor Impact in NLCM

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 $\mathbf{v}^* \in L\left(\mathbb{T}_{\mathbf{x}}\mathbb{M}\,;\,\mathcal{R}
ight) \in \mathbb{T}_{\mathbf{v}}^*\mathbb{M}$

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 $\mathbf{v}^* \in L\left(\mathbb{T}_{\mathbf{x}}\mathbb{M}\,;\mathcal{R}
ight) \in \mathbb{T}_{\mathbf{x}}^*\mathbb{M}$

Tangent map

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Cotangent vector

$$\mathbf{v}^* \in L\left(\mathbb{T}_{\mathbf{x}}\mathbb{M}\,;\,\mathcal{R}
ight)\in\mathbb{T}_{\mathbf{x}}^*\mathbb{M}$$

Tangent map

► A map $\boldsymbol{\zeta} \in \mathrm{C}^1(\mathbb{M}\,;\mathbb{N})$ sends a curve $\mathbf{c} \in \mathrm{C}^1([a,b]\,;\mathbb{M})$ into a curve $\boldsymbol{\zeta} \circ \mathbf{c} \in \mathrm{C}^1([a,b]\,;\mathbb{N})$. The G-Factor Impact in NLCM

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Cotangent vector

$$\mathbf{v}^* \in L\left(\mathbb{T}_{\mathbf{x}}\mathbb{M}\,;\mathcal{R}
ight) \in \mathbb{T}^*_{\mathbf{x}}\mathbb{M}$$

Tangent map

► A map $\boldsymbol{\zeta} \in \mathrm{C}^1(\mathbb{M}\,;\mathbb{N})$ sends a curve $\mathbf{c} \in \mathrm{C}^1([a,b]\,;\mathbb{M})$ into a curve $\boldsymbol{\zeta} \circ \mathbf{c} \in \mathrm{C}^1([a,b]\,;\mathbb{N})$.

► The tangent map $T_{\mathbf{x}}\boldsymbol{\zeta} \in C^{0}(\mathbb{T}_{\mathbf{x}}\mathbb{M}\,;\mathbb{T}_{\boldsymbol{\zeta}(\mathbf{x})}\mathbb{N})$ sends a tangent vector at $\mathbf{x} \in \mathbb{M}$ $\mathbf{v} \in \mathbb{T}_{\mathbf{x}}(\mathbb{M}) := \partial_{\mu=\lambda} \mathbf{c}(\mu)$ into a tangent vector at $\boldsymbol{\zeta}(\mathbf{x}) \in \mathbb{N}$ $T_{\mathbf{x}}\boldsymbol{\zeta} \cdot \mathbf{v} \in \mathbb{T}_{\boldsymbol{\zeta}(\mathbf{x})}(\mathbb{N}) := \partial_{\mu=\lambda} (\boldsymbol{\zeta} \circ \mathbf{c})(\mu)$ The G-Factor Impact in NLCM

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disjoint union of tangent spaces:

 $\mathbb{TM}:=\cup_{\textbf{x}\in\mathbb{M}}\mathbb{T}_{\textbf{x}}\mathbb{M}$





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Tangent bundle

disjoint union of tangent spaces:

$$\mathbb{TM} := \cup_{\mathbf{x} \in \mathbb{M}} \mathbb{T}_{\mathbf{x}} \mathbb{M}$$

• Projection:
$$oldsymbol{ au}_{\mathbb{M}}\in\mathrm{C}^{1}(\mathbb{TM}\,;\mathbb{M})$$

$$\mathbf{v} \in \mathbb{T}_{\mathbf{x}}\mathbb{M}, \quad \boldsymbol{ au}_{\mathbb{M}}(\mathbf{v}) := \mathbf{x} \quad ext{base point}$$





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disjoint union of tangent spaces:

 $\mathbb{TM}:=\cup_{\textbf{x}\in\mathbb{M}}\mathbb{T}_{\textbf{x}}\mathbb{M}$

• Projection:
$${m au}_{\mathbb M}\in {
m C}^1({\mathbb T}{\mathbb M}\,;{\mathbb M})$$

 $\mathbf{v} \in \mathbb{T}_{\mathbf{x}}\mathbb{M}\,, \quad \boldsymbol{ au}_{\mathbb{M}}(\mathbf{v}) := \mathbf{x} \quad ext{base point}$

Surjective submersion:

 ${\mathcal T}_{{\tt v}} {\bm \tau}_{\mathbb M} \in {\rm C}^1({\mathbb T}_{{\tt v}} {\mathbb T} {\mathbb M}\, ; {\mathbb T}_{{\tt x}} {\mathbb M})$ is surjective





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disjoint union of tangent spaces:

$$\mathbb{TM}:=\cup_{\textbf{x}\in\mathbb{M}}\mathbb{T}_{\textbf{x}}\mathbb{M}$$

• Projection:
$$au_{\mathbb{M}} \in C^1(\mathbb{TM}; \mathbb{M})$$

 $\mathbf{v} \in \mathbb{T}_{\mathbf{x}}\mathbb{M}, \quad \boldsymbol{\tau}_{\mathbb{M}}(\mathbf{v}) := \mathbf{x} \quad \text{base point}$

Surjective submersion:

 ${\mathcal T}_{{\sf v}} {\boldsymbol \tau}_{\mathbb M} \in {\rm C}^1({\mathbb T}_{{\sf v}} {\mathbb T} {\mathbb M}\, ; {\mathbb T}_{{\sf x}} {\mathbb M})$ is surjective

Tangent functor

 $oldsymbol{\zeta} \in \mathrm{C}^1(\mathbb{M}\,;\mathbb{N}) \quad \mapsto \quad Toldsymbol{\zeta} \in \mathrm{C}^0(\mathbb{TM}\,;\mathbb{TN})$



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fiber bundle

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base manifold

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▶ E, M manifolds

base martfold

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Time derivatives

- ► E, M manifolds
- Fiber bundle projection: $\pi^{E}_{\mathbb{M}} \in C^{1}(E; \mathbb{M})$ surjective submersion

base manifold

fiber bundle

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Time derivatives

- ► E, M manifolds
- Fiber bundle projection: $\pi^{E}_{\mathbb{M}} \in C^{1}(E; \mathbb{M})$ surjective submersion
- ► Total space: E
- ► Base space: M
- \blacktriangleright Fiber manifold: $({m \pi}_{\mathbb M}^{
 m E})^{-1}({m x})$ based at ${m x}\in {\mathbb M}$

base manifold

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- ► E, M manifolds
- Fiber bundle projection: $\pi^{E}_{\mathbb{M}} \in C^{1}(E; \mathbb{M})$ surjective submersion
- ► Total space: E
- ► Base space: M
- ▶ Fiber manifold: $(\pi^{\mathrm{E}}_{\mathbb{M}})^{-1}(\mathsf{x})$ based at $\mathsf{x} \in \mathbb{M}$
- ▶ Tangent bundle $T\pi^{\mathrm{E}}_{\mathbb{M}} \in \mathrm{C}^{0}(\mathbb{T}\mathrm{E}\,;\mathbb{T}\mathbb{M})$

base manifold

fiber bundle

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- ► Total space: E
- ► Base space: M
- ▶ Fiber manifold: $(\pi^{\mathrm{E}}_{\mathbb{M}})^{-1}(\mathsf{x})$ based at $\mathsf{x} \in \mathbb{M}$
- ▶ Tangent bundle $T\pi^{\mathrm{E}}_{\mathbb{M}} \in \mathrm{C}^{0}(\mathbb{T}\mathrm{E}\,;\mathbb{T}\mathbb{M})$
- ▶ Vertical tangent subbundle $T\pi_{\mathbb{M}}^{\mathbb{E}} \in \mathrm{C}^{0}(\mathbb{VE}; \mathbb{TM})$

base manifold

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- ► Total space: E
- ► Base space: M
- \blacktriangleright Fiber manifold: $(\pi^{\mathrm{E}}_{\mathbb{M}})^{-1}(\mathsf{x})$ based at $\mathsf{x} \in \mathbb{M}$
- ▶ Tangent bundle $T\pi^{\mathrm{E}}_{\mathbb{M}} \in \mathrm{C}^0(\mathbb{T}\mathrm{E}\,;\mathbb{T}\mathbb{M})$
- ► Vertical tangent subbundle $T\pi_{\mathbb{M}}^{\mathrm{E}} \in \mathrm{C}^{0}(\mathbb{V}\mathrm{E};\mathbb{T}\mathbb{M})$ with: $\delta \mathbf{e} \in \mathbb{V}\mathrm{E} \subset \mathbb{T}\mathrm{E} \implies T_{\mathbf{e}}\pi_{\mathbb{M}}^{\mathrm{E}} \cdot \delta \mathbf{e} = 0$

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Torus



Listing-Möbius strip



Klein Bottle

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▶ Fiber bundle $\pi^{\mathrm{E}}_{\mathbb{M}} \in \mathrm{C}^1(\mathrm{E}\,;\mathbb{M})$



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Math₅

Sections of fiber bundles

 $\boldsymbol{\pi}^{\mathrm{E}}_{\mathbb{M}} \in \mathrm{C}^{1}(\mathrm{E}; \mathbb{M})$ Fiber bundle



 $\mathbf{s}^{\mathrm{E}}_{\mathbb{M}} \in \mathrm{C}^1(\mathbb{M}\,;\mathrm{E})\,, \quad oldsymbol{\pi}^{\mathrm{E}}_{\mathbb{M}} \circ \mathbf{s}^{\mathrm{E}}_{\mathbb{M}} = \mathrm{ID}_{\mathbb{M}}$



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Sections

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Sections of fiber bundles

- ▶ Fiber bundle $\pi^{\mathrm{E}}_{\mathbb{M}} \in \mathrm{C}^{1}(\mathrm{E}; \mathbb{M})$
- ▶ Sections $\mathbf{s}_{\mathbb{M}}^{\mathrm{E}} \in \mathrm{C}^{1}(\mathbb{M}\,;\,\mathrm{E})\,, \quad \boldsymbol{\pi}_{\mathbb{M}}^{\mathrm{E}} \circ \mathbf{s}_{\mathbb{M}}^{\mathrm{E}} = \mathrm{ID}_{\mathbb{M}}$
- ▶ Tangent v.f. $\mathbf{v}_{\rm E} \in \mathbf{0}$
- $\mathbf{v}_{\mathrm{E}} \in \mathrm{C}^1(\mathrm{E}\,;\mathbb{T}\mathrm{E})\,, \quad oldsymbol{ au}_{\mathrm{E}} \circ \mathbf{v}_{\mathrm{E}} = \mathrm{ID}_{\mathrm{E}}$



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Sections of fiber bundles

- ▶ Fiber bundle $\pi^{\mathrm{E}}_{\mathbb{M}} \in \mathrm{C}^1(\mathrm{E}\,;\mathbb{M})$
- ▶ Sections $\mathbf{s}_{\mathbb{M}}^{\mathrm{E}} \in \mathrm{C}^{1}(\mathbb{M}; \mathrm{E}), \quad \boldsymbol{\pi}_{\mathbb{M}}^{\mathrm{E}} \circ \mathbf{s}_{\mathbb{M}}^{\mathrm{E}} = \mathrm{ID}_{\mathbb{M}}$
- ▶ Tangent v.f. $\mathbf{v}_{\mathrm{E}} \in \mathrm{C}^{1}(\mathrm{E}\,; \mathbb{T}\mathrm{E})\,, \quad \boldsymbol{\tau}_{\mathrm{E}} \circ \mathbf{v}_{\mathrm{E}} = \mathrm{ID}_{\mathrm{E}}$
- \blacktriangleright Vertical tangent sections ${\cal T} m{\pi}^{
 m E}_{\mathbb{M}} \circ m{v}_{
 m E} = 0$



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- ▶ Vertical tangent sections $T \boldsymbol{\pi}_{\mathbb{M}}^{\mathrm{E}} \circ \mathbf{v}_{\mathrm{E}} = 0$

Sections of tangent and bi-tangent bundles



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- ▶ Vertical tangent sections $T \boldsymbol{\pi}_{\mathbb{M}}^{\mathrm{E}} \circ \mathbf{v}_{\mathrm{E}} = 0$

Sections of tangent and bi-tangent bundles

Tangent vector fields:

 $\mathbf{v} \in \mathrm{C}^1(\mathbb{M}\,;\mathbb{TM})\,:\, oldsymbol{ au}_\mathbb{M}\circ\mathbf{v} = \mathrm{ID}_\mathbb{M}$



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- Vertical tangent sections $T \boldsymbol{\pi}_{\mathbb{M}}^{\mathrm{E}} \circ \mathbf{v}_{\mathrm{E}} = 0$

Sections of tangent and bi-tangent bundles

Tangent vector fields:

 $\mathbf{v}\in\mathrm{C}^1(\mathbb{M}\,;\mathbb{TM})\,:\,oldsymbol{ au}_\mathbb{M}\circ\mathbf{v}=\mathrm{ID}_\mathbb{M}$

Bi-tangent vector fields:

 $\mathsf{X} \in \mathrm{C}^1(\mathbb{TM}\,;\mathbb{TTM})\,:\, oldsymbol{ au}_{\mathbb{TM}}\circ\mathsf{X}=\mathrm{Id}_{\mathbb{TM}}$



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- ▶ Vertical tangent sections $T \boldsymbol{\pi}_{\mathbb{M}}^{\mathrm{E}} \circ \mathbf{v}_{\mathrm{E}} = 0$

Sections of tangent and bi-tangent bundles

Tangent vector fields:

 $\mathbf{v}\in\mathrm{C}^1(\mathbb{M}\,;\mathbb{TM})\,:\,oldsymbol{ au}_\mathbb{M}\circ\mathbf{v}=\mathrm{ID}_\mathbb{M}$

Bi-tangent vector fields:

 $\mathsf{X} \in \mathrm{C}^1(\mathbb{TM}\,;\mathbb{TTM})\,:\, oldsymbol{ au}_{\mathbb{TM}}\circ\mathsf{X}=\mathrm{Id}_{\mathbb{TM}}$

• Vertical bi-tangent vectors $X \in \operatorname{Ker} \mathcal{T}_{v} \boldsymbol{\tau}_{\mathbb{M}}$



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Covariant

$$\mathbf{s}_{\mathbf{x}}^{\mathrm{COV}} \in \mathrm{COV}_{\mathbf{x}}(\mathbb{TM}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}^2; \mathcal{R}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}; \mathbb{T}_{\mathbf{x}}^*\mathbb{M})$$

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► Covariant

$$\mathbf{s}_{\mathbf{x}}^{\text{COV}} \in \text{COV}_{\mathbf{x}}(\mathbb{TM}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}^2; \mathcal{R}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}; \mathbb{T}_{\mathbf{x}}^*\mathbb{M})$$

Contravariant

$$\mathbf{s}_{\mathbf{x}}^{\mathrm{CON}} \in \mathrm{CON}_{\mathbf{x}}(\mathbb{TM}) = L(\mathbb{T}_{\mathbf{x}}^{*}\mathbb{M}^{2}; \mathcal{R}) = L(\mathbb{T}_{\mathbf{x}}^{*}\mathbb{M}; \mathbb{T}_{\mathbf{x}}\mathbb{M})$$

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► Covariant

$$\mathbf{s}_{\mathbf{x}}^{\text{COV}} \in \text{COV}_{\mathbf{x}}(\mathbb{TM}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}^2; \mathcal{R}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}; \mathbb{T}_{\mathbf{x}}^*\mathbb{M})$$

- ► Contravariant $\mathbf{s}_{\mathbf{x}}^{\text{CON}} \in \text{CON}_{\mathbf{x}}(\mathbb{TM}) = L(\mathbb{T}_{\mathbf{x}}^*\mathbb{M}^2; \mathcal{R}) = L(\mathbb{T}_{\mathbf{x}}^*\mathbb{M}; \mathbb{T}_{\mathbf{x}}\mathbb{M})$
- ► Mixed $\mathbf{s}_{\mathbf{x}}^{\text{MIX}} \in \text{MIX}_{\mathbf{x}}(\mathbb{TM}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}, \mathbb{T}_{\mathbf{x}}^{*}\mathbb{M}; \mathcal{R}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}; \mathbb{T}_{\mathbf{x}}\mathbb{M})$
- with the alteration rules:

$$\mathbf{s}_{\mathbf{x}}^{\mathrm{COV}} = \mathbf{g}_{\mathbf{x}} \circ \mathbf{s}_{\mathbf{x}}^{\mathrm{MIX}} \,, \quad \mathbf{s}_{\mathbf{x}}^{\mathrm{CON}} = \mathbf{s}_{\mathbf{x}}^{\mathrm{MIX}} \circ \mathbf{g}_{\mathbf{x}}^{-1} \,,$$

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- ► Covariant $\mathbf{s}_{\mathbf{x}}^{\text{COV}} \in \text{COV}_{\mathbf{x}}(\mathbb{TM}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}^2; \mathcal{R}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}; \mathbb{T}_{\mathbf{x}}^*\mathbb{M})$
- ► Contravariant $\mathbf{s}_{\mathbf{x}}^{\text{CON}} \in \text{CON}_{\mathbf{x}}(\mathbb{TM}) = L(\mathbb{T}_{\mathbf{x}}^*\mathbb{M}^2; \mathcal{R}) = L(\mathbb{T}_{\mathbf{x}}^*\mathbb{M}; \mathbb{T}_{\mathbf{x}}\mathbb{M})$
- ► Mixed $\mathbf{s}_{\mathbf{x}}^{\text{MIX}} \in \text{MIX}_{\mathbf{x}}(\mathbb{TM}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}, \mathbb{T}_{\mathbf{x}}^{*}\mathbb{M}; \mathcal{R}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}; \mathbb{T}_{\mathbf{x}}\mathbb{M})$
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Tensor bundles and sections

▶ Tensor bundle $au_{\mathbb{M}}^{ ext{TENS}} \in \mathrm{C}^1(ext{TENS}(\mathbb{TM}); \mathbb{M})$

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Tensor spaces

- ► Covariant $\mathbf{s}_{\mathbf{x}}^{\text{COV}} \in \text{COV}_{\mathbf{x}}(\mathbb{TM}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}^2; \mathcal{R}) = L(\mathbb{T}_{\mathbf{x}}\mathbb{M}; \mathbb{T}_{\mathbf{x}}^*\mathbb{M})$
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Tensor bundles and sections

- ▶ Tensor bundle $au_{\mathbb{M}}^{ ext{TENS}} \in \mathrm{C}^1(ext{TENS}(\mathbb{TM}); \mathbb{M})$
- ► Tensor field $\mathbf{s}_{\mathbb{M}}^{\text{TENS}} \in C^1(\mathbb{M}; \text{TENS}(\mathbb{TM}))$

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Tensor spaces

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Tensor bundles and sections

- ▶ Tensor bundle $au_{\mathbb{M}}^{ ext{TENS}} \in \mathrm{C}^1(ext{TENS}(\mathbb{TM}); \mathbb{M})$
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• with:
$$\boldsymbol{ au}_{\mathbb{M}}^{\mathrm{TENS}} \circ \boldsymbol{\mathsf{s}}_{\mathbb{M}}^{\mathrm{TENS}} = \mathrm{ID}_{\mathbb{M}}$$

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Push and pull Given a map $\boldsymbol{\zeta} \in \mathrm{C}^1(\mathbb{M}\,;\mathbb{N})$

Pull-back of a scalar field

$$f: \mathbb{N} \mapsto \mathrm{FUN}(\mathbb{N}) \quad \mapsto \quad \boldsymbol{\zeta} \! \downarrow \! f: \mathbb{M} \mapsto \mathrm{FUN}(\mathbb{M})$$

defined by:

$$(\zeta \downarrow f)_{\mathsf{x}} := \zeta \downarrow f_{\zeta(\mathsf{x})} := f_{\zeta(\mathsf{x})} \in \mathrm{FUN}_{\mathsf{x}}(\mathbb{M})$$

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Covariance Paradigm

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Push-forward of a tangent vector field

 $\mathbf{v}\in\mathrm{C}^1(\mathbb{M}\,;\mathbb{TM})\quad\mapsto\quad \pmb{\zeta}\!\uparrow\!\mathbf{v}:\mathbb{N}\mapsto\mathbb{TN}$

defined by:

$$(\zeta \uparrow \mathbf{v})_{\zeta(\mathbf{x})} := \zeta \uparrow \mathbf{v}_{\mathbf{x}} = T_{\mathbf{x}} \zeta \cdot \mathbf{v}_{\mathbf{x}} \in \mathbb{T}_{\zeta(\mathbf{x})} \mathbb{N}$$
.

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Covectors

$$\langle \zeta {\downarrow} \mathbf{v}^*_{\zeta(\mathbf{x})}, \mathbf{v}_{\mathbf{x}} \rangle = \langle \mathbf{v}^*_{\zeta(\mathbf{x})}, \zeta {\uparrow} \mathbf{v}_{\mathbf{x}} \rangle = \langle T^*_{\zeta(\mathbf{x})} \zeta \circ \mathbf{v}^*_{\zeta(\mathbf{x})}, \mathbf{v}_{\mathbf{x}} \rangle$$

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Covariant tensors

$$\zeta \! \downarrow \! \mathsf{s}^{\mathrm{COV}}_{\zeta(\mathsf{x})} = \mathit{T}^*_{\zeta(\mathsf{x})} \zeta \circ \mathsf{s}^{\mathrm{COV}}_{\zeta(\mathsf{x})} \circ \mathit{T}_{\mathsf{x}} \zeta \in \mathrm{COV}(\mathbb{TM})_{\mathsf{x}}$$

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Covectors

$$\langle \zeta \! \downarrow \! \mathbf{v}^*_{\zeta(\mathbf{x})}, \mathbf{v}_{\mathbf{x}} \rangle = \langle \mathbf{v}^*_{\zeta(\mathbf{x})}, \zeta \! \uparrow \! \mathbf{v}_{\mathbf{x}} \rangle = \langle \, \mathcal{T}^*_{\zeta(\mathbf{x})} \zeta \circ \mathbf{v}^*_{\zeta(\mathbf{x})}, \mathbf{v}_{\mathbf{x}} \rangle$$

Covariant tensors

$$\boldsymbol{\zeta} \! \downarrow \! \boldsymbol{\mathsf{s}}_{\boldsymbol{\zeta}(\boldsymbol{\mathsf{x}})}^{\mathrm{COV}} = \mathit{T}_{\boldsymbol{\zeta}(\boldsymbol{\mathsf{x}})}^{*} \! \boldsymbol{\zeta} \circ \boldsymbol{\mathsf{s}}_{\boldsymbol{\zeta}(\boldsymbol{\mathsf{x}})}^{\mathrm{COV}} \circ \mathit{T}_{\boldsymbol{\mathsf{x}}} \boldsymbol{\zeta} \in \mathrm{COV}(\mathbb{TM})_{\boldsymbol{\mathsf{x}}}$$

Contravariant tensors

$$\zeta \uparrow \boldsymbol{s}_{\boldsymbol{x}}^{\mathrm{CON}} = \mathcal{T}_{\boldsymbol{x}} \boldsymbol{\zeta} \circ \boldsymbol{s}_{\boldsymbol{x}}^{\mathrm{CON}} \circ \mathcal{T}_{\boldsymbol{\zeta}(\boldsymbol{x})}^{*} \boldsymbol{\zeta} \in \mathrm{CON}(\mathbb{TN})_{\boldsymbol{\zeta}(\boldsymbol{x})}$$

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Covectors

$$\langle \zeta {\downarrow} \mathbf{v}^*_{\zeta(\mathbf{x})}, \mathbf{v}_{\mathbf{x}} \rangle = \langle \mathbf{v}^*_{\zeta(\mathbf{x})}, \zeta {\uparrow} \mathbf{v}_{\mathbf{x}} \rangle = \langle \, \mathcal{T}^*_{\zeta(\mathbf{x})} \zeta \circ \mathbf{v}^*_{\zeta(\mathbf{x})}, \mathbf{v}_{\mathbf{x}} \rangle$$

Covariant tensors

$$\boldsymbol{\zeta} \! \downarrow \! \boldsymbol{s}^{\mathrm{COV}}_{\boldsymbol{\zeta}(\boldsymbol{x})} = \boldsymbol{\mathcal{T}}^*_{\boldsymbol{\zeta}(\boldsymbol{x})} \boldsymbol{\zeta} \circ \boldsymbol{s}^{\mathrm{COV}}_{\boldsymbol{\zeta}(\boldsymbol{x})} \circ \boldsymbol{\mathcal{T}}_{\boldsymbol{x}} \boldsymbol{\zeta} \in \mathrm{COV}(\mathbb{TM})_{\boldsymbol{x}}$$

Contravariant tensors

$$\zeta \uparrow \mathbf{s}_{\mathbf{x}}^{\mathrm{CON}} = \mathit{T}_{\mathbf{x}} \zeta \circ \mathbf{s}_{\mathbf{x}}^{\mathrm{CON}} \circ \mathit{T}_{\zeta(\mathbf{x})}^{*} \zeta \in \mathrm{CON}(\mathbb{TN})_{\zeta(\mathbf{x})}$$

Mixed tensors

$$\zeta \uparrow \boldsymbol{s}_{\boldsymbol{x}}^{\mathrm{MIX}} = \mathcal{T}_{\boldsymbol{x}} \zeta \circ \boldsymbol{s}_{\boldsymbol{x}}^{\mathrm{MIX}} \circ \mathcal{T}_{\boldsymbol{\zeta}(\boldsymbol{x})} \boldsymbol{\zeta}^{-1} \in \mathrm{MIX}(\mathbb{TN})_{\boldsymbol{\zeta}(\boldsymbol{x})}$$

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Parallel transport along a curve $\mathbf{c} \in \mathrm{C}^1([a,b];\mathbb{M})$

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Covariance Paradigm

Parallel transport along a curve $\mathbf{c} \in \mathrm{C}^1([a,b];\mathbb{M})$

Vector fields

$$\begin{split} \mathbf{x} &= \mathbf{c}(\mu) \,, \quad \mathbf{v}_{\mathbf{x}} \in \mathbb{T}_{\mathbf{x}} \mathbb{M} \quad \mapsto \quad \mathbf{c}_{\lambda,\mu} \Uparrow \mathbf{v}_{\mathbf{x}} \in \mathbb{T}_{\mathbf{c}(\lambda)} \mathbb{M} \\ & \mathbf{c}_{\mu,\mu} \Uparrow \mathbf{v}_{\mathbf{x}} = \mathbf{v}_{\mathbf{x}} \\ & \mathbf{c}_{\lambda,\mu} \Uparrow \circ \mathbf{c}_{\mu,\nu} \Uparrow = \mathbf{c}_{\lambda,\nu} \Uparrow \end{split}$$

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Covariance Paradigm

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Vector fields

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• Covector fields $\mathbf{v}^*_{\mathbf{x}} \in \mathbb{T}^*_{\mathbf{x}}\mathbb{M}$ (naturality)

$$\langle \mathbf{c}_{\lambda,\mu} \Uparrow \mathbf{v}_{\mathbf{x}}^{*}, \mathbf{c}_{\lambda,\mu} \Uparrow \mathbf{v}_{\mathbf{x}} \rangle = \mathbf{c}_{\lambda,\mu} \Uparrow \langle \mathbf{v}_{\mathbf{x}}^{*}, \mathbf{v}_{\mathbf{x}} \rangle$$

Tensor fields (naturality)

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► Events manifold: E – four dimensional

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Covariance Paradigm

How to play the game

Kinematics

- ► Events manifold: E four dimensional
- Observer split into space-time: $S \times I$

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Covariance Paradigm

How to play the game

Kinematics

- ► Events manifold: E four dimensional
- Observer split into space-time: $S \times I$
- time is absolute (Classical Mechanics)

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Covariance Paradigm

How to play the game

Kinematics

- ► Events manifold: E four dimensional
- Observer split into space-time: $S \times I$
- time is absolute (Classical Mechanics)
- \blacktriangleright distance between simultaneous events \mapsto metric tensor

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• Time and space fibrations: $\gamma : E \mapsto S \times I$ (observer)

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Example

Covariance Paradigm

• Time and space fibrations: $\gamma : E \mapsto S \times I$ (observer)



 $egin{aligned} \pi_{I,\mathrm{E}} &= \pi_{I,(\mathcal{S} imes I)} \circ oldsymbol{\gamma} \ \pi_{\mathcal{S},\mathrm{E}} &= \pi_{\mathcal{S},(\mathcal{S} imes I)} \circ oldsymbol{\gamma} \end{aligned}$

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Covariance Paradigm

• Time and space fibrations: $\gamma : E \mapsto S \times I$ (observer)



$$\pi_{I,\mathrm{E}} = \pi_{I,(\mathcal{S} imes I)} \circ \gamma$$

 $\pi_{\mathcal{S},\mathrm{E}} = \pi_{\mathcal{S},(\mathcal{S} imes I)} \circ \gamma$

Time-vertical subbundle: spatial vectors

$$\mathbf{v} \in \mathbb{V}_{\mathbf{e}} \mathbf{E} \quad \Longleftrightarrow \quad T_{\mathbf{e}} \boldsymbol{\pi}_{I,\mathbf{E}} \cdot \mathbf{v} = \mathbf{0}$$

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Covariance Paradigm

• Time and space fibrations: $\gamma : E \mapsto S \times I$ (observer)



$$\pi_{I,\mathrm{E}} = \pi_{I,(\mathcal{S} imes I)} \circ \gamma$$
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Time-vertical subbundle: spatial vectors

$$\mathbf{v} \in \mathbb{V}_{\mathbf{e}} \mathbf{E} \quad \Longleftrightarrow \quad T_{\mathbf{e}} \boldsymbol{\pi}_{I,\mathbf{E}} \cdot \mathbf{v} = \mathbf{0}$$

 $\blacktriangleright \ \mathbf{v}_{\mathbf{e}} \in \mathbb{V}_{\mathbf{e}} \mathbf{E} \quad \Longleftrightarrow \quad \gamma \uparrow \mathbf{v}_{\mathbf{e}} = (v_{\mathbf{x},t}, \mathbf{0}_t) \in \mathbb{T}_{\mathbf{x}} \mathcal{S} \times \mathbb{T}_t I$

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Covariance Paradigm

 Trajectory: T_φ ⊂ E; subbundle of the events time-bundle ball (dim=3+1) membrane (dim=2+1) wire (dim=1+1) The G-Factor Impact in NLCM

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Covariance Paradigm

- Trajectory: T_φ ⊂ E; subbundle of the events time-bundle ball (dim=3+1) membrane (dim=2+1) wire (dim=1+1)
- ▶ time fibration \mapsto fibers: body placements $\mathbf{\Omega}_t$



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Covariance Paradigm

- Trajectory: T_φ ⊂ E; subbundle of the events time-bundle ball (dim=3+1) membrane (dim=2+1) wire (dim=1+1)
- \blacktriangleright time fibration \mapsto fibers: body placements Ω_t
- vertical tangent fibration \mapsto material vectors \mathbf{v}_{φ}



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- Evolution operator: φ



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- \blacktriangleright time fibration \mapsto fibers: body placements Ω_t
- vertical tangent fibration \mapsto material vectors \mathbf{v}_{φ}
- Evolution operator: φ
- ► Law of determinism (CHAPMAN-KOLMOGOROV):

$$\varphi_{ au,s} = \varphi_{ au,t} \circ \varphi_{t,s}$$

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- \blacktriangleright time fibration \mapsto fibers: body placements Ω_t
- vertical tangent fibration \mapsto material vectors \mathbf{v}_{φ}
- Evolution operator: φ
- ► Law of determinism (CHAPMAN-KOLMOGOROV):

$$arphi_{ au,s} = arphi_{ au,t} \circ arphi_{t,s}$$

Displacements: diffeomorphisms between placements

$$oldsymbol{arphi}_{ au,t}\in\mathrm{C}^1(oldsymbol{\Omega}_t\,;oldsymbol{\Omega}_ au)\,,\quad au,t\in I$$

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Covariance Paradigm



Equivalence relation on the trajectory:

$$(\mathbf{e}_1\,,\mathbf{e}_2)\in\mathcal{T}_{\boldsymbol{arphi}} imes\mathcal{T}_{\boldsymbol{arphi}}\,:\,\mathbf{e}_2=\boldsymbol{arphi}_{t_2,t_1}(\mathbf{e}_1)\,.$$

with $t_i = \boldsymbol{\pi}_{I,\mathrm{E}}(\mathbf{e}_i)$, i = 1, 2.

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Covariance Paradigm



• Equivalence relation on the trajectory:

$$(\mathbf{e}_1\,,\mathbf{e}_2)\in\mathcal{T}_{oldsymbol{arphi}} imes\mathcal{T}_{oldsymbol{arphi}}\,:\,\mathbf{e}_2=oldsymbol{arphi}_{t_2,t_1}(\mathbf{e}_1)\,.$$

with $t_i = \boldsymbol{\pi}_{I,\mathrm{E}}(\mathbf{e}_i)$, i = 1, 2.

Body = quotient manifold (foliation)

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Covariance Paradigm



• Equivalence relation on the trajectory:

$$(\mathbf{e}_1\,,\mathbf{e}_2)\in\mathcal{T}_{oldsymbol{arphi}} imes\mathcal{T}_{oldsymbol{arphi}}\,:\,\mathbf{e}_2=oldsymbol{arphi}_{t_2,t_1}(\mathbf{e}_1)\,.$$

with $t_i = \boldsymbol{\pi}_{I,\mathrm{E}}(\mathbf{e}_i)$, i = 1, 2.

Body = quotient manifold (foliation) Particles = equivalence classes (folia) The G-Factor Impact in NLCM

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Covariance Paradigm



• Equivalence relation on the trajectory:

$$(\mathbf{e}_1\,,\mathbf{e}_2)\in\mathcal{T}_{oldsymbol{arphi}} imes\mathcal{T}_{oldsymbol{arphi}}\,:\,\mathbf{e}_2=oldsymbol{arphi}_{t_2,t_1}(\mathbf{e}_1)\,.$$

with $t_i = \boldsymbol{\pi}_{I,\mathrm{E}}(\mathbf{e}_i)$, i = 1, 2.

Body = quotient manifold (foliation) Particles = equivalence classes (folia)

mass conservation

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Tensor bundles

spatial tensor bundles:

 $oldsymbol{ au}_{ ext{E}}^{ ext{TENS}} \in ext{C}^1(ext{TENS}(extsf{VE}) ext{; E})$

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Tensor bundles

- spatial tensor bundles:
- material tensor bundles:

 $oldsymbol{ au}_{ ext{E}}^{ ext{TENS}} \in ext{C}^{1}(ext{TENS}(extsf{VE}); ext{E}) \ oldsymbol{ au}_{oldsymbol{arphi}}^{ ext{TENS}} \in ext{C}^{1}(ext{TENS}(extsf{V}\mathcal{T}_{oldsymbol{arphi}}); \mathcal{T}_{oldsymbol{arphi}})$

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Covariance Paradigm

Tensor bundles

- spatial tensor bundles:
- material tensor bundles: $\tau_{\varphi}^{\text{TENS}} \in C^1(\text{TENS}(\mathbb{V}\mathcal{T}_{\varphi});\mathcal{T}_{\varphi})$
- material-based spatial tensor bundles:

 $oldsymbol{ au}_{\mathrm{E},oldsymbol{arphi}}^{\mathrm{TENS}} \in \mathrm{C}^1(\mathrm{TENS}(\mathbb{V}\mathrm{E})_{\mathcal{T}_{oldsymbol{arphi}}}\,;\mathcal{T}_{oldsymbol{arphi}})$

 $\boldsymbol{\tau}_{\mathrm{E}}^{\mathrm{TENS}} \in \mathrm{C}^{1}(\mathrm{TENS}(\mathbb{V}\mathrm{E});\mathrm{E})$

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Tensor bundles

- ▶ spatial tensor bundles: $au_{
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 m E})$
- material tensor bundles: $\tau_{\omega}^{\text{TENS}} \in C^1(\text{TENS}(\mathbb{V}\mathcal{T}_{\omega});\mathcal{T}_{\omega})$
- material-based spatial tensor bundles:

 $oldsymbol{ au}_{\mathrm{E},oldsymbol{arphi}}^{\mathrm{TENS}} \in \mathrm{C}^1(\mathrm{TENS}(\mathbb{V}\mathrm{E})_{\mathcal{T}_{oldsymbol{arphi}}}\,;\mathcal{T}_{oldsymbol{arphi}})$

Tensor fields (sections of the bundles)

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Tensor fields (sections of the bundles)

▶ spatial tensor fields: $s_E^{TENS} \in C^1(E; TENS(VE))$

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Tensor fields (sections of the bundles)

- ▶ spatial tensor fields: $s_E^{TENS} \in C^1(E; TENS(VE))$
- ▶ material tensor fields: $\mathbf{s}_{\boldsymbol{\varphi}}^{\mathrm{TENS}} \in \mathrm{C}^{1}(\mathcal{T}_{\boldsymbol{\varphi}}; \mathrm{TENS}(\mathbb{V}\mathcal{T}_{\boldsymbol{\varphi}}))$

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Tensor fields (sections of the bundles)

- ▶ spatial tensor fields: $s_E^{TENS} \in C^1(E; TENS(VE))$
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- material-based spatial tensor fields:

 $\mathbf{s}_{\mathrm{E}, oldsymbol{arphi}}^{\mathrm{TENS}} \in \mathrm{C}^1(\mathcal{T}_{oldsymbol{arphi}} \,; \mathrm{TENS}(\mathbb{V}\mathrm{E}))$

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- ▶ spatial tensor bundles: $au_{
 m E}^{
 m TENS} \in {
 m C}^1({
 m TENS}({
 m VE});{
 m E})$
- material tensor bundles: $\tau_{\varphi}^{\text{TENS}} \in C^1(\text{TENS}(\mathbb{V}\mathcal{T}_{\varphi});\mathcal{T}_{\varphi})$
- material-based spatial tensor bundles:

 $oldsymbol{ au}_{\mathrm{E},oldsymbol{arphi}}^{\mathrm{TENS}} \in \mathrm{C}^1(\mathrm{TENS}(\mathbb{V}\mathrm{E})_{\mathcal{T}_{oldsymbol{arphi}}}\,;\mathcal{T}_{oldsymbol{arphi}})$

Tensor fields (sections of the bundles)

- ▶ spatial tensor fields: $s_E^{TENS} \in C^1(E; TENS(VE))$
- ▶ material tensor fields: $\mathbf{s}_{\boldsymbol{\varphi}}^{\text{TENS}} \in C^1(\mathcal{T}_{\boldsymbol{\varphi}}; \text{TENS}(\mathbb{V}\mathcal{T}_{\boldsymbol{\varphi}}))$
- material-based spatial tensor fields:

 $\mathbf{s}_{\mathrm{E}, oldsymbol{arphi}}^{\mathrm{TENS}} \in \mathrm{C}^1(\mathcal{T}_{oldsymbol{arphi}}\,; \mathrm{TENS}(\mathbb{V}\mathrm{E}))$

such that: $\boldsymbol{ au}_{\mathrm{E}}^{\mathrm{TENS}} \circ \boldsymbol{s}_{\mathrm{E}}^{\mathrm{TENS}} = \mathrm{ID}_{\mathrm{E}}$.

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spatial field: metric tensor field

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Examples

- spatial field: metric tensor field
- material fields: strain, stress, stretching, stressing, thermal gradient, temperature, free energy, entropy etc.

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Examples

- spatial field: metric tensor field
- material fields: strain, stress, stretching, stressing, thermal gradient, temperature, free energy, entropy etc.
- material-based spatial fields: velocity, acceleration, kinetic momentum.

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Covariance Paradigm

Material fields at different times along a trajectory must be compared by push along the material displacement. Material fields on push-related trajectories must be compared by push along the relative motion. The G-Factor Impact in NLCM

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Material fields at different times along a trajectory must be compared by push along the material displacement. Material fields on push-related trajectories must be compared by push along the relative motion.

Push and parallel transport along the motion

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Material fields at different times along a trajectory must be compared by push along the material displacement. Material fields on push-related trajectories must be compared by push along the relative motion.

Push and parallel transport along the motion



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Material fields at different times along a trajectory must be compared by push along the material displacement. Material fields on push-related trajectories must be compared by push along the relative motion.

Push and parallel transport along the motion



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Time derivatives along the motion

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Covariance Paradigm
Time derivatives along the motion

Lie time derivative - LTD (Convective time derivative - CTD)

Material tensor field

$$\dot{\mathbf{s}}_{oldsymbol{arphi},t} := \mathcal{L}_{oldsymbol{arphi},t} \, \mathbf{s}_{oldsymbol{arphi}} = \partial_{ au=t} \left(oldsymbol{arphi}_{ au,t} {oldsymbol{arphi}}_{ au, au}
ight)$$

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Time derivatives along the motion

Lie time derivative - LTD (Convective time derivative - CTD)

Material tensor field

Material time-derivative - MTD (Parallel time-derivative - PTD)

Material-based spatial fields

$$\dot{\mathbf{s}}_{\mathrm{E}, \boldsymbol{\varphi}, t} :=
abla_{\boldsymbol{\varphi}, t} \, \mathbf{s}_{\mathrm{E}, \boldsymbol{\varphi}} = \partial_{ au = t} \, \boldsymbol{\varphi}_{ au, t} \Downarrow \, \mathbf{s}_{\mathrm{E}, \boldsymbol{\varphi}, au}$$

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LTD of a material field

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LTD of a material field

$$\begin{split} \dot{\mathbf{s}}_{\boldsymbol{\varphi},t}(\mathbf{x}) &:= (\mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi}})_{\mathbf{x}} = \partial_{\tau=t} \left(\boldsymbol{\varphi}_{\tau,t} \downarrow \mathbf{s}_{\boldsymbol{\varphi},\tau} \right)_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \downarrow (\mathbf{s}_{\boldsymbol{\varphi},\tau} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{s}_{\boldsymbol{\varphi},\tau}(\mathbf{x}) + \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \downarrow (\mathbf{s}_{\boldsymbol{\varphi},t} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{s}_{\boldsymbol{\varphi},\tau}(\mathbf{x}) + \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi},t}(\mathbf{x}) \end{split}$$

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LTD of a material field

$$\begin{split} \dot{\mathbf{s}}_{\boldsymbol{\varphi},t}(\mathbf{x}) &:= (\mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi}})_{\mathbf{x}} = \partial_{\tau=t} \, (\boldsymbol{\varphi}_{\tau,t} \! \mid \! \mathbf{s}_{\boldsymbol{\varphi},\tau})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \! \mid \! (\mathbf{s}_{\boldsymbol{\varphi},\tau} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{s}_{\boldsymbol{\varphi},\tau}(\mathbf{x}) + \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \! \mid \! (\mathbf{s}_{\boldsymbol{\varphi},t} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{s}_{\boldsymbol{\varphi},\tau}(\mathbf{x}) + \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi},t}(\mathbf{x}) \end{split}$$

MTD of the velocity field - Acceleration

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LTD of a material field

$$\begin{split} \dot{\mathbf{s}}_{\boldsymbol{\varphi},t}(\mathbf{x}) &:= (\mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi}})_{\mathbf{x}} = \partial_{\tau=t} \, (\boldsymbol{\varphi}_{\tau,t} \! \mid \! \mathbf{s}_{\boldsymbol{\varphi},\tau})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \! \mid \! (\mathbf{s}_{\boldsymbol{\varphi},\tau} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{s}_{\boldsymbol{\varphi},\tau}(\mathbf{x}) + \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \! \mid \! (\mathbf{s}_{\boldsymbol{\varphi},t} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{s}_{\boldsymbol{\varphi},\tau}(\mathbf{x}) + \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi},t}(\mathbf{x}) \end{split}$$

MTD of the velocity field - Acceleration

$$\begin{split} \mathbf{a}_{\mathrm{E},\boldsymbol{\varphi},t}(\mathbf{x}) &:= (\nabla_{\boldsymbol{\varphi},t} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi}})_{\mathbf{x}} = \partial_{\tau=t} \, (\boldsymbol{\varphi}_{\tau,t} \Downarrow \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \Downarrow (\mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau}(\mathbf{x}) + \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \Downarrow (\mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},t} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau}(\mathbf{x}) + \nabla_{\mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},t}} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},t}(\mathbf{x}) \end{split}$$

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LTD of a material field

$$\begin{split} \dot{\mathbf{s}}_{\boldsymbol{\varphi},t}(\mathbf{x}) &:= (\mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi}})_{\mathbf{x}} = \partial_{\tau=t} \, (\boldsymbol{\varphi}_{\tau,t} \! \mid \! \mathbf{s}_{\boldsymbol{\varphi},\tau})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \! \mid \! (\mathbf{s}_{\boldsymbol{\varphi},\tau} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{s}_{\boldsymbol{\varphi},\tau}(\mathbf{x}) + \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \! \mid \! (\mathbf{s}_{\boldsymbol{\varphi},t} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{s}_{\boldsymbol{\varphi},\tau}(\mathbf{x}) + \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi},t}(\mathbf{x}) \end{split}$$

MTD of the velocity field - Acceleration

$$\begin{aligned} \mathbf{a}_{\mathrm{E},\boldsymbol{\varphi},t}(\mathbf{x}) &:= (\nabla_{\boldsymbol{\varphi},t} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi}})_{\mathbf{x}} = \partial_{\tau=t} \left(\varphi_{\tau,t} \Downarrow \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau} \right)_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \varphi_{\tau,t} \Downarrow \left(\mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau} \circ \varphi_{\tau,t} \right)_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau}(\mathbf{x}) + \partial_{\tau=t} \, \varphi_{\tau,t} \Downarrow \left(\mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},t} \circ \varphi_{\tau,t} \right)_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau}(\mathbf{x}) + \nabla_{\mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},t}} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},t}(\mathbf{x}) \end{aligned}$$

The latter is D'ALEMBERT-D. BERNOULLI formula, applicable only in special problems in hydrodynamics, where it was conceived. This eventually led to the NAVIER-STOKES-ST.VENANT differential equation of motion

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Covariance Paradigm Time derivatives



LTD of a material field

$$\begin{split} \dot{\mathbf{s}}_{\boldsymbol{\varphi},t}(\mathbf{x}) &:= (\mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi}})_{\mathbf{x}} = \partial_{\tau=t} \left(\boldsymbol{\varphi}_{\tau,t} \! \downarrow \! \mathbf{s}_{\boldsymbol{\varphi},\tau} \right)_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \! \downarrow \! \left(\mathbf{s}_{\boldsymbol{\varphi},\tau} \circ \boldsymbol{\varphi}_{\tau,t} \right)_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{s}_{\boldsymbol{\varphi},\tau}(\mathbf{x}) + \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \! \downarrow \! \left(\mathbf{s}_{\boldsymbol{\varphi},t} \circ \boldsymbol{\varphi}_{\tau,t} \right)_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{s}_{\boldsymbol{\varphi},\tau}(\mathbf{x}) + \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi},t}(\mathbf{x}) \end{split}$$

MTD of the velocity field - Acceleration

$$\begin{split} \mathbf{a}_{\mathrm{E},\boldsymbol{\varphi},t}(\mathbf{x}) &:= (\nabla_{\boldsymbol{\varphi},t} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi}})_{\mathbf{x}} = \partial_{\tau=t} \, (\boldsymbol{\varphi}_{\tau,t} \Downarrow \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \Downarrow (\mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau}(\mathbf{x}) + \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \Downarrow (\mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},t} \circ \boldsymbol{\varphi}_{\tau,t})_{\mathbf{x}} \\ &= \partial_{\tau=t} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau}(\mathbf{x}) + \nabla_{\mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},t}} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},t}(\mathbf{x}) \end{split}$$

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$$\dot{\mathbf{s}}_{oldsymbol{arphi},t}(\mathbf{x}) := (\mathcal{L}_{oldsymbol{arphi},t} \, \mathbf{s}_{oldsymbol{arphi}})_{\mathbf{x}}$$

 $\mathbf{a}_{\mathrm{E},oldsymbol{arphi},t}(\mathbf{x}) := (
abla_{oldsymbol{arphi},t} \, \mathbf{v}_{\mathrm{E},oldsymbol{arphi}})_{\mathbf{x}}$



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$$\begin{split} \dot{\mathbf{s}}_{\boldsymbol{\varphi},t}(\mathbf{x}) &:= (\mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi}})_{\mathbf{x}} &= \partial_{\tau=t} \, \mathbf{s}_{\boldsymbol{\varphi},\tau}(\mathbf{x}) + \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\boldsymbol{\varphi},t}(\mathbf{x}) \\ \mathbf{a}_{\mathrm{E},\boldsymbol{\varphi},t}(\mathbf{x}) &:= (\nabla_{\boldsymbol{\varphi},t} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi}})_{\mathbf{x}} &= \partial_{\tau=t} \, \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},\tau}(\mathbf{x}) + \nabla_{\mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},t}} \, \mathbf{v}_{\mathrm{E},\boldsymbol{\varphi},t}(\mathbf{x}) \end{split}$$

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$$\begin{split} \dot{\mathbf{s}}_{\varphi,t}(\mathbf{x}) &:= (\mathcal{L}_{\varphi,t} \, \mathbf{s}_{\varphi})_{\mathbf{x}} \quad = \partial_{\tau=t} \, \mathbf{s}_{\varphi,\tau}(\mathbf{x}) + \mathcal{L}_{\varphi,t} \, \mathbf{s}_{\varphi,t}(\mathbf{x}) \\ \mathbf{a}_{\mathrm{E},\varphi,t}(\mathbf{x}) &:= (\nabla_{\varphi,t} \, \mathbf{v}_{\mathrm{E},\varphi})_{\mathbf{x}} \quad = \partial_{\tau=t} \, \, \mathbf{v}_{\mathrm{E},\varphi,\tau}(\mathbf{x}) + \nabla_{\mathbf{v}_{\mathrm{E},\varphi,t}} \, \mathbf{v}_{\mathrm{E},\varphi,t}(\mathbf{x}) \end{split}$$

In fact LEIBNIZ rule cannot be applied unless the following special properties of the trajectory hold true:

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$$\begin{split} \dot{\mathbf{s}}_{\varphi,t}(\mathbf{x}) &:= (\mathcal{L}_{\varphi,t} \, \mathbf{s}_{\varphi})_{\mathbf{x}} &= \partial_{\tau=t} \, \mathbf{s}_{\varphi,\tau}(\mathbf{x}) + \mathcal{L}_{\varphi,t} \, \mathbf{s}_{\varphi,t}(\mathbf{x}) \\ \mathbf{a}_{\mathrm{E},\varphi,t}(\mathbf{x}) &:= (\nabla_{\varphi,t} \, \mathbf{v}_{\mathrm{E},\varphi})_{\mathbf{x}} &= \partial_{\tau=t} \, \mathbf{v}_{\mathrm{E},\varphi,\tau}(\mathbf{x}) + \nabla_{\mathbf{v}_{\mathrm{E},\varphi,t}} \, \mathbf{v}_{\mathrm{E},\varphi,t}(\mathbf{x}) \end{split}$$

In fact LEIBNIZ rule cannot be applied unless the following special properties of the trajectory hold true:

$$\begin{aligned} & (\mathbf{x},t) \in \mathcal{T}_{\boldsymbol{\varphi}} & \Longrightarrow \quad (\mathbf{x},\tau) \in \mathcal{T}_{\boldsymbol{\varphi}} \quad \forall \, \tau \in I_t \\ & (\mathbf{x},t) \in \mathcal{T}_{\boldsymbol{\varphi}} & \Longrightarrow \quad (\boldsymbol{\varphi}_{\tau,t}(\mathbf{x}),t) \in \mathcal{T}_{\boldsymbol{\varphi}} \end{aligned}$$

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$$\begin{split} \dot{\mathbf{s}}_{\varphi,t}(\mathbf{x}) &:= (\mathcal{L}_{\varphi,t} \, \mathbf{s}_{\varphi})_{\mathbf{x}} &= \partial_{\tau=t} \, \mathbf{s}_{\varphi,\tau}(\mathbf{x}) + \mathcal{L}_{\varphi,t} \, \mathbf{s}_{\varphi,t}(\mathbf{x}) \\ \mathbf{a}_{\mathrm{E},\varphi,t}(\mathbf{x}) &:= (\nabla_{\varphi,t} \, \mathbf{v}_{\mathrm{E},\varphi})_{\mathbf{x}} &= \partial_{\tau=t} \, \mathbf{v}_{\mathrm{E},\varphi,\tau}(\mathbf{x}) + \nabla_{\mathbf{v}_{\mathrm{E},\varphi,t}} \, \mathbf{v}_{\mathrm{E},\varphi,t}(\mathbf{x}) \end{split}$$

In fact LEIBNIZ rule cannot be applied unless the following special properties of the trajectory hold true:

$$(\mathbf{x},t) \in \mathcal{T}_{\boldsymbol{\varphi}} \implies (\mathbf{x},\tau) \in \mathcal{T}_{\boldsymbol{\varphi}} \quad \forall \tau \in I_t$$

$$(\mathbf{x}\,,t)\in\mathcal{T}_{oldsymbol{arphi}} \implies (oldsymbol{arphi}_{ au,t}(\mathbf{x})\,,t)\in\mathcal{T}_{oldsymbol{arphi}}$$

Both conditions are not fulfilled in solid mechanics, as a rule.

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A sample of objective stress tensors

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A sample of objective stress tensors

Convective stress tensor rates in TRUESDELL & NOLL (1965):

 $\stackrel{\Delta}{\mathbf{T}} = \dot{\mathbf{T}} + \mathbf{L}^T \mathbf{T} + \mathbf{T} \mathbf{L}$

Co-rotational stress tensor rates in TRUESDELL & NOLL (1965):

$$\overset{\circ}{\mathbf{T}} = \dot{\mathbf{T}} - \mathbf{W}\mathbf{T} + \mathbf{T}\mathbf{W}$$

with $\dot{\mathbf{T}}$ material time derivative

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A sample of objective stress tensors

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with $\dot{\mathbf{T}}$ material time derivative

Both formulas rely on LEIBNIZ rule and on treating the material stress tensor field as a spatial valued tensor field

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Norm axioms



$\begin{aligned} \|\mathbf{a}\| \ge 0, \quad \|\mathbf{a}\| = 0 \implies \mathbf{a} = 0\\ \|\mathbf{a}\| + \|\mathbf{b}\| \ge \|\mathbf{c}\| \quad \text{triangle inequality,}\\ \|\alpha \, \mathbf{a}\| = |\alpha| \|\mathbf{a}\| \end{aligned}$

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Norm axioms

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Parallelogram rule



$$\|\mathbf{a} + \mathbf{b}\|^2 + \|\mathbf{a} - \mathbf{b}\|^2 = 2 [\|\mathbf{a}\|^2 + \|\mathbf{b}\|^2]$$

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The metric tensor

Theorem (Fréchet – von Neumann – Jordan)

$$\mathbf{g}(\mathbf{a},\mathbf{b}) := \frac{1}{4} \big[\|\mathbf{a} + \mathbf{b}\|^2 - \|\mathbf{a} - \mathbf{b}\|^2 \big]$$

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Theorem (Fréchet – von Neumann – Jordan)

$$\mathbf{g}(\mathbf{a}\,,\mathbf{b}):=rac{1}{4}ig[\|\mathbf{a}+\mathbf{b}\|^2-\|\mathbf{a}-\mathbf{b}\|^2ig]$$



Kosaku Josida (1909 - 1990)

$$t \begin{bmatrix} g(e_1, e_1) \cdots g(e_1, e_3) \\ \vdots \\ g(e_3, e_1) \cdots g(e_3, e_3) \\ g(e_3, e_1) \cdots g(e_3, e_3) \\ \vdots \\ g(e_3, e_1) \cdots g(e_3, e_3) \\ \vdots \\ g(e_3, e_1) \cdots g(e_3, e_3) \\ \vdots \\ g(e_3, e_1) \\ g(e_1, e_1) \\ g(e_3, e_1) \\ g(e_3, e_1$$

The metric tensor

Theorem (Fréchet – von Neumann – Jordan)

$$\mathbf{g}(\mathbf{a}\,,\mathbf{b}):=rac{1}{4}ig[\|\mathbf{a}+\mathbf{b}\|^2-\|\mathbf{a}-\mathbf{b}\|^2ig]$$



Maurice René Fréchet (1878 - 1973)

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The metric tensor

Theorem (Fréchet – von Neumann – Jordan)

$$\mathbf{g}(\mathbf{a}\,,\mathbf{b}):=rac{1}{4}ig[\|\mathbf{a}+\mathbf{b}\|^2-\|\mathbf{a}-\mathbf{b}\|^2ig]$$





John von Neumann (1903 - 1957)

 $)^{2} = \mathsf{det} \begin{bmatrix} \mathbf{g}(\mathbf{e}_{1}\,,\mathbf{e}_{1})\cdots\mathbf{g}(\mathbf{e}_{1}\,,\mathbf{e}_{3})\\ \cdots\\ \mathbf{g}(\mathbf{e}_{3}\,,\mathbf{e}_{1})\cdots\mathbf{g}(\mathbf{e}_{3}\,,\mathbf{e}_{3}) \end{bmatrix}$

The metric tensor

Theorem (Fréchet – von Neumann – Jordan)

$$\mathbf{g}(\mathbf{a}\,,\mathbf{b}):=rac{1}{4}ig[\|\mathbf{a}+\mathbf{b}\|^2-\|\mathbf{a}-\mathbf{b}\|^2ig]$$





Pascual Jordan (1902 - 1980)

$$= \operatorname{Jordan})$$

$$= \|\mathbf{a} - \mathbf{b}\|^{2}]$$

$$= \operatorname{det} \begin{bmatrix} \mathbf{g}(\mathbf{e}_{1}, \mathbf{e}_{1}) \cdots \mathbf{g}(\mathbf{e}_{1}, \mathbf{e}_{3}) \\ \cdots \\ \mathbf{g}(\mathbf{e}_{3}, \mathbf{e}_{1}) \cdots \mathbf{g}(\mathbf{e}_{3}, \mathbf{e}_{3}) \end{bmatrix}$$

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• Metric tensor field: $\mathbf{g} \in \mathrm{C}^1(\mathcal{S}; \mathrm{COV}(\mathbb{T}\mathcal{S}))$

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Covariance Paradigm

- Metric tensor field: $\mathbf{g} \in \mathrm{C}^1(\mathcal{S}; \mathrm{COV}(\mathbb{T}\mathcal{S}))$
- ► Spatial metric tensor field (on the events manifold) $\mathbf{g}_{\mathrm{E}}(\mathbf{a}, \mathbf{b}) := \mathbf{g}(T\pi_{\mathcal{S},\mathrm{E}} \cdot \mathbf{a}, T\pi_{\mathcal{S},\mathrm{E}} \cdot \mathbf{b}), \quad \mathbf{a}, \mathbf{b} \in \mathbb{V}\mathrm{E}$

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- Spatial immersion of material vectors

$$\mathbf{i}_{\mathrm{E},\mathcal{T}_{oldsymbol{arphi}}}\in\mathrm{C}^{1}(\mathcal{T}_{oldsymbol{arphi}}\,;\mathrm{E})$$

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- Spatial immersion of material vectors

 $\mathbf{i}_{\mathrm{E},\mathcal{T}_{oldsymbol{arphi}}}\in\mathrm{C}^{1}(\mathcal{T}_{oldsymbol{arphi}}\,;\mathrm{E})$

Material metric tensor field (pull back)

$$\mathbf{g}_{\boldsymbol{\varphi}} := \mathbf{i}_{\mathrm{E},\mathcal{T}_{\boldsymbol{\varphi}}} \! \! \downarrow \! \mathbf{g}_{\mathrm{E}} := \mathcal{T}^* \mathbf{i}_{\mathrm{E},\mathcal{T}_{\boldsymbol{\varphi}}} \circ \mathbf{g}_{\mathrm{E}} \circ \mathcal{T} \mathbf{i}_{\mathrm{E},\mathcal{T}_{\boldsymbol{\varphi}}}$$

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Time derivatives

Leonhard Euler (1707 - 1783)

Lie (convective) time derivative

Leonhard Euler (1707 - 1783)



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Lie (convective) time derivative

• Stretching: $\dot{\boldsymbol{\varepsilon}}_{\boldsymbol{\varphi},t} := \frac{1}{2} \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{g}_{\boldsymbol{\varphi}} = \frac{1}{2} \partial_{\tau=t} \left(\boldsymbol{\varphi}_{\tau,t} \downarrow \mathbf{g}_{\boldsymbol{\varphi},\tau} \right)$



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Leonhard Euler (1707 - 1783)

Lie (convective) time derivative

- $\blacktriangleright \text{ Stretching:} \quad \dot{\boldsymbol{\varepsilon}}_{\boldsymbol{\varphi},t} := \frac{1}{2} \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{g}_{\boldsymbol{\varphi}} = \frac{1}{2} \partial_{\tau=t} \left(\boldsymbol{\varphi}_{\tau,t} \downarrow \mathbf{g}_{\boldsymbol{\varphi},\tau} \right)$
- Euler's formula (generalized)

$$\frac{1}{2}\mathcal{L}_{\boldsymbol{\varphi},t}\,\mathbf{g}_{\boldsymbol{\varphi}}=\frac{1}{2}\nabla^{\text{MAT}}_{\mathbf{v}_{\boldsymbol{\varphi},t}}\,\mathbf{g}_{\boldsymbol{\varphi},t}+\text{sym}\,(\mathbf{g}_{\boldsymbol{\varphi},t}\circ(\text{Tors}^{\text{MAT}}+\nabla^{\text{MAT}})(\mathbf{v}_{\boldsymbol{\varphi},t}))$$

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Lie (convective) time derivative

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• where
$$\mathbf{g}_{\boldsymbol{\varphi},t} \circ \nabla^{\text{MAT}} \mathbf{v}_{\boldsymbol{\varphi},t} = \mathbf{i}_{\boldsymbol{\varphi},t} \downarrow (\mathbf{g} \circ \nabla \mathbf{v}_{\boldsymbol{\varphi},t})$$

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Lie (convective) time derivative

- $\blacktriangleright \text{ Stretching:} \quad \dot{\boldsymbol{\varepsilon}}_{\boldsymbol{\varphi},t} := \frac{1}{2} \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{g}_{\boldsymbol{\varphi}} = \frac{1}{2} \partial_{\tau=t} \left(\boldsymbol{\varphi}_{\tau,t} \downarrow \mathbf{g}_{\boldsymbol{\varphi},\tau} \right)$
- Euler's formula (generalized)

$$\frac{1}{2}\mathcal{L}_{\boldsymbol{\varphi},t}\,\mathbf{g}_{\boldsymbol{\varphi}}=\frac{1}{2}\nabla^{\text{MAT}}_{\mathbf{v}_{\boldsymbol{\varphi},t}}\,\mathbf{g}_{\boldsymbol{\varphi},t}+\text{sym}\left(\mathbf{g}_{\boldsymbol{\varphi},t}\circ(\text{Tors}^{\text{MAT}}+\nabla^{\text{MAT}})(\mathbf{v}_{\boldsymbol{\varphi},t})\right)$$

 where g_{φ,t} ◦ ∇^{MAT} v_{φ,t} = i_{φ,t}↓(g ◦ ∇ v_{φ,t})
 with g_{φ,t} ◦ ∇^{MAT} v_{φ,t} ∈ C¹(Ω_t; COV(TΩ_t)) g ◦ ∇ v_{φ,t} ∈ C¹(Ω_t; COV(TΩ_tS))

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Lie (convective) time derivative

- $\blacktriangleright \text{ Stretching:} \quad \dot{\boldsymbol{\varepsilon}}_{\boldsymbol{\varphi},t} := \frac{1}{2} \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{g}_{\boldsymbol{\varphi}} = \frac{1}{2} \partial_{\tau=t} \left(\boldsymbol{\varphi}_{\tau,t} \downarrow \mathbf{g}_{\boldsymbol{\varphi},\tau} \right)$
- Euler's formula (generalized)

$$\frac{1}{2}\mathcal{L}_{\boldsymbol{\varphi},t}\,\mathbf{g}_{\boldsymbol{\varphi}}=\frac{1}{2}\nabla^{\text{MAT}}_{\mathbf{v}_{\boldsymbol{\varphi},t}}\,\mathbf{g}_{\boldsymbol{\varphi},t}+\text{sym}\left(\mathbf{g}_{\boldsymbol{\varphi},t}\circ(\text{Tors}^{\text{MAT}}+\nabla^{\text{MAT}})(\mathbf{v}_{\boldsymbol{\varphi},t})\right)$$

• where
$$\mathbf{g}_{\varphi,t} \circ \nabla^{\mathrm{MAT}} \mathbf{v}_{\varphi,t} = \mathbf{i}_{\varphi,t} \downarrow (\mathbf{g} \circ \nabla \mathbf{v}_{\varphi,t})$$

with

$$\begin{split} & \mathbf{g}_{\boldsymbol{\varphi},t} \circ \nabla^{\text{MAT}} \mathbf{v}_{\boldsymbol{\varphi},t} \in \mathrm{C}^{1}(\boldsymbol{\Omega}_{t}\,; \mathrm{COV}(\mathbb{T}\boldsymbol{\Omega}_{t})) \\ & \mathbf{g} \circ \nabla \mathbf{v}_{\boldsymbol{\varphi},t} \in \mathrm{C}^{1}(\boldsymbol{\Omega}_{t}\,; \mathrm{COV}(\mathbb{T}_{\boldsymbol{\Omega}_{t}}\mathcal{S})) \end{split}$$

Mixed form of the stretching tensor (standard):

$$\mathbf{D}_{\boldsymbol{\varphi},t} := \mathbf{g}_{\boldsymbol{\varphi},t}^{-1} \circ \frac{1}{2} \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{g}_{\boldsymbol{\varphi}} = \operatorname{sym}\left(\nabla^{\text{MAT}} \mathbf{v}_{\boldsymbol{\varphi},t}\right)$$

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Lie (convective) time derivative

Stress: σ_φ ∈ C¹(T_φ; CON(VT_φ)) contravariant material tensor field in duality with the stretching covariant material tensor field: έ_{φ,t} := ½L_{φ,t} g_φ ∈ C¹(T_φ; COV(VT_φ)) The G-Factor Impact in NLCM

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Stressing:

$$\dot{\sigma}_{oldsymbol{arphi},t} \coloneqq \mathcal{L}_{oldsymbol{arphi},t} \, \sigma_{oldsymbol{arphi}} = \partial_{ au=t} \left(oldsymbol{arphi}_{ au,t} {\downarrow} oldsymbol{\sigma}_{oldsymbol{arphi}, au}
ight)$$

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Stressing:

$$\dot{\boldsymbol{\sigma}}_{\boldsymbol{arphi},t} \coloneqq \mathcal{L}_{\boldsymbol{arphi},t} \, \boldsymbol{\sigma}_{\boldsymbol{arphi}} = \partial_{ au=t} \left(\boldsymbol{arphi}_{ au,t} {oldsymbol{arphi}}_{ au, au}
ight)$$

▶ Spatial contravariant tensor: $s_E \in C^1(E; CON(VE))$

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Stressing:

$$\dot{\sigma}_{oldsymbol{arphi},t} := \mathcal{L}_{oldsymbol{arphi},t} \, \sigma_{oldsymbol{arphi}} = \partial_{ au = t} \left(oldsymbol{arphi}_{ au,t} {oldsymbol{arphi}}_{oldsymbol{arphi}, au}
ight)$$

- ► Spatial contravariant tensor: s_E ∈ C¹(E; CON(VE))
- Leibniz rule (applicable to spatial tensor fields)

$$\begin{split} \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\mathrm{E}} &:= \partial_{\tau=t} \left(\boldsymbol{\varphi}_{\tau,t} \! \downarrow \! \mathbf{s}_{\mathrm{E},\tau} \right) = \partial_{\tau=t} \, \, \mathbf{s}_{\mathrm{E},\tau} + \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \! \downarrow \! \left(\mathbf{s}_{\mathrm{E},t} \circ \boldsymbol{\varphi}_{\tau,t} \right) \\ &= \partial_{\tau=t} \, \, \mathbf{s}_{\mathrm{E},\tau} + \mathcal{L}_{\mathbf{v}_{\boldsymbol{\varphi},t}} \, \mathbf{s}_{\mathrm{E},t} \end{split}$$

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Stressing:

$$\dot{\sigma}_{oldsymbol{arphi},t} := \mathcal{L}_{oldsymbol{arphi},t} \, \sigma_{oldsymbol{arphi}} = \partial_{ au = t} \left(oldsymbol{arphi}_{ au,t} {oldsymbol{arphi}}_{oldsymbol{arphi}, au}
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- ▶ Spatial contravariant tensor: $s_E \in C^1(E; CON(VE))$
- Leibniz rule (applicable to spatial tensor fields)

$$\begin{split} \mathcal{L}_{\boldsymbol{\varphi},t} \, \mathbf{s}_{\mathrm{E}} &:= \partial_{\tau=t} \left(\boldsymbol{\varphi}_{\tau,t} \!\downarrow \! \mathbf{s}_{\mathrm{E},\tau} \right) = \partial_{\tau=t} \, \, \mathbf{s}_{\mathrm{E},\tau} + \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,t} \!\downarrow \! \left(\mathbf{s}_{\mathrm{E},t} \circ \boldsymbol{\varphi}_{\tau,t} \right) \\ &= \partial_{\tau=t} \, \, \mathbf{s}_{\mathrm{E},\tau} + \mathcal{L}_{\mathbf{v}\boldsymbol{\varphi},t} \, \mathbf{s}_{\mathrm{E},t} \end{split}$$

• Expression of Lie derivative in terms of parallel derivative $\mathcal{L}_{\mathbf{v}_{\varphi,t}} \mathbf{s}_{\mathrm{E},t} = \nabla_{\mathbf{v}_{\varphi,t}} \mathbf{s}_{\mathrm{E},t} - 2 \operatorname{sym} (\nabla \mathbf{v}_{\varphi,t} \circ \mathbf{s}_{\mathrm{E},t})$

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▶ Change of observer $oldsymbol{\zeta}_{\mathrm{E}} \in \mathrm{C}^1(\mathrm{E}\,;\mathrm{E})\,,$ automorphism

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- Relative motion:

• Change of observer $\boldsymbol{\zeta}_{\mathrm{E}} \in \mathrm{C}^1(\mathrm{E}\,;\mathrm{E})\,,$ automorphism $\boldsymbol{\zeta} \in \mathrm{C}^1(\mathcal{T}_{\boldsymbol{\omega}}; \mathcal{T}_{\boldsymbol{\zeta} \uparrow \boldsymbol{\omega}}), \quad \mathsf{diffeomorphism}$



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- Relative motion:

• Change of observer $\boldsymbol{\zeta}_{\mathrm{E}} \in \mathrm{C}^1(\mathrm{E}\,;\mathrm{E})\,,$ automorphism $oldsymbol{\zeta}\in\mathrm{C}^1(\mathcal{T}_{oldsymbol{arphi}}\,;\,\mathcal{T}_{oldsymbol{\mathcal{L}}\uparrowoldsymbol{arphi}})\,,$ diffeomorphism



Pushed motion:

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Time independence and Invariance of material fields

• Time independence $\mathbf{s}_{\boldsymbol{\varphi},\tau} = \boldsymbol{\varphi}_{\tau,t} \uparrow \mathbf{s}_{\boldsymbol{\varphi},t}$

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Time independence and Invariance of material fields

- $\blacktriangleright \text{ Time independence } \mathbf{s}_{\boldsymbol{\varphi},\tau} = \boldsymbol{\varphi}_{\tau,t} \uparrow \mathbf{s}_{\boldsymbol{\varphi},t}$
- Invariance $\mathbf{s}_{\zeta \uparrow \varphi}$

 $\mathsf{s}_{oldsymbol{\zeta} egin{smallmatrix} \mathsf{s}_{oldsymbol{\zeta}} = oldsymbol{\zeta} egin{smallmatrix} \mathsf{s}_{oldsymbol{arphi}} & \mathsf{s}_{oldsymbol{$

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- $\blacktriangleright \text{ Time independence } \mathbf{s}_{\boldsymbol{\varphi},\tau} = \boldsymbol{\varphi}_{\tau,t} \uparrow \mathbf{s}_{\boldsymbol{\varphi},t}$
- ► Invariance $\mathbf{s}_{\boldsymbol{\zeta} \uparrow \boldsymbol{\varphi}} = \boldsymbol{\zeta} \uparrow \mathbf{s}_{\boldsymbol{\varphi}}$
- Push of Lie time derivative to reference

$$arphi_{t,\mathrm{REF}} \!\!\downarrow \! (\mathcal{L}_{oldsymbol{arphi},t} \, \mathbf{s}_{oldsymbol{arphi}}) = \partial_{ au = t} \, arphi_{ au,\mathrm{REF}} \!\!\downarrow \! \mathbf{s}_{oldsymbol{arphi}, au}$$

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- Invariance $\mathbf{s}_{\boldsymbol{\zeta}\uparrow\boldsymbol{arphi}} = \boldsymbol{\zeta}\uparrow\mathbf{s}_{\boldsymbol{arphi}}$
- Push of Lie time derivative to reference

$$arphi_{t,\mathrm{REF}} \!\!\downarrow \! (\mathcal{L}_{oldsymbol{arphi},t} \, \mathbf{s}_{oldsymbol{arphi}}) = \partial_{ au = t} \, arphi_{ au,\mathrm{REF}} \!\!\downarrow \! \mathbf{s}_{oldsymbol{arphi}, au}$$

Lie time derivative along pushed motions

$$\mathcal{L}_{(\zeta \uparrow arphi),t}\left(\zeta \uparrow \mathbf{s}_{arphi}
ight) = \zeta_t \uparrow \mathcal{L}_{arphi,t} \, \mathbf{s}_{arphi}$$

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• Constitutive operator \mathbf{H}_{φ}

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• Constitutive operator \mathbf{H}_{φ}

A material bundle morphism whose domain and codomain are Whitney products of material tensor bundles

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Covariance Paradigm

• Constitutive operator \mathbf{H}_{φ}

A material bundle morphism whose domain and codomain are Whitney products of material tensor bundles

Constitutive time independence

$$\mathbf{H}_{oldsymbol{arphi}, au}=oldsymbol{arphi}_{ au,t}\!\!\uparrow\!\mathbf{H}_{oldsymbol{arphi},t}$$

$$(arphi_{ au,t} {\uparrow} \mathbf{H}_{arphi,t}) (arphi_{ au,t} {\uparrow} \mathbf{s}_{arphi,t}) = arphi_{ au,t} {\uparrow} (\mathbf{H}_{arphi,t} (\mathbf{s}_{arphi,t}))$$

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• Constitutive operator $\mathbf{H}_{\boldsymbol{\varphi}}$

A material bundle morphism whose domain and codomain are Whitney products of material tensor bundles

Constitutive time independence

$$\mathbf{H}_{oldsymbol{arphi}, au}=oldsymbol{arphi}_{ au,t}\!\!\uparrow\!\mathbf{H}_{oldsymbol{arphi},t}$$

$$(arphi_{ au,t}\!\!\uparrow\!\mathbf{H}_{arphi,t})(arphi_{ au,t}\!\!\uparrow\!\mathbf{s}_{arphi,t}) = arphi_{ au,t}\!\!\uparrow\!(\mathbf{H}_{arphi,t}(\mathbf{s}_{arphi,t}))$$

Constitutive invariance under relative motions

 ${f H}_{\zeta \uparrow arphi} = \zeta \uparrow {f H}_arphi$ $(\zeta \uparrow {f H}_arphi)(\zeta \uparrow {f s}_arphi) = \zeta \uparrow ({f H}_arphi({f s}_arphi))$

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Constitutive hypo-elastic law

$$\left\{egin{array}{ll} \dot{m{arepsilon}}_{m{arphi}} &= {f e}_{m{arphi}} \ {f e}_{m{arphi}} &= {f H}_{m{arphi}}^{
m HYPO}({m{\sigma}}_{m{arphi}}) \cdot \dot{m{\sigma}}_{m{arphi}}
ight.$$

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Constitutive hypo-elastic law

$$egin{aligned} \dot{oldsymbol{arepsilon}}_{oldsymbol{arphi}} &= \mathbf{e}_{oldsymbol{arphi}} \ \mathbf{e}_{oldsymbol{arphi}} &= \mathbf{H}_{oldsymbol{arphi}}^{ ext{HYPO}}(oldsymbol{\sigma}_{oldsymbol{arphi}}) \cdot \dot{oldsymbol{\sigma}}_{oldsymbol{arphi}} \end{aligned}$$

► CAUCHY integrability

 $\langle d_{\mathsf{F}} \mathbf{H}_{\varphi}^{\text{HYPO}}(\boldsymbol{\sigma}_{\varphi}) \cdot \delta \boldsymbol{\sigma}_{\varphi} \cdot \delta_{1} \boldsymbol{\sigma}_{\varphi}, \delta_{2} \boldsymbol{\sigma}_{\varphi} \rangle = \text{symmetric}$

 $\mathsf{H}_{arphi}^{\scriptscriptstyle\mathrm{HYPO}}(\pmb{\sigma}_{arphi}) = d_{F} \mathbf{\Phi}_{arphi}(\pmb{\sigma}_{arphi})$

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Constitutive hypo-elastic law

$$egin{aligned} \dot{oldsymbol{arphi}}_{oldsymbol{arphi}} &= \mathbf{e}_{oldsymbol{arphi}} \ \mathbf{e}_{oldsymbol{arphi}} &= \mathbf{H}_{oldsymbol{arphi}}^{ ext{HYPO}}(oldsymbol{\sigma}_{oldsymbol{arphi}}) \cdot \dot{oldsymbol{\sigma}}_{oldsymbol{arphi}} \end{aligned}$$

► CAUCHY integrability

$$\langle d_{\mathsf{F}} \mathsf{H}^{_{\mathrm{HYPO}}}_{\boldsymbol{\varphi}}(\boldsymbol{\sigma}_{\boldsymbol{\varphi}}) \cdot \delta \boldsymbol{\sigma}_{\boldsymbol{\varphi}} \cdot \delta_{1} \boldsymbol{\sigma}_{\boldsymbol{\varphi}}, \delta_{2} \boldsymbol{\sigma}_{\boldsymbol{\varphi}} \rangle = \mathrm{symmetric} =$$

 $\mathsf{H}_{arphi}^{\scriptscriptstyle\mathrm{HYPO}}(\sigma_{arphi}) = d_{F} \mathbf{\Phi}_{arphi}(\sigma_{arphi})$

► GREEN integrability

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Covariance Paradigm

Elasticity

 Elastic constitutive operator: hypo-elastic constitutive operator which is integrable and time independent



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Covariance Paradigm

Elasticity

- Elastic constitutive operator: hypo-elastic constitutive operator which is integrable and time independent
- Constitutive elastic law:

$$\left\{ egin{array}{ll} \dot{m{arepsilon}}_{m{arphi}} &= {f e}_{m{arphi}} \ egin{array}{ll} {m{e e}} &= d_F^2 E_{m{arphi}}^*(m{\sigma}_{m{arphi}}) \cdot \dot{m{\sigma}}_{m{arphi}} \end{array}
ight.$$

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Covariance Paradigm

Elasticity

- Elastic constitutive operator: hypo-elastic constitutive operator which is integrable and time independent
- Constitutive elastic law:

$$egin{array}{lll} \left(egin{array}{lll} \dot{m{arepsilon}}_{m{arphi}} = {f e}_{m{arphi}} \ \left({f e}_{m{arphi}} = d_F^2 E_{m{arphi}}^*(m{\sigma}_{m{arphi}}) \cdot \dot{m{\sigma}}_{m{arphi}}
ight) \end{array}
ight) \end{array}$$

pull-back to reference:

$$egin{aligned} oldsymbol{arphi}_{t, ext{REF}} & igl\downarrow \mathbf{e}_{oldsymbol{arphi},t} \, = \, d_F^2 E_{ ext{REF}}^* (oldsymbol{arphi}_{t, ext{REF}} igl\downarrow oldsymbol{\sigma}_{oldsymbol{arphi},t}) \cdot \partial_{ au=t} \, oldsymbol{arphi}_{\tau, ext{REF}} igl\downarrow oldsymbol{\sigma}_{oldsymbol{arphi}, au}) \ & = \, \partial_{ au=t} \, d_F E_{ ext{REF}}^* (oldsymbol{arphi}_{\tau, ext{REF}} igl\downarrow oldsymbol{\sigma}_{oldsymbol{arphi}, au}) \end{aligned}$$

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Covariance Paradigm
Elasticity

- Elastic constitutive operator: hypo-elastic constitutive operator which is integrable and time independent
- Constitutive elastic law:

$$egin{aligned} & egin{aligned} & egi$$

pull-back to reference:

$$egin{aligned} oldsymbol{arphi}_{t, ext{REF}} & igl\downarrow \mathbf{e}_{oldsymbol{arphi},t} &= d_F^2 E_{ ext{REF}}^* (oldsymbol{arphi}_{t, ext{REF}} igl\downarrow oldsymbol{\sigma}_{oldsymbol{arphi},t}) \cdot \partial_{ au=t} \; oldsymbol{arphi}_{ au, ext{REF}} igl\downarrow oldsymbol{\sigma}_{oldsymbol{arphi},t}) \ &= \partial_{ au=t} \; d_F E_{ ext{REF}}^* (oldsymbol{arphi}_{ au, ext{REF}} igl\downarrow oldsymbol{\sigma}_{oldsymbol{arphi},t}) \end{aligned}$$

$$egin{aligned} &arphi_{ au, ext{REF}} := arphi_{ au, ext{REF}} \circ arphi_{ au, ext{REF}} \ & \mathcal{E}_{ ext{REF}}^* := arphi_{ au, ext{REF}} {\downarrow} \mathcal{E}_{oldsymbol{arphi}, ext{t}}^* & ext{time independent} \end{aligned}$$

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Conservativeness of hyper-elasticity

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Conservativeness of hyper-elasticity

GREEN integrability of the elastic operator $\mathbf{H}_{\boldsymbol{\varphi}}$ implies conservativeness:

$$\oint_{I} \int_{\Omega_{t}} \langle \boldsymbol{\sigma}_{\boldsymbol{\varphi},t}, \mathbf{e}_{\boldsymbol{\varphi},t} \rangle \, \mathbf{m}_{\boldsymbol{\varphi},t} \, dt = 0$$

for any cycle in the stress time bundle, i.e. for any stress path $\sigma_{\varphi} \in C^{1}(I; CON(\mathbb{V}T_{\varphi}))$ such that:

$$oldsymbol{\sigma}_{oldsymbol{arphi},t_2} = oldsymbol{arphi}_{t_2,t_1} \! \! \uparrow \! oldsymbol{\sigma}_{oldsymbol{arphi},t_1} \,, \quad I = [t_1,t_2]$$

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Elasto-visco-plasticity

Constitutive law

$$\left\{ egin{array}{ll} \dot{m{\varepsilon}}_{m{arphi}} &= {m{e}}_{m{arphi}} + {m{p}}_{m{arphi}} \ {m{e}}_{m{arphi}} &= d_F^2 E_{m{arphi}}^*(\sigma_{m{arphi}}) \cdot \dot{\sigma}_{m{arphi}} \ {m{p}}_{m{arphi}} &\in \partial_F \mathcal{F}_{m{arphi}}(\sigma_{m{arphi}}) \end{array}
ight.$$

stretching additivity hyper-elastic law visco-plastic flow rule

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Covariance Paradigm

• total strain in the time interval I = [s, t]:

$$oldsymbol{arphi}_{oldsymbol{arphi},t,s} := oldsymbol{arphi}_{t,s} {\downarrow} oldsymbol{g}_{oldsymbol{arphi},t} - oldsymbol{g}_{oldsymbol{arphi},s}$$

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Example

Covariance Paradigm

• total strain in the time interval I = [s, t]:

$$arepsilon_{oldsymbol{arphi},t,s} := arphi_{t,s} {\downarrow} \mathbf{g}_{oldsymbol{arphi},t} - \mathbf{g}_{oldsymbol{arphi},s}$$

reference total strain:

$$\begin{split} \boldsymbol{\varepsilon}_{\boldsymbol{\varphi},l}^{\text{REF}} &\coloneqq \frac{1}{2} \int_{I} \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,\text{REF}} \! \downarrow \! \mathbf{g}_{\boldsymbol{\varphi},\tau} \, dt \\ &= \frac{1}{2} \boldsymbol{\varphi}_{t,\text{REF}} \! \downarrow \! \mathbf{g}_{\boldsymbol{\varphi},t} - \frac{1}{2} \boldsymbol{\varphi}_{s,\text{REF}} \! \downarrow \! \mathbf{g}_{\boldsymbol{\varphi},s} \\ &= \frac{1}{2} \boldsymbol{\varphi}_{s,\text{REF}} \! \downarrow \! (\boldsymbol{\varphi}_{t,s} \! \downarrow \! \mathbf{g}_{\boldsymbol{\varphi},t} - \mathbf{g}_{\boldsymbol{\varphi},s}) = \frac{1}{2} \boldsymbol{\varphi}_{s,\text{REF}} \! \downarrow \! \boldsymbol{\varepsilon}_{\boldsymbol{\varphi},t,s} \end{split}$$

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Covariance Paradigm

• total strain in the time interval I = [s, t]:

$$arepsilon_{oldsymbol{arphi},t,s} := arphi_{t,s} {\downarrow} \mathbf{g}_{oldsymbol{arphi},t} - \mathbf{g}_{oldsymbol{arphi},s}$$

reference total strain:

$$\begin{split} \boldsymbol{\varepsilon}_{\boldsymbol{\varphi},\boldsymbol{I}}^{\text{REF}} &:= \frac{1}{2} \int_{\boldsymbol{I}} \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,\text{REF}} \!\!\downarrow \! \mathbf{g}_{\boldsymbol{\varphi},\tau} \, dt \\ &= \frac{1}{2} \boldsymbol{\varphi}_{t,\text{REF}} \!\!\downarrow \! \mathbf{g}_{\boldsymbol{\varphi},t} - \frac{1}{2} \boldsymbol{\varphi}_{s,\text{REF}} \!\!\downarrow \! \mathbf{g}_{\boldsymbol{\varphi},s} \\ &= \frac{1}{2} \boldsymbol{\varphi}_{s,\text{REF}} \!\!\downarrow \! (\boldsymbol{\varphi}_{t,s} \!\!\downarrow \! \mathbf{g}_{\boldsymbol{\varphi},t} - \mathbf{g}_{\boldsymbol{\varphi},s}) = \frac{1}{2} \boldsymbol{\varphi}_{s,\text{REF}} \!\!\downarrow \! \boldsymbol{\varepsilon}_{\boldsymbol{\varphi},t,s} \end{split}$$

reference elastic and visco-plastic strain:

$$\mathbf{e}_{\boldsymbol{\varphi},l}^{\text{REF}} := \int_{l} \boldsymbol{\varphi}_{t,\text{REF}} \downarrow \mathbf{e}_{\boldsymbol{\varphi},t} \, dt \,, \qquad \mathbf{p}_{\boldsymbol{\varphi},l}^{\text{REF}} := \int_{l} \boldsymbol{\varphi}_{t,\text{REF}} \downarrow \mathbf{p}_{\boldsymbol{\varphi},t} \, dt$$

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• total strain in the time interval I = [s, t]:

$$arepsilon_{oldsymbol{arphi},t,s} := arphi_{t,s} {\downarrow} \mathbf{g}_{oldsymbol{arphi},t} - \mathbf{g}_{oldsymbol{arphi},s}$$

reference total strain:

$$\begin{split} \boldsymbol{\varepsilon}_{\boldsymbol{\varphi},\boldsymbol{I}}^{\text{REF}} &:= \frac{1}{2} \int_{\boldsymbol{I}} \partial_{\tau=t} \, \boldsymbol{\varphi}_{\tau,\text{REF}} \!\!\downarrow \! \mathbf{g}_{\boldsymbol{\varphi},\tau} \, dt \\ &= \frac{1}{2} \boldsymbol{\varphi}_{t,\text{REF}} \!\!\downarrow \! \mathbf{g}_{\boldsymbol{\varphi},t} - \frac{1}{2} \boldsymbol{\varphi}_{s,\text{REF}} \!\!\downarrow \! \mathbf{g}_{\boldsymbol{\varphi},s} \\ &= \frac{1}{2} \boldsymbol{\varphi}_{s,\text{REF}} \!\!\downarrow \! (\boldsymbol{\varphi}_{t,s} \!\!\downarrow \! \mathbf{g}_{\boldsymbol{\varphi},t} - \mathbf{g}_{\boldsymbol{\varphi},s}) = \frac{1}{2} \boldsymbol{\varphi}_{s,\text{REF}} \!\!\downarrow \! \boldsymbol{\varepsilon}_{\boldsymbol{\varphi},t,s} \end{split}$$

reference elastic and visco-plastic strain:

$$\mathbf{e}_{\varphi,I}^{\text{REF}} := \int_{I} \varphi_{t,\text{REF}} \downarrow \mathbf{e}_{\varphi,t} \, dt \,, \qquad \mathbf{p}_{\varphi,I}^{\text{REF}} := \int_{I} \varphi_{t,\text{REF}} \downarrow \mathbf{p}_{\varphi,t} \, dt$$

additivity of reference strains:

$$\boldsymbol{\varepsilon}_{\boldsymbol{arphi},l}^{\scriptscriptstyle ext{REF}} = \mathbf{e}_{\boldsymbol{arphi},l}^{\scriptscriptstyle ext{REF}} + \mathbf{p}_{\boldsymbol{arphi},l}^{\scriptscriptstyle ext{REF}}$$

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Any constitutive law must conform to the principle of MFI which requires that material fields, fulfilling the law, will still fulfill it when evaluated by another Euclid observer The G-Factor Impact in NLCM

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Covariance Paradigm

Principle of MFI

Any constitutive law must conform to the principle of MFI which requires that material fields, fulfilling the law, will still fulfill it when evaluated by another Euclid observer

$$arepsilon_{arphi} = \mathsf{H}_{arphi}(\mathsf{s}_{arphi}) \quad \Longleftrightarrow \quad arepsilon_{\zeta^{\mathrm{iso}} \uparrow arphi} = \mathsf{H}_{\zeta^{\mathrm{iso}} \uparrow arphi}(\mathsf{s}_{\zeta^{\mathrm{iso}} \uparrow arphi}),$$

► for any isometric relative motion $\zeta^{iso} \in C^1(\mathcal{I}_{\varphi}; \mathcal{I}_{\zeta^{iso} \uparrow \varphi})$ induced by a change of Euclid observer $\zeta_{\mathsf{E}}^{iso} \in C^1(\mathsf{E}; \mathsf{E})$.

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Sufficient conditions

Material fields must be frame invariant

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Sufficient conditions

- Material fields must be frame invariant
- Constitutive operators must be frame invariant

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Covariance Paradigm

► Frame invariance of the hypo-elastic operator

$$\mathsf{H}^{\scriptscriptstyle\mathrm{HYPO}}_{\boldsymbol{\zeta}^{\scriptscriptstyle\mathrm{ISO}}\uparrowoldsymbol{arphi}}=\boldsymbol{\zeta}^{\scriptscriptstyle\mathrm{ISO}}\!\uparrow\!\mathsf{H}^{\scriptscriptstyle\mathrm{HYPO}}_{oldsymbol{arphi}}$$

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Pushed operator

$$(\boldsymbol{\zeta}^{ ext{ISO}}\!\!\uparrow\!\boldsymbol{\mathsf{H}}_{arphi}^{ ext{HYPO}})(\boldsymbol{\zeta}^{ ext{ISO}}\!\!\uparrow\!\boldsymbol{\sigma}_{arphi})\cdot\boldsymbol{\zeta}^{ ext{ISO}}\!\!\uparrow\!\!\dot{\boldsymbol{\sigma}}_{arphi}=\boldsymbol{\zeta}^{ ext{ISO}}\!\!\uparrow\!(\boldsymbol{\mathsf{H}}_{arphi}^{ ext{HYPO}}(\boldsymbol{\sigma}_{arphi})\cdot\dot{\boldsymbol{\sigma}}_{arphi})$$

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Examples:

the simplest hypo-elastic operator is frame invariant:

$$\mathbf{H}_{\boldsymbol{\varphi},t}^{\text{HYPO}}(\mathbf{T}_{\boldsymbol{\varphi},t}) := \frac{1}{2\,\mu}\,\mathbb{I}_{\boldsymbol{\varphi},t} - \frac{\nu}{E}\,\mathbf{I}_{\boldsymbol{\varphi},t}\otimes\mathbf{I}_{\boldsymbol{\varphi},t}\,,$$

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$$(\boldsymbol{\zeta}^{ ext{ISO}}\!\!\uparrow\!\mathsf{H}^{ ext{HYPO}}_{oldsymbol{arphi}})(\boldsymbol{\zeta}^{ ext{ISO}}\!\!\uparrow\!\boldsymbol{\sigma}_{oldsymbol{arphi}})\cdot\boldsymbol{\zeta}^{ ext{ISO}}\!\!\uparrow\!\!\dot{oldsymbol{\sigma}}_{oldsymbol{arphi}}=\boldsymbol{\zeta}^{ ext{ISO}}\!\!\uparrow\!(\mathsf{H}^{ ext{HYPO}}_{oldsymbol{arphi}}(oldsymbol{\sigma}_{oldsymbol{arphi}})\cdot\dot{oldsymbol{\sigma}}_{oldsymbol{arphi}})$$

Examples:

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the visco-plastic flow rule is frame invariant

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▶ theoretical: spatial, material and material based spatial fields

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- ▶ theoretical: spatial, material and material based spatial fields
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Covariance Paradigm

- ▶ theoretical: spatial, material and material based spatial fields
- theoretical: covariance paradigm
- theoretical: stretching and stressing are Lie time-derivatives



- ▶ theoretical: spatial, material and material based spatial fields
- theoretical: covariance paradigm
- theoretical: stretching and stressing are Lie time-derivatives
- theoretical: covariant formulation of constitutive laws

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- theoretical: covariance paradigm
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- ▶ theoretical: covariant formulation of Material Frame Indifference

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- theoretical: covariant theory of elasto-visco-plasticity
- computational: integrability of simplest hypo-elasticity

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- computational: integrability of simplest hypo-elasticity
- computational: finite elastic (anelastic) strains are time integrals of strain rates pull-back to a reference placement

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- theoretical: covariant theory of elasto-visco-plasticity
- computational: integrability of simplest hypo-elasticity
- computational: finite elastic (anelastic) strains are time integrals of strain rates pull-back to a reference placement
- computational: constitutive relations in the nonlinear range are governed by rate laws which may be got from linearized ones by substituting Lie time-derivatives to partial time derivatives

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Consequences

treatments of constitutive behaviors in the nonlinear range should be revised and reformulated

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Consequences

- treatments of constitutive behaviors in the nonlinear range should be revised and reformulated
- algorithms for numerical computations must be modified to comply with the covariant theory; multiplicative decomposition of the deformation gradient is geometrically inconsistent

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