

## DINAMICA

- equazioni del moto del manipolatore in funzione delle forze e momenti agenti su di esso

### Formulazione di Lagrange

### Proprietà notevoli del modello dinamico

### Identificazione dei parametri dinamici

### Formulazione di Newton–Eulero

### Dinamica diretta e dinamica inversa

### Modello dinamico nello spazio operativo

### Ellissoide di manipolabilità dinamica

### Scalatura dinamica di traiettorie

# FORMULAZIONE DI LAGRANGE

- Lagrangiana  $\equiv$  (energia cinetica) – (energia potenziale)

$$\mathcal{L} = \mathcal{T} - \mathcal{U}$$

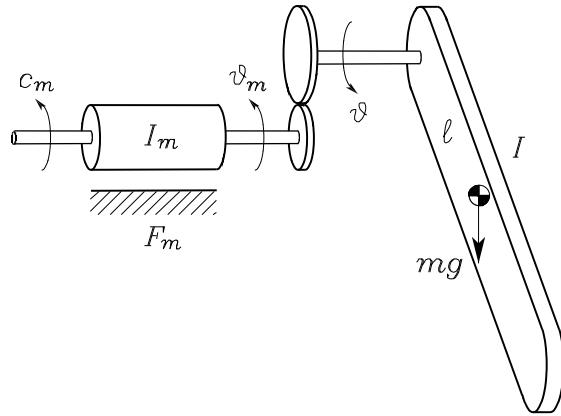
- Equazioni di Lagrange

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\lambda}_i} - \frac{\partial \mathcal{L}}{\partial \lambda_i} = \xi_i \quad i = 1, \dots, n$$

\* *coordinate generalizzate*

$$\begin{bmatrix} \lambda_1 \\ \vdots \\ \lambda_n \end{bmatrix} = \mathbf{q}$$

- Esempio



- ★ energia cinetica

$$\mathcal{T} = \frac{1}{2}I\dot{\vartheta}^2 + \frac{1}{2}I_m k_r^2 \dot{\vartheta}^2$$

- ★ energia potenziale

$$\mathcal{U} = mg\ell(1 - \cos \vartheta)$$

- ★ lagrangiana

$$\mathcal{L} = \frac{1}{2}I\dot{\vartheta}^2 + \frac{1}{2}I_m k_r^2 \dot{\vartheta}^2 - mg\ell(1 - \cos \vartheta)$$

- ★ equazione del moto

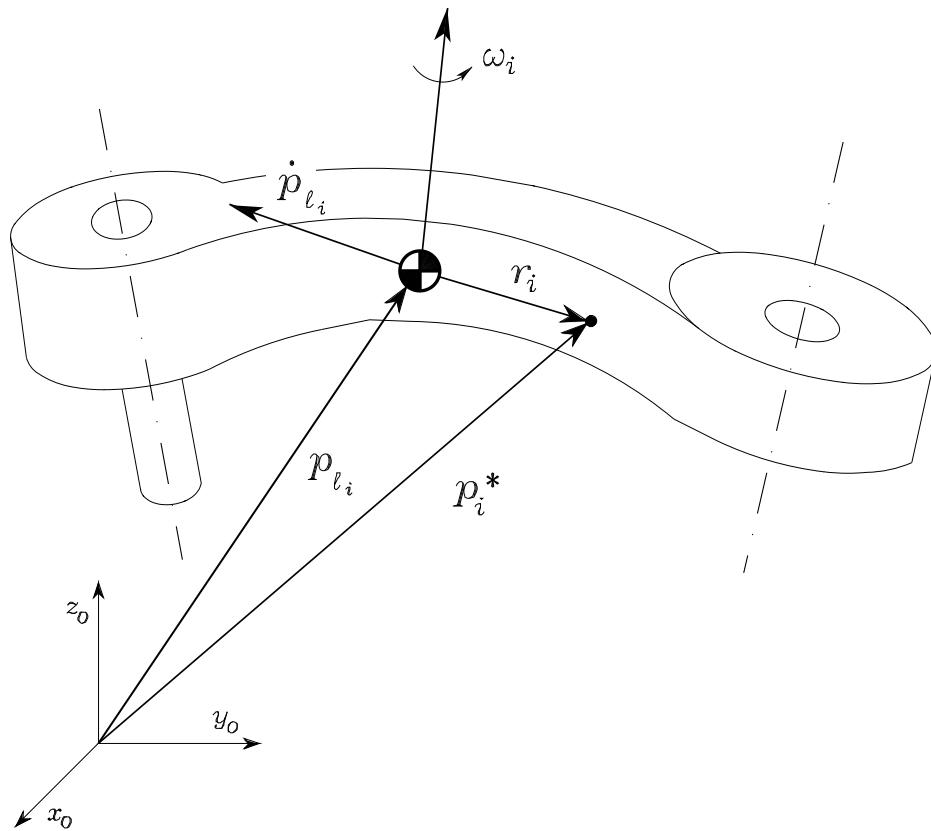
$$(I + I_m k_r^2)\ddot{\vartheta} + mg\ell \sin \vartheta = \xi \\ = \tau - F\dot{\vartheta} - F_m k_r \dot{\vartheta}$$

## Energia cinetica

- Contributo bracci + attuatori

$$\mathcal{T} = \sum_{i=1}^n (\mathcal{T}_{\ell_i} + \mathcal{T}_{m_i})$$

- Braccio  $i$



$$\mathcal{T}_{\ell_i} = \frac{1}{2} \int_{V_{\ell_i}} \dot{\mathbf{p}}_i^{*T} \dot{\mathbf{p}}_i^* \rho dV$$

★ baricentro

$$\mathbf{p}_{\ell_i} = \frac{1}{m_{\ell_i}} \int_{V_{\ell_i}} \mathbf{p}_i^* \rho dV$$

★ velocità lineare della particella elementare

$$\dot{\mathbf{p}}_i^* = \dot{\mathbf{p}}_{\ell_i} + \boldsymbol{\omega}_i \times \mathbf{r}_i = \dot{\mathbf{p}}_{\ell_i} + \mathbf{S}(\boldsymbol{\omega}_i) \mathbf{r}_i$$

- Traslazionale

$$\frac{1}{2} \int_{V_{\ell_i}} \dot{\mathbf{p}}_{\ell_i}^T \dot{\mathbf{p}}_{\ell_i} \rho dV = \frac{1}{2} m_{\ell_i} \dot{\mathbf{p}}_{\ell_i}^T \dot{\mathbf{p}}_{\ell_i}$$

- Mutuo

$$\begin{aligned} 2 \left( \frac{1}{2} \int_{V_{\ell_i}} \dot{\mathbf{p}}_{\ell_i}^T \mathbf{S}(\boldsymbol{\omega}_i) \mathbf{r}_i \rho dV \right) &= \\ 2 \left( \frac{1}{2} \dot{\mathbf{p}}_{\ell_i}^T \mathbf{S}(\boldsymbol{\omega}_i) \int_{V_{\ell_i}} (\mathbf{p}_i^* - \mathbf{p}_{\ell_i}) \rho dV \right) &= 0 \end{aligned}$$

- Rotazionale

$$\begin{aligned} \frac{1}{2} \int_{V_{\ell_i}} \mathbf{r}_i^T \mathbf{S}^T(\boldsymbol{\omega}_i) \mathbf{S}(\boldsymbol{\omega}_i) \mathbf{r}_i \rho dV &= \frac{1}{2} \boldsymbol{\omega}_i^T \left( \int_{V_{\ell_i}} \mathbf{S}^T(\mathbf{r}_i) \mathbf{S}(\mathbf{r}_i) \rho dV \right) \boldsymbol{\omega}_i \\ &= \frac{1}{2} \boldsymbol{\omega}_i^T \mathbf{I}_{\ell_i} \boldsymbol{\omega}_i \end{aligned}$$

★ *tensore di inerzia*

$$\begin{aligned} \mathbf{I}_{\ell_i} &= \begin{bmatrix} \int (r_{iy}^2 + r_{iz}^2) \rho dV & -\int r_{ix} r_{iy} \rho dV & -\int r_{ix} r_{iz} \rho dV \\ * & \int (r_{ix}^2 + r_{iz}^2) \rho dV & -\int r_{iy} r_{iz} \rho dV \\ * & * & \int (r_{ix}^2 + r_{iy}^2) \rho dV \end{bmatrix} \\ &= \begin{bmatrix} I_{\ell_i xx} & -I_{\ell_i xy} & -I_{\ell_i xz} \\ * & I_{\ell_i yy} & -I_{\ell_i yz} \\ * & * & I_{\ell_i zz} \end{bmatrix} \end{aligned}$$

- Energia cinetica del braccio  $i$

$$\begin{aligned}\mathcal{T}_{\ell_i} &= \frac{1}{2}m_{\ell_i}\dot{\mathbf{p}}_{\ell_i}^T\dot{\mathbf{p}}_{\ell_i} + \frac{1}{2}\boldsymbol{\omega}_i^T\mathbf{R}_i\mathbf{I}_{\ell_i}^i\mathbf{R}_i^T\boldsymbol{\omega}_i \\ &= \frac{1}{2}m_{\ell_i}\dot{\mathbf{q}}^T\mathbf{J}_P^{(\ell_i)T}\mathbf{J}_P^{(\ell_i)}\dot{\mathbf{q}} + \frac{1}{2}\dot{\mathbf{q}}^T\mathbf{J}_O^{(\ell_i)T}\mathbf{R}_i\mathbf{I}_{\ell_i}^i\mathbf{R}_i^T\mathbf{J}_O^{(\ell_i)}\dot{\mathbf{q}}\end{aligned}$$

★ velocità lineare

$$\dot{\mathbf{p}}_{\ell_i} = \mathbf{J}_{P1}^{(\ell_i)}\dot{q}_1 + \dots + \mathbf{J}_{Pi}^{(\ell_i)}\dot{q}_i = \mathbf{J}_P^{(\ell_i)}\dot{\mathbf{q}}$$

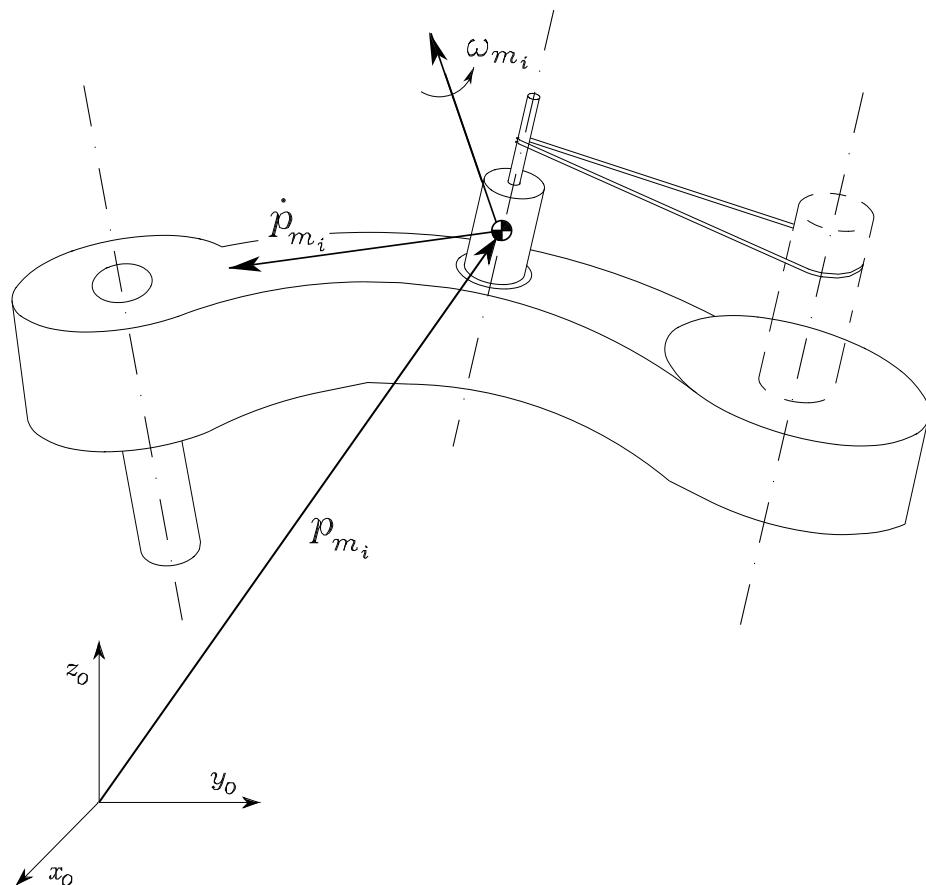
$$\mathbf{J}_{Pj}^{(\ell_i)} = \begin{cases} \mathbf{z}_{j-1} & \text{per un giunto prismatico} \\ \mathbf{z}_{j-1} \times (\mathbf{p}_{\ell_i} - \mathbf{p}_{j-1}) & \text{per un giunto rotoidale} \end{cases}$$

★ velocità angolare

$$\boldsymbol{\omega}_i = \mathbf{J}_{O1}^{(\ell_i)}\dot{q}_1 + \dots + \mathbf{J}_{Oi}^{(\ell_i)}\dot{q}_i = \mathbf{J}_O^{(\ell_i)}\dot{\mathbf{q}}$$

$$\mathbf{J}_{Oj}^{(\ell_i)} = \begin{cases} \mathbf{0} & \text{per un giunto prismatico} \\ \mathbf{z}_{j-1} & \text{per un giunto rotoidale} \end{cases}$$

- Motore  $i$  (elettrico rotante)



★ trasmissione rigida

$$\boldsymbol{\omega}_{m_i} = \boldsymbol{\omega}_{i-1} + k_{ri} \dot{q}_i \boldsymbol{z}_{m_i}$$

- Energia cinetica del rotore  $i$

$$\begin{aligned}\mathcal{T}_{m_i} &= \frac{1}{2} m_{m_i} \dot{\mathbf{p}}_{m_i}^T \dot{\mathbf{p}}_{m_i} + \frac{1}{2} \boldsymbol{\omega}_{m_i}^T \mathbf{I}_{m_i} \boldsymbol{\omega}_{m_i} \\ &= \frac{1}{2} m_{m_i} \dot{\mathbf{q}}^T \mathbf{J}_P^{(m_i)T} \mathbf{J}_P^{(m_i)} \dot{\mathbf{q}} + \frac{1}{2} \dot{\mathbf{q}}^T \mathbf{J}_O^{(m_i)T} \mathbf{R}_{m_i} \mathbf{I}_{m_i}^{m_i} \mathbf{R}_{m_i}^T \mathbf{J}_O^{(m_i)} \dot{\mathbf{q}}\end{aligned}$$

★ velocità lineare

$$\mathbf{J}_P^{(m_i)} = [\mathbf{j}_{P1}^{(m_i)} \quad \dots \quad \mathbf{j}_{P,i-1}^{(m_i)} \quad \mathbf{0} \quad \dots \quad \mathbf{0}]$$

$$\mathbf{j}_{Pj}^{(m_i)} = \begin{cases} \mathbf{z}_{j-1} & \text{per un giunto prismatico} \\ \mathbf{z}_{j-1} \times (\mathbf{p}_{m_i} - \mathbf{p}_{j-1}) & \text{per un giunto rotoidale} \end{cases}$$

★ velocità angolare

$$\mathbf{J}_O^{(m_i)} = [\mathbf{j}_{O1}^{(m_i)} \quad \dots \quad \mathbf{j}_{O,i-1}^{(m_i)} \quad \mathbf{j}_{Oi}^{(m_i)} \quad \mathbf{0} \quad \dots \quad \mathbf{0}]$$

$$\mathbf{j}_{Oj}^{(m_i)} = \begin{cases} \mathbf{j}_{Oj}^{(\ell_i)} & j = 1, \dots, i-1 \\ k_{ri} \mathbf{z}_{m_i} & j = i \end{cases}$$

- Energia cinetica totale

$$\mathcal{T} = \frac{1}{2} \sum_{i=1}^n \sum_{j=1}^n b_{ij}(\mathbf{q}) \dot{q}_i \dot{q}_j = \frac{1}{2} \dot{\mathbf{q}}^T \mathbf{B}(\mathbf{q}) \dot{\mathbf{q}}$$

- *Matrice di inerzia*

$$\begin{aligned} \mathbf{B}(\mathbf{q}) = & \sum_{i=1}^n \left( m_{\ell_i} \mathbf{J}_P^{(\ell_i)T} \mathbf{J}_P^{(\ell_i)} + \mathbf{J}_O^{(\ell_i)T} \mathbf{R}_i \mathbf{I}_{\ell_i}^i \mathbf{R}_i^T \mathbf{J}_O^{(\ell_i)} \right. \\ & \left. + m_{m_i} \mathbf{J}_P^{(m_i)T} \mathbf{J}_P^{(m_i)} + \mathbf{J}_O^{(m_i)T} \mathbf{R}_{m_i} \mathbf{I}_{m_i}^{m_i} \mathbf{R}_{m_i}^T \mathbf{J}_O^{(m_i)} \right) \end{aligned}$$

- ★ simmetrica
- ★ definita positiva
- ★ dipendente dalla configurazione (in generale)

## Energia potenziale

$$\mathcal{U} = \sum_{i=1}^n (\mathcal{U}_{\ell_i} + \mathcal{U}_{m_i})$$

- Braccio  $i$

$$\mathcal{U}_{\ell_i} = - \int_{V_{\ell_i}} \mathbf{g}_0^T \mathbf{p}_i^* \rho dV = -m_{\ell_i} \mathbf{g}_0^T \mathbf{p}_{\ell_i}$$

- Rotore  $i$

$$\mathcal{U}_{m_i} = -m_{m_i} \mathbf{g}_0^T \mathbf{p}_{m_i}$$

- Energia potenziale totale

$$\mathcal{U} = - \sum_{i=1}^n (m_{\ell_i} \mathbf{g}_0^T \mathbf{p}_{\ell_i} + m_{m_i} \mathbf{g}_0^T \mathbf{p}_{m_i})$$

## Equazioni del moto

$$\mathcal{L}(\mathbf{q}, \dot{\mathbf{q}}) = \mathcal{T}(\mathbf{q}, \dot{\mathbf{q}}) - \mathcal{U}(\mathbf{q})$$

$$= \frac{1}{2} \sum_{i=1}^n \sum_{j=1}^n b_{ij}(\mathbf{q}) \dot{q}_i \dot{q}_j + \sum_{i=1}^n (m_{\ell_i} \mathbf{g}_0^T \mathbf{p}_{\ell_i}(\mathbf{q}) + m_{m_i} \mathbf{g}_0^T \mathbf{p}_{m_i}(\mathbf{q}))$$

- Equazioni di Lagrange

$$\mathbf{B}(\mathbf{q}) \ddot{\mathbf{q}} + \mathbf{n}(\mathbf{q}, \dot{\mathbf{q}}) = \boldsymbol{\xi}$$

$$\mathbf{n}(\mathbf{q}, \dot{\mathbf{q}}) = \dot{\mathbf{B}}(\mathbf{q}) \dot{\mathbf{q}} - \frac{1}{2} \left( \frac{\partial}{\partial \mathbf{q}} (\dot{\mathbf{q}}^T \mathbf{B}(\mathbf{q}) \dot{\mathbf{q}}) \right)^T + \left( \frac{\partial \mathcal{U}(\mathbf{q})}{\partial \mathbf{q}} \right)^T$$

$$\frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{q}_i} \right) = \frac{d}{dt} \left( \frac{\partial \mathcal{T}}{\partial \dot{q}_i} \right) = \sum_{j=1}^n b_{ij}(\mathbf{q}) \ddot{q}_j + \sum_{j=1}^n \frac{db_{ij}(\mathbf{q})}{dt} \dot{q}_j$$

$$= \sum_{j=1}^n b_{ij}(\mathbf{q}) \ddot{q}_j + \sum_{j=1}^n \sum_{k=1}^n \frac{\partial b_{ij}(\mathbf{q})}{\partial q_k} \dot{q}_k \dot{q}_j$$

$$\frac{\partial \mathcal{T}}{\partial q_i} = \frac{1}{2} \sum_{j=1}^n \sum_{k=1}^n \frac{\partial b_{jk}(\mathbf{q})}{\partial q_i} \dot{q}_k \dot{q}_j$$

$$\begin{aligned} \frac{\partial \mathcal{U}}{\partial q_i} &= - \sum_{j=1}^n \left( m_{\ell_j} \mathbf{g}_0^T \frac{\partial \mathbf{p}_{\ell_j}}{\partial q_i} + m_{m_j} \mathbf{g}_0^T \frac{\partial \mathbf{p}_{m_j}}{\partial q_i} \right) \\ &= - \sum_{j=1}^n \left( m_{\ell_j} \mathbf{g}_0^T \boldsymbol{\jmath}_{Pi}^{(\ell_j)}(\mathbf{q}) + m_{m_j} \mathbf{g}_0^T \boldsymbol{\jmath}_{Pi}^{(m_j)}(\mathbf{q}) \right) \\ &= g_i(\mathbf{q}) \end{aligned}$$

- Equazioni del moto

$$\sum_{j=1}^n b_{ij}(\mathbf{q}) \ddot{q}_j + \sum_{j=1}^n \sum_{k=1}^n h_{ijk}(\mathbf{q}) \dot{q}_k \dot{q}_j + g_i(\mathbf{q}) = \xi_i \quad i = 1, \dots, n$$

ove

$$h_{ijk} = \frac{\partial b_{ij}}{\partial q_k} - \frac{1}{2} \frac{\partial b_{jk}}{\partial q_i}$$

- ★ *termini in accelerazione*

il coefficiente  $b_{ii}$  rappresenta il momento di inerzia visto all'asse del giunto  $i$ , nella configurazione corrente del manipolatore, quando gli altri giunti sono bloccati

il coefficiente  $b_{ij}$  tiene conto dell'effetto dell'accelerazione del giunto  $j$  sul giunto  $i$

- ★ *termini quadratici in velocità*

il termine  $h_{ijj} \dot{q}_j^2$  rappresenta l'effetto *centrifugo* indotto al giunto  $i$  dalla velocità del giunto  $j$ ; si noti che  $h_{iii} = 0$ , poiché  $\partial b_{ii}/\partial q_i = 0$

il termine  $h_{ijk} \dot{q}_j \dot{q}_k$  rappresenta l'effetto di *Coriolis* indotto al giunto  $i$  dalle velocità dei giunti  $j$  e  $k$

- ★ *termini dipendenti solo dalla configurazione*

il termine  $g_i$  rappresenta la coppia generata all'asse del giunto  $i$  nella configurazione corrente del manipolatore per effetto della gravità

- Forze non conservative
  - ★ coppie di attuazione  $\tau$
  - ★ coppie di attrito viscoso  $-F_v \dot{q}$
  - ★ coppie di attrito statico  $-f_s(q, \dot{q}) \approx -F_s \text{sgn}(\dot{q})$
  - ★ coppie di bilanciamento di forze di contatto  $-J^T(q)h$
- Modello dinamico nello spazio dei giunti

$$\mathbf{B}(q)\ddot{q} + \mathbf{C}(q, \dot{q})\dot{q} + \mathbf{F}_v\dot{q} + \mathbf{f}_s(\dot{q}) + \mathbf{g}(q) = \boldsymbol{\tau} - \mathbf{J}^T(q)\mathbf{h}_e$$

$$\sum_{j=1}^n c_{ij} \dot{q}_j = \sum_{j=1}^n \sum_{k=1}^n h_{ijk} \dot{q}_k \dot{q}_j$$

# PROPRIETÀ NOTEVOLI DEL MODELLO DINAMICO

## Anti-simmetria della matrice $\dot{B} - 2C$

$$\begin{aligned} \sum_{j=1}^n c_{ij} \dot{q}_j &= \sum_{j=1}^n \sum_{k=1}^n h_{ijk} \dot{q}_k \dot{q}_j = \sum_{j=1}^n \sum_{k=1}^n \left( \frac{\partial b_{ij}}{\partial q_k} - \frac{1}{2} \frac{\partial b_{jk}}{\partial q_i} \right) \dot{q}_k \dot{q}_j \\ &= \frac{1}{2} \sum_{j=1}^n \sum_{k=1}^n \frac{\partial b_{ij}}{\partial q_k} \dot{q}_k \dot{q}_j + \frac{1}{2} \sum_{j=1}^n \sum_{k=1}^n \left( \frac{\partial b_{ik}}{\partial q_j} - \frac{\partial b_{jk}}{\partial q_i} \right) \dot{q}_k \dot{q}_j \end{aligned}$$

- Elementi di  $C$

$$c_{ij} = \sum_{k=1}^n c_{ijk} \dot{q}_k$$

★ simboli di Christoffel del primo tipo

$$c_{ijk} = \frac{1}{2} \left( \frac{\partial b_{ij}}{\partial q_k} + \frac{\partial b_{ik}}{\partial q_j} - \frac{\partial b_{jk}}{\partial q_i} \right)$$

- Proprietà notevole

$$\mathbf{N}(\mathbf{q}, \dot{\mathbf{q}}) = \dot{\mathbf{B}}(\mathbf{q}) - 2\mathbf{C}(\mathbf{q}, \dot{\mathbf{q}}) = -\mathbf{N}^T(\mathbf{q}, \dot{\mathbf{q}})$$

- Elementi di  $\mathbf{C}$

$$\begin{aligned} c_{ij} &= \frac{1}{2} \sum_{k=1}^n \frac{\partial b_{ij}}{\partial q_k} \dot{q}_k + \frac{1}{2} \sum_{k=1}^n \left( \frac{\partial b_{ik}}{\partial q_j} - \frac{\partial b_{jk}}{\partial q_i} \right) \dot{q}_k \\ &= \frac{1}{2} \dot{b}_{ij} + \frac{1}{2} \sum_{k=1}^n \left( \frac{\partial b_{ik}}{\partial q_j} - \frac{\partial b_{jk}}{\partial q_i} \right) \dot{q}_k \end{aligned}$$

- Elementi di  $\mathbf{N}$

$$n_{ij} = \dot{b}_{ij} - 2c_{ij} = \sum_{k=1}^n \left( \frac{\partial b_{jk}}{\partial q_i} - \frac{\partial b_{ik}}{\partial q_j} \right) \dot{q}_k = -n_{ji}$$

↓

$$\mathbf{w}^T \mathbf{N}(\mathbf{q}, \dot{\mathbf{q}}) \mathbf{w} = 0 \quad \forall \mathbf{w}$$

\* se  $\mathbf{w} = \dot{\mathbf{q}}$ :

$$\dot{\mathbf{q}}^T \mathbf{N}(\mathbf{q}, \dot{\mathbf{q}}) \dot{\mathbf{q}} = \mathbf{0} \quad \forall \mathbf{C}$$

- Principio di conservazione dell'energia (*Hamilton*)

$$\frac{1}{2} \frac{d}{dt} (\dot{\mathbf{q}}^T \mathbf{B}(\mathbf{q}) \dot{\mathbf{q}}) = \dot{\mathbf{q}}^T (\boldsymbol{\tau} - \mathbf{F}_v \dot{\mathbf{q}} - \mathbf{f}_s(\dot{\mathbf{q}}) - \mathbf{g}(\mathbf{q}) - \mathbf{J}^T(\mathbf{q}) \mathbf{h}_e)$$

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} (\dot{\mathbf{q}}^T \mathbf{B}(\mathbf{q}) \dot{\mathbf{q}}) &= \frac{1}{2} \dot{\mathbf{q}}^T \dot{\mathbf{B}}(\mathbf{q}) \dot{\mathbf{q}} + \dot{\mathbf{q}}^T \mathbf{B}(\mathbf{q}) \ddot{\mathbf{q}} \\ &= \frac{1}{2} \dot{\mathbf{q}}^T (\dot{\mathbf{B}}(\mathbf{q}) - 2\mathbf{C}(\mathbf{q}, \dot{\mathbf{q}})) \dot{\mathbf{q}} \\ &\quad + \dot{\mathbf{q}}^T (\boldsymbol{\tau} - \mathbf{F}_v \dot{\mathbf{q}} - \mathbf{f}_s(\dot{\mathbf{q}}) - \mathbf{g}(\mathbf{q}) - \mathbf{J}^T(\mathbf{q}) \mathbf{h}_e) \end{aligned}$$

## Linearità nei parametri dinamici

- Insieme braccio  $i$  e rotore  $i + 1$

$$\mathcal{T}_i = \mathcal{T}_{\ell_i} + \mathcal{T}_{m_{i+1}}$$

$$\mathcal{T}_{\ell_i} = \frac{1}{2} m_{\ell_i} \dot{\mathbf{p}}_{\ell_i}^T \dot{\mathbf{p}}_{\ell_i} + \frac{1}{2} \boldsymbol{\omega}_i^T \mathbf{I}_{\ell_i} \boldsymbol{\omega}_i$$

$$\mathcal{T}_{m_{i+1}} = \frac{1}{2} m_{m_{i+1}} \dot{\mathbf{p}}_{m_{i+1}}^T \dot{\mathbf{p}}_{m_{i+1}} + \frac{1}{2} \boldsymbol{\omega}_{m_{i+1}}^T \mathbf{I}_{m_{i+1}} \boldsymbol{\omega}_{m_{i+1}}$$

- ★ con riferimento al baricentro  $\mathbf{p}_{C_i}$  dell'insieme braccio–rotore

$$\dot{\mathbf{p}}_{\ell_i} = \dot{\mathbf{p}}_{C_i} + \boldsymbol{\omega}_i \times \mathbf{r}_{C_i, \ell_i}$$

$$\dot{\mathbf{p}}_{m_{i+1}} = \dot{\mathbf{p}}_{C_i} + \boldsymbol{\omega}_i \times \mathbf{r}_{C_i, m_{i+1}}$$

$$\mathbf{r}_{C_i, \ell_i} = \mathbf{p}_{\ell_i} - \mathbf{p}_{C_i}$$

$$\mathbf{r}_{C_i, m_{i+1}} = \mathbf{p}_{m_{i+1}} - \mathbf{p}_{C_i}$$

- Braccio

$$\begin{aligned}
 \mathcal{T}_{\ell_i} &= \frac{1}{2} m_{\ell_i} \dot{\mathbf{p}}_{C_i}^T \dot{\mathbf{p}}_{C_i} + \dot{\mathbf{p}}_{C_i}^T \mathbf{S}(\boldsymbol{\omega}_i) m_{\ell_i} \mathbf{r}_{C_i, \ell_i} \\
 &\quad + \frac{1}{2} m_{\ell_i} \boldsymbol{\omega}_i^T \mathbf{S}^T(\mathbf{r}_{C_i, \ell_i}) \mathbf{S}(\mathbf{r}_{C_i, \ell_i}) \boldsymbol{\omega}_i + \frac{1}{2} \boldsymbol{\omega}_i^T \mathbf{I}_{\ell_i} \boldsymbol{\omega}_i \\
 &= \frac{1}{2} m_{\ell_i} \dot{\mathbf{p}}_{C_i}^T \dot{\mathbf{p}}_{C_i} + \dot{\mathbf{p}}_{C_i}^T \mathbf{S}(\boldsymbol{\omega}_i) m_{\ell_i} \mathbf{r}_{C_i, \ell_i} + \frac{1}{2} \boldsymbol{\omega}_i^T \bar{\mathbf{I}}_{\ell_i} \boldsymbol{\omega}_i
 \end{aligned}$$

★ per il teorema di Steiner

$$\bar{\mathbf{I}}_{\ell_i} = \mathbf{I}_{\ell_i} + m_{\ell_i} \mathbf{S}^T(\mathbf{r}_{C_i, \ell_i}) \mathbf{S}(\mathbf{r}_{C_i, \ell_i})$$

- Rotore

$$\begin{aligned}
 \mathcal{T}_{m_{i+1}} &= \frac{1}{2} m_{m_{i+1}} \dot{\mathbf{p}}_{C_i}^T \dot{\mathbf{p}}_{C_i} + \dot{\mathbf{p}}_{C_i}^T \mathbf{S}(\boldsymbol{\omega}_i) m_{m_{i+1}} \mathbf{r}_{C_i, m_{i+1}} + \frac{1}{2} \boldsymbol{\omega}_i^T \bar{\mathbf{I}}_{m_{i+1}} \boldsymbol{\omega}_i \\
 &\quad + k_{r,i+1} \dot{q}_{i+1} \mathbf{z}_{m_{i+1}}^T \mathbf{I}_{m_{i+1}} \boldsymbol{\omega}_i + \frac{1}{2} k_{r,i+1}^2 \dot{q}_{i+1}^2 \mathbf{z}_{m_{i+1}}^T \mathbf{I}_{m_{i+1}} \mathbf{z}_{m_{i+1}}
 \end{aligned}$$

★ per il teorema di Steiner

$$\bar{\mathbf{I}}_{m_{i+1}} = \mathbf{I}_{m_{i+1}} + m_{m_{i+1}} \mathbf{S}^T(\mathbf{r}_{C_i, m_{i+1}}) \mathbf{S}(\mathbf{r}_{C_i, m_{i+1}})$$

- Energia cinetica

$$\begin{aligned}\mathcal{T}_i = & \frac{1}{2}m_i \dot{\mathbf{p}}_{C_i}^T \dot{\mathbf{p}}_{C_i} + \frac{1}{2}\boldsymbol{\omega}_i^T \bar{\mathbf{I}}_i \boldsymbol{\omega}_i + k_{r,i+1} \dot{q}_{i+1} \mathbf{z}_{m_{i+1}}^T \mathbf{I}_{m_{i+1}} \boldsymbol{\omega}_i \\ & + \frac{1}{2}k_{r,i+1}^2 \dot{q}_{i+1}^2 \mathbf{z}_{m_{i+1}}^T \mathbf{I}_{m_{i+1}} \mathbf{z}_{m_{i+1}}\end{aligned}$$

- ★ rotore con distribuzione di massa simmetrica intorno al suo asse di rotazione

$$\mathbf{I}_{m_i}^{m_i} = \begin{bmatrix} I_{m_i xx} & 0 & 0 \\ 0 & I_{m_i xx} & 0 \\ 0 & 0 & I_{m_i zz} \end{bmatrix}$$



$$\mathbf{I}_{m_{i+1}} \mathbf{z}_{m_{i+1}} = \mathbf{R}_{m_{i+1}} \mathbf{I}_{m_{i+1}}^{m_{i+1}} \mathbf{R}_{m_{i+1}}^T \mathbf{z}_{m_{i+1}} = I_{m_{i+1}} \mathbf{z}_{m_{i+1}}$$

- Quantità riferite alle terne solidali ai bracci

$$\begin{aligned}
 \mathcal{T}_i &= \frac{1}{2}m_i \dot{\mathbf{p}}_{C_i}^{iT} \dot{\mathbf{p}}_{C_i}^i + \frac{1}{2}\boldsymbol{\omega}_i^{iT} \bar{\mathbf{I}}_i^i \boldsymbol{\omega}_i^i + k_{r,i+1} \dot{q}_{i+1} I_{m_{i+1}} \mathbf{z}_{m_{i+1}}^{iT} \boldsymbol{\omega}_i^i \\
 &\quad + \frac{1}{2}k_{r,i+1}^2 \dot{q}_{i+1}^2 I_{m_{i+1}} \\
 &= \frac{1}{2}m_i \dot{\mathbf{p}}_i^{iT} \dot{\mathbf{p}}_i^i + \dot{\mathbf{p}}_i^{iT} \mathbf{S}(\boldsymbol{\omega}_i^i) m_i \mathbf{r}_{i,C_i}^i + \frac{1}{2}\boldsymbol{\omega}_i^{iT} \hat{\mathbf{I}}_i^i \boldsymbol{\omega}_i^i \\
 &\quad + k_{r,i+1} \dot{q}_{i+1} I_{m_{i+1}} \mathbf{z}_{m_{i+1}}^{iT} \boldsymbol{\omega}_i^i + \frac{1}{2}k_{r,i+1}^2 \dot{q}_{i+1}^2 I_{m_{i+1}}
 \end{aligned}$$

\* per il teorema di Steiner

$$\hat{\mathbf{I}}_i^i = \bar{\mathbf{I}}_i^i + m_i \mathbf{S}^T(\mathbf{r}_{i,C_i}^i) \mathbf{S}(\mathbf{r}_{i,C_i}^i)$$

- Energia cinetica *lineare* rispetto a:

- ★ massa  $m_i$
- ★ momento primo di inerzia

$$m_i \mathbf{r}_{i,C_i}^i = \begin{bmatrix} m_i \ell_{C_i x} \\ m_i \ell_{C_i y} \\ m_i \ell_{C_i z} \end{bmatrix}$$

- ★ tensore di inerzia

$$\hat{\mathbf{I}}_i^i = \begin{bmatrix} \bar{I}_{ixx} + m_i(\ell_{C_i y}^2 + \ell_{C_i z}^2) & -\bar{I}_{ixy} - m_i \ell_{C_i x} \ell_{C_i y} & \\ * & \bar{I}_{iyy} + m_i(\ell_{C_i z}^2 + \ell_{C_i x}^2) & \\ * & & * \\ & -\bar{I}_{ixz} - m_i \ell_{C_i x} \ell_{C_i z} & \\ & -\bar{I}_{iyz} - m_i \ell_{C_i y} \ell_{C_i z} & \\ & \bar{I}_{izz} + m_i(\ell_{C_i x}^2 + \ell_{C_i y}^2) & \end{bmatrix}$$

$$= \begin{bmatrix} \hat{I}_{ixx} & -\hat{I}_{ixy} & -\hat{I}_{ixz} \\ * & \hat{I}_{iyy} & \hat{I}_{iyz} \\ * & * & \hat{I}_{izz} \end{bmatrix}$$

- Energia potenziale

$$\begin{aligned}\mathcal{U}_i &= -m_i \mathbf{g}_0^{iT} \mathbf{p}_{C_i}^i \\ &= -\mathbf{g}_0^{iT} (m_i \mathbf{p}_i^i + m_i \mathbf{r}_{i,C_i}^i)\end{aligned}$$

*lineare* rispetto a:

- ★ massa  $m_i$
- ★ momento primo di inerzia

$$m_i \mathbf{r}_{i,C_i}^i = \begin{bmatrix} m_i \ell_{C_i x} \\ m_i \ell_{C_i y} \\ m_i \ell_{C_i z} \end{bmatrix}$$

- Lagrangiana

$$\mathcal{L} = \sum_{i=1}^n (\boldsymbol{\beta}_{\mathcal{T}i}^T - \boldsymbol{\beta}_{\mathcal{U}i}^T) \boldsymbol{\pi}_i$$

★ vettore ( $11 \times 1$ ) di parametri dinamici

$$\boldsymbol{\pi}_i = [m_i \ m_i \ell_{C_i x} \ m_i \ell_{C_i y} \ m_i \ell_{C_i z} \hat{I}_{ixx} \ \hat{I}_{ixy} \ \hat{I}_{ixz} \ \hat{I}_{iyx} \ \hat{I}_{iyz} \ \hat{I}_{izz} \ I_{m_i}]^T$$

- Le operazioni richieste dalle equazioni di Lagrange non alterano la linearità

$$\xi_i = \sum_{j=1}^n \boldsymbol{y}_{ij}^T \boldsymbol{\pi}_j$$

ove

$$\boldsymbol{y}_{ij} = \frac{d}{dt} \frac{\partial \boldsymbol{\beta}_{\mathcal{T}j}}{\partial \dot{q}_i} - \frac{\partial \boldsymbol{\beta}_{\mathcal{T}j}}{\partial q_i} + \frac{\partial \boldsymbol{\beta}_{\mathcal{U}j}}{\partial q_i}$$

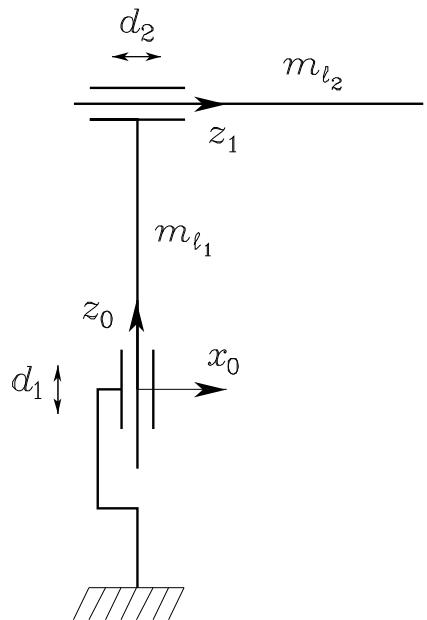
- Proprietà notevole

$$\begin{bmatrix} \xi_1 \\ \xi_2 \\ \vdots \\ \xi_n \end{bmatrix} = \begin{bmatrix} \mathbf{y}_{11}^T & \mathbf{y}_{12}^T & \dots & \mathbf{y}_{1n}^T \\ \mathbf{0}^T & \mathbf{y}_{22}^T & \dots & \mathbf{y}_{2n}^T \\ \vdots & \vdots & \ddots & \vdots \\ \mathbf{0}^T & \mathbf{0}^T & \dots & \mathbf{y}_{nn}^T \end{bmatrix} \begin{bmatrix} \boldsymbol{\pi}_1 \\ \boldsymbol{\pi}_2 \\ \vdots \\ \boldsymbol{\pi}_n \end{bmatrix}$$

★ in forma compatta:

$$\boldsymbol{\tau} = \mathbf{Y}(\mathbf{q}, \dot{\mathbf{q}}, \ddot{\mathbf{q}})\boldsymbol{\pi}$$

# Manipolatore cartesiano a due bracci



- Motori
  - ★ masse  $m_{m_i}$
  - ★ momenti di inerzia  $I_{m_i}$
  - ★  $\mathbf{p}_{m_i} = \mathbf{p}_{i-1} \quad z_{m_i} = z_{i-1}$

## ★ Jacobiani

$$\mathbf{J}_P^{(\ell_1)} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ 1 & 0 \end{bmatrix} \quad \mathbf{J}_P^{(\ell_2)} = \begin{bmatrix} 0 & 1 \\ 0 & 0 \\ 1 & 0 \end{bmatrix}$$

$$\mathbf{J}_P^{(m_1)} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ 0 & 0 \end{bmatrix} \quad \mathbf{J}_P^{(m_2)} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ 1 & 0 \end{bmatrix}$$

$$\mathbf{J}_O^{(m_1)} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ k_{r1} & 0 \end{bmatrix} \quad \mathbf{J}_O^{(m_2)} = \begin{bmatrix} 0 & k_{r2} \\ 0 & 0 \\ 0 & 0 \end{bmatrix}$$

- Matrice di inerzia

$$\mathbf{B} = \begin{bmatrix} m_{\ell_1} + m_{m_2} + k_{r1}^2 I_{m_1} + m_{\ell_2} & 0 \\ 0 & m_{\ell_2} + k_{r2}^2 I_{m_2} \end{bmatrix}$$

- Forze gravitazionali

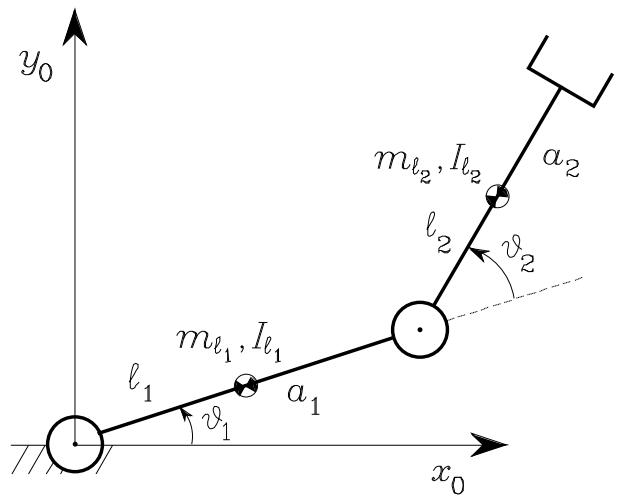
$$g_1 = (m_{\ell_1} + m_{m_2} + m_{\ell_2})g \quad g_2 = 0$$

- Equazioni del moto

$$(m_{\ell_1} + m_{m_2} + k_{r1}^2 I_{m_1} + m_{\ell_2})\ddot{d}_1 + (m_{\ell_1} + m_{m_2} + m_{\ell_2})g = f_1$$

$$(m_{\ell_2} + k_{r2}^2 I_{m_2})\ddot{d}_2 = f_2$$

## Manipolatore planare a due bracci



- Motori
  - ★ masse  $m_{m_i}$
  - ★ momenti di inerzia  $I_{m_i}$
  - ★  $p_{m_i} = p_{i-1} \quad z_{m_i} = z_{i-1}$

★ Jacobiani

$$\mathbf{J}_P^{(\ell_1)} = \begin{bmatrix} -\ell_1 s_1 & 0 \\ \ell_1 c_1 & 0 \\ 0 & 0 \end{bmatrix} \quad \mathbf{J}_P^{(\ell_2)} = \begin{bmatrix} -a_1 s_1 - \ell_2 s_{12} & -\ell_2 s_{12} \\ a_1 c_1 + \ell_2 c_{12} & \ell_2 c_{12} \\ 0 & 0 \end{bmatrix}$$

$$\mathbf{J}_O^{(\ell_1)} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ 1 & 0 \end{bmatrix} \quad \mathbf{J}_O^{(\ell_2)} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ 1 & 1 \end{bmatrix}$$

$$\mathbf{J}_P^{(m_1)} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ 0 & 0 \end{bmatrix} \quad \mathbf{J}_P^{(m_2)} = \begin{bmatrix} -a_1 s_1 & 0 \\ a_1 c_1 & 0 \\ 0 & 0 \end{bmatrix}$$

$$\mathbf{J}_O^{(m_1)} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ k_{r1} & 0 \end{bmatrix} \quad \mathbf{J}_O^{(m_2)} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ 1 & k_{r2} \end{bmatrix}$$

- Matrice di inerzia

$$\mathbf{B}(\mathbf{q}) = \begin{bmatrix} b_{11}(\vartheta_2) & b_{12}(\vartheta_2) \\ b_{21}(\vartheta_2) & b_{22} \end{bmatrix}$$

$$\begin{aligned} b_{11} &= I_{\ell_1} + m_{\ell_1} \ell_1^2 + k_{r1}^2 I_{m_1} + I_{\ell_2} + m_{\ell_2} (a_1^2 + \ell_2^2 + 2a_1 \ell_2 c_2) \\ &\quad + I_{m_2} + m_{m_2} a_1^2 \end{aligned}$$

$$b_{12} = b_{21} = I_{\ell_2} + m_{\ell_2} (\ell_2^2 + a_1 \ell_2 c_2) + k_{r2} I_{m_2}$$

$$b_{22} = I_{\ell_2} + m_{\ell_2} \ell_2^2 + k_{r2}^2 I_{m_2}$$

- Forze centrifughe e di Coriolis

$$c_{111} = \frac{1}{2} \frac{\partial b_{11}}{\partial q_1} = 0$$

$$c_{112} = c_{121} = \frac{1}{2} \frac{\partial b_{11}}{\partial q_2} = -m_{\ell_2} a_1 \ell_2 s_2 = h$$

$$c_{122} = \frac{\partial b_{12}}{\partial q_2} - \frac{1}{2} \frac{\partial b_{22}}{\partial q_1} = h$$

$$c_{211} = \frac{\partial b_{21}}{\partial q_1} - \frac{1}{2} \frac{\partial b_{11}}{\partial q_2} = -h$$

$$c_{212} = c_{221} = \frac{1}{2} \frac{\partial b_{22}}{\partial q_1} = 0$$

$$c_{222} = \frac{1}{2} \frac{\partial b_{22}}{\partial q_2} = 0$$

$$\mathbf{C}(\mathbf{q}, \dot{\mathbf{q}}) = \begin{bmatrix} h\dot{\vartheta}_2 & h(\dot{\vartheta}_1 + \dot{\vartheta}_2) \\ -h\dot{\vartheta}_1 & 0 \end{bmatrix}$$

★ anti-simmetria

$$\begin{aligned} \mathbf{N}(\mathbf{q}, \dot{\mathbf{q}}) &= \dot{\mathbf{B}}(\mathbf{q}) - 2\mathbf{C}(\mathbf{q}, \dot{\mathbf{q}}) \\ &= \begin{bmatrix} 2h\dot{\vartheta}_2 & h\dot{\vartheta}_2 \\ h\dot{\vartheta}_2 & 0 \end{bmatrix} - 2 \begin{bmatrix} h\dot{\vartheta}_2 & h(\dot{\vartheta}_1 + \dot{\vartheta}_2) \\ -h\dot{\vartheta}_1 & 0 \end{bmatrix} \\ &= \begin{bmatrix} 0 & -2h\dot{\vartheta}_1 - h\dot{\vartheta}_2 \\ 2h\dot{\vartheta}_1 + h\dot{\vartheta}_2 & 0 \end{bmatrix} \end{aligned}$$

- Forze gravitazionali

$$\begin{aligned} g_1 &= (m_{\ell_1}\ell_1 + m_{m_2}a_1 + m_{\ell_2}a_1)gc_1 + m_{\ell_2}\ell_2gc_{12} \\ g_2 &= m_{\ell_2}\ell_2gc_{12} \end{aligned}$$

- Equazioni del moto

$$\begin{aligned}
 & (I_{\ell_1} + m_{\ell_1}\ell_1^2 + k_{r1}^2 I_{m_1} + I_{\ell_2} + m_{\ell_2}(a_1^2 + \ell_2^2 + 2a_1\ell_2c_2) \\
 & \quad + I_{m_2} + m_{m_2}a_1^2) \ddot{\vartheta}_1 \\
 & + (I_{\ell_2} + m_{\ell_2}(\ell_2^2 + a_1\ell_2c_2) + k_{r2}I_{m_2}) \ddot{\vartheta}_2 \\
 & - 2m_{\ell_2}a_1\ell_2s_2\dot{\vartheta}_1\dot{\vartheta}_2 - m_{\ell_2}a_1\ell_2s_2\dot{\vartheta}_2^2 \\
 & + (m_{\ell_1}\ell_1 + m_{m_2}a_1 + m_{\ell_2}a_1)gc_1 + m_{\ell_2}\ell_2gc_{12} = \tau_1 \\
 & (I_{\ell_2} + m_{\ell_2}\ell_2^2 + k_{r2}^2 I_{m_2}) \ddot{\vartheta}_2 + (I_{\ell_2} + m_{\ell_2}(\ell_2^2 + a_1\ell_2c_2) + k_{r2}I_{m_2}) \ddot{\vartheta}_1 \\
 & + m_{\ell_2}a_1\ell_2s_2\dot{\vartheta}_1^2 + m_{\ell_2}\ell_2gc_{12} = \tau_2
 \end{aligned}$$

- Parametrizzazione del modello

$$\boldsymbol{\pi} = [\pi_1 \quad \pi_2 \quad \pi_3 \quad \pi_4 \quad \pi_5 \quad \pi_6 \quad \pi_7 \quad \pi_8]^T$$

$$\pi_1 = m_1 = m_{\ell_1} + m_{m_2}$$

$$\pi_2 = m_1 \ell_{C_1} = m_{\ell_1} (\ell_1 - a_1)$$

$$\pi_3 = \hat{I}_1 = I_{\ell_1} + m_{\ell_1} (\ell_1 - a_1)^2 + I_{m_2}$$

$$\pi_4 = I_{m_1}$$

$$\pi_5 = m_2 = m_{\ell_2}$$

$$\pi_6 = m_{\ell_2} \ell_{C_2} = m_{\ell_2} (\ell_2 - a_2)$$

$$\pi_7 = \hat{I}_2 = I_{\ell_2} + m_{\ell_2} (\ell_2 - a_2)^2$$

$$\pi_8 = I_{m_2}$$

$$\mathbf{Y} = \begin{bmatrix} y_{11} & y_{12} & y_{13} & y_{14} & y_{15} & y_{16} & y_{17} & y_{18} \\ y_{21} & y_{22} & y_{23} & y_{24} & y_{25} & y_{26} & y_{27} & y_{28} \end{bmatrix}$$

$$y_{11} = a_1^2 \ddot{\vartheta}_1 + a_1 g c_1$$

$$y_{12} = 2a_1 \ddot{\vartheta}_1 + g c_1$$

$$y_{13} = \ddot{\vartheta}_1$$

$$y_{14} = k_{r1}^2 \ddot{\vartheta}_1$$

$$y_{15} = (a_1^2 + 2a_1 a_2 c_2 + a_2^2) \ddot{\vartheta}_1 + (a_1 a_2 c_2 + a_2^2) \ddot{\vartheta}_2$$

$$- 2a_1 a_2 s_2 \dot{\vartheta}_1 \dot{\vartheta}_2 - a_1 a_2 s_2 \dot{\vartheta}_2^2 + a_1 g c_1 + a_2 g c_{12}$$

$$y_{16} = (2a_1 c_2 + 2a_2) \ddot{\vartheta}_1 + (a_1 c_2 + 2a_2) \ddot{\vartheta}_2 - 2a_1 s_2 \dot{\vartheta}_1 \dot{\vartheta}_2$$

$$- a_1 s_2 \dot{\vartheta}_2^2 + g c_{12}$$

$$y_{17} = \ddot{\vartheta}_1 + \ddot{\vartheta}_2$$

$$y_{18} = \ddot{\vartheta}_1 + k_{r2} \ddot{\vartheta}_2$$

$$y_{21} = 0$$

$$y_{22} = 0$$

$$y_{23} = 0$$

$$y_{24} = 0$$

$$y_{25} = (a_1 a_2 c_2 + a_2^2) \ddot{\vartheta}_1 + a_2^2 \ddot{\vartheta}_2 + a_1 a_2 s_2 \dot{\vartheta}_1^2 + a_2 g c_{12}$$

$$y_{26} = (a_1 c_2 + 2a_2) \ddot{\vartheta}_1 + 2a_2 \ddot{\vartheta}_2 + a_1 s_2 \dot{\vartheta}_1^2 + g c_{12}$$

$$y_{27} = \ddot{\vartheta}_1 + \ddot{\vartheta}_2$$

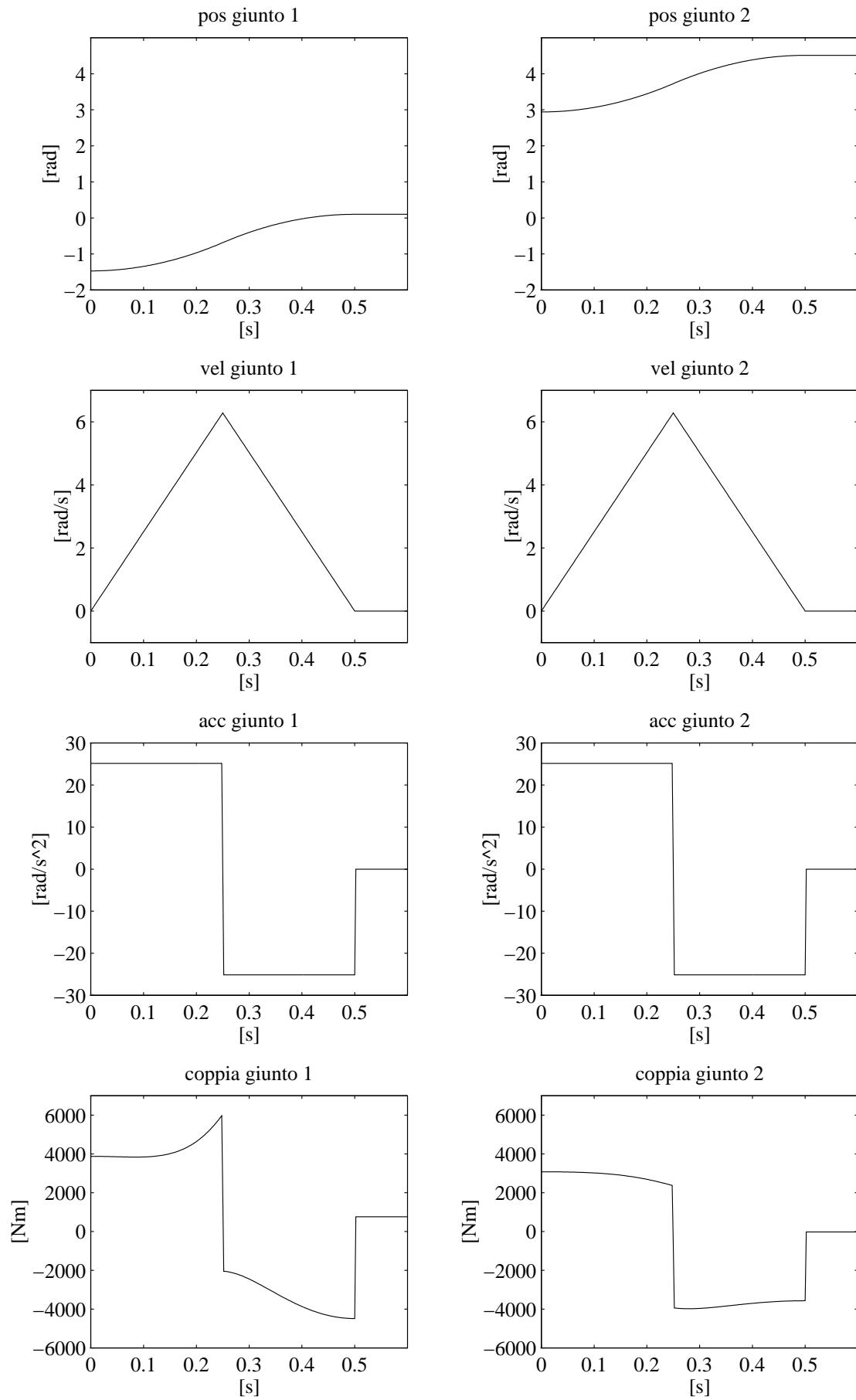
$$y_{28} = k_{r2} \ddot{\vartheta}_1 + k_{r2}^2 \ddot{\vartheta}_2.$$

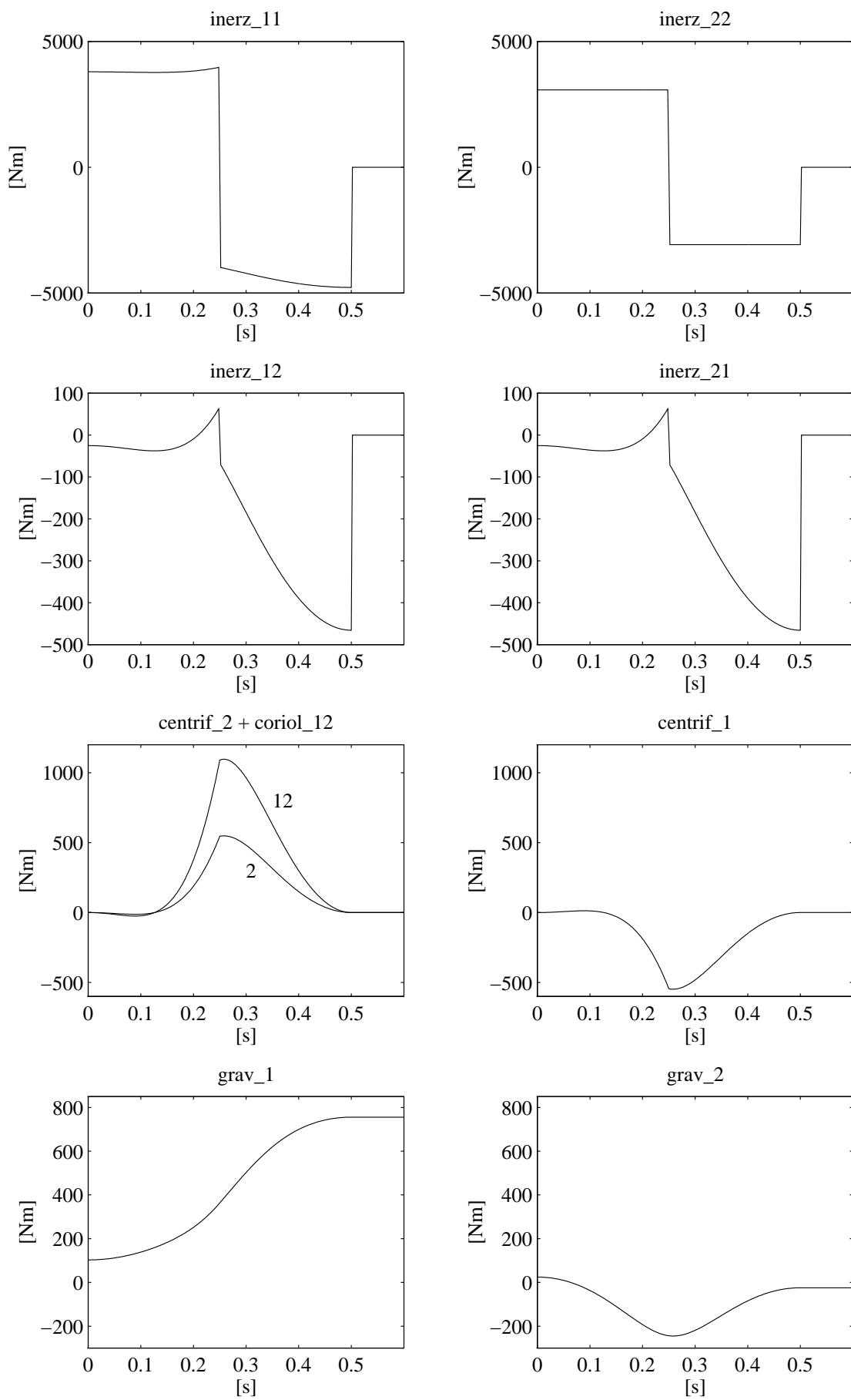
## Esempio

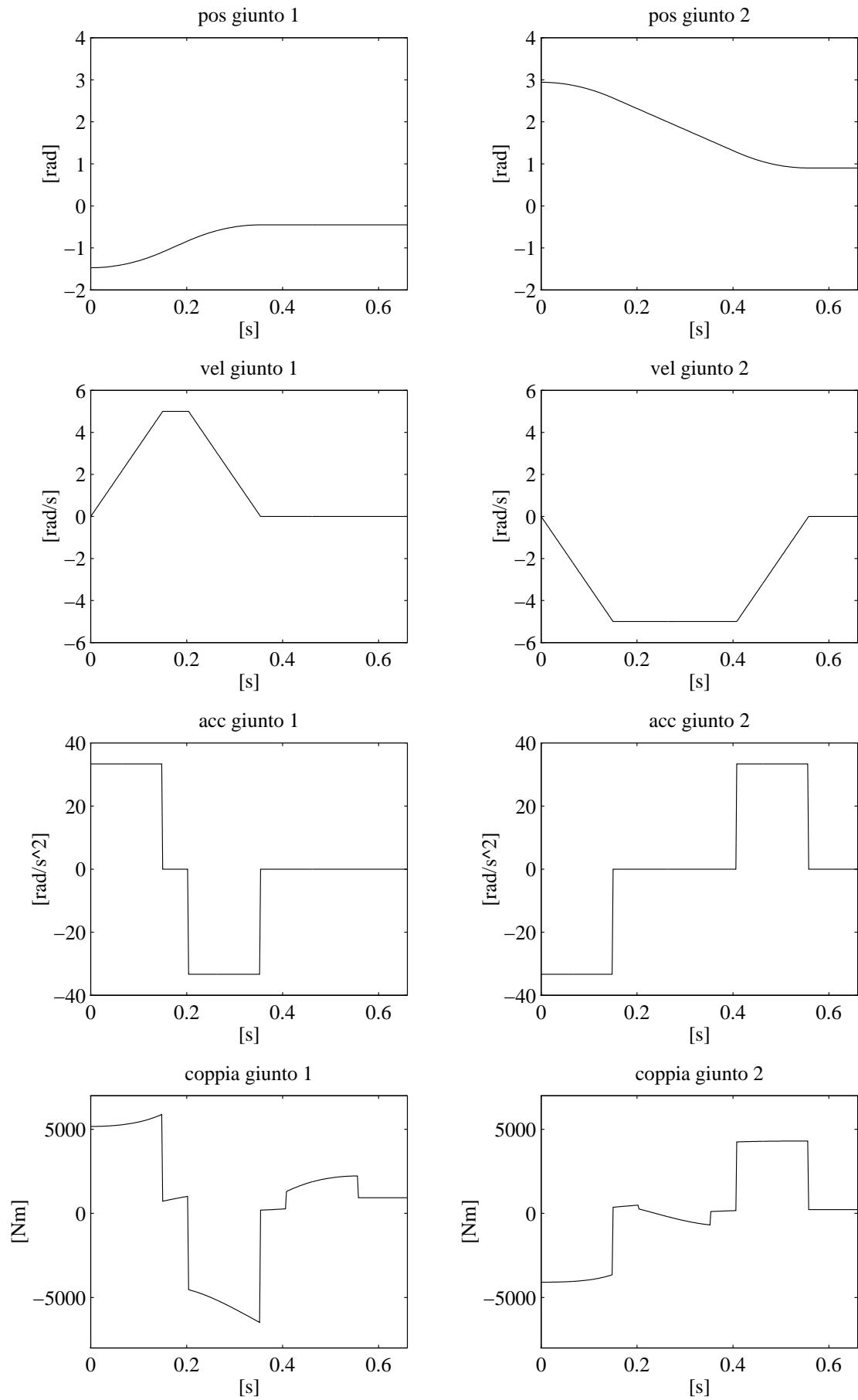
- Manipolatore planare a due bracci

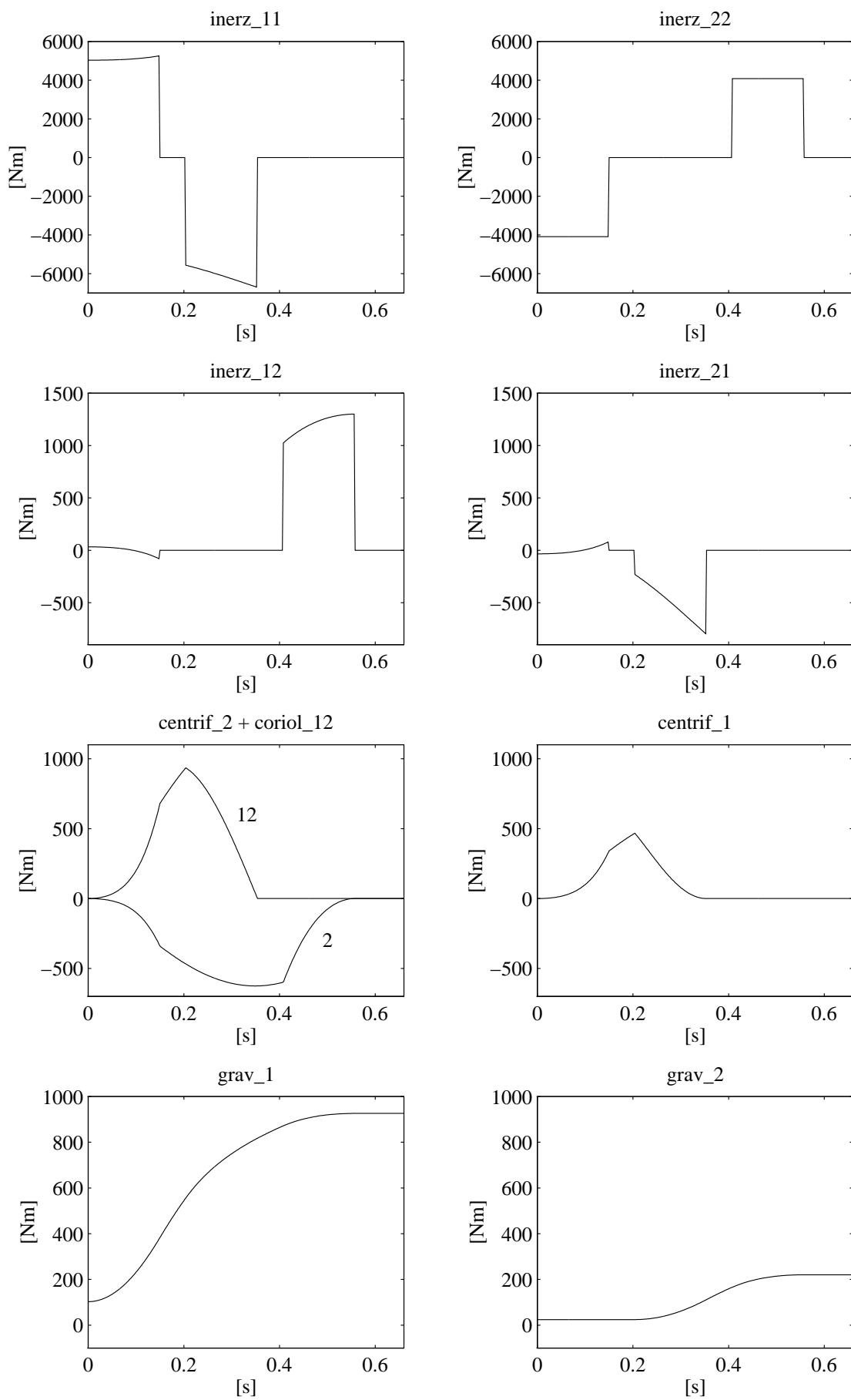
$$\star \quad a_1 = a_2 = 1 \text{ m} \quad \ell_1 = \ell_2 = 0.5 \text{ m}$$
$$m_{\ell_1} = m_{\ell_2} = 50 \text{ kg} \quad I_{\ell_1} = I_{\ell_2} = 10 \text{ kg} \cdot \text{m}^2$$

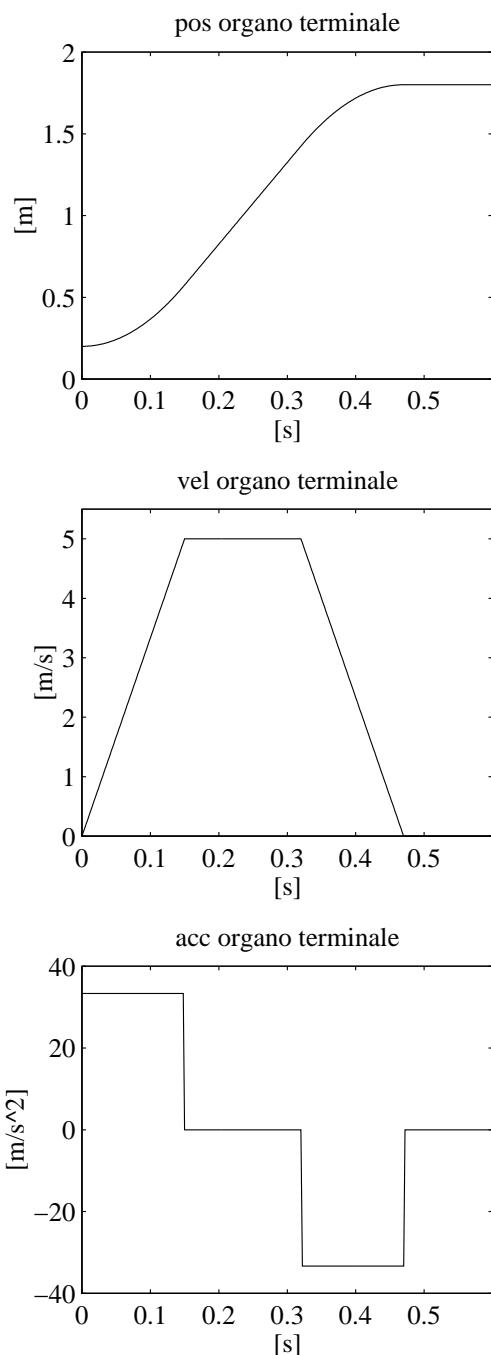
$$\star \quad k_{r1} = k_{r2} = 100$$
$$m_{m_1} = m_{m_2} = 5 \text{ kg} \quad I_{m_1} = I_{m_2} = 0.01 \text{ kg} \cdot \text{m}^2$$

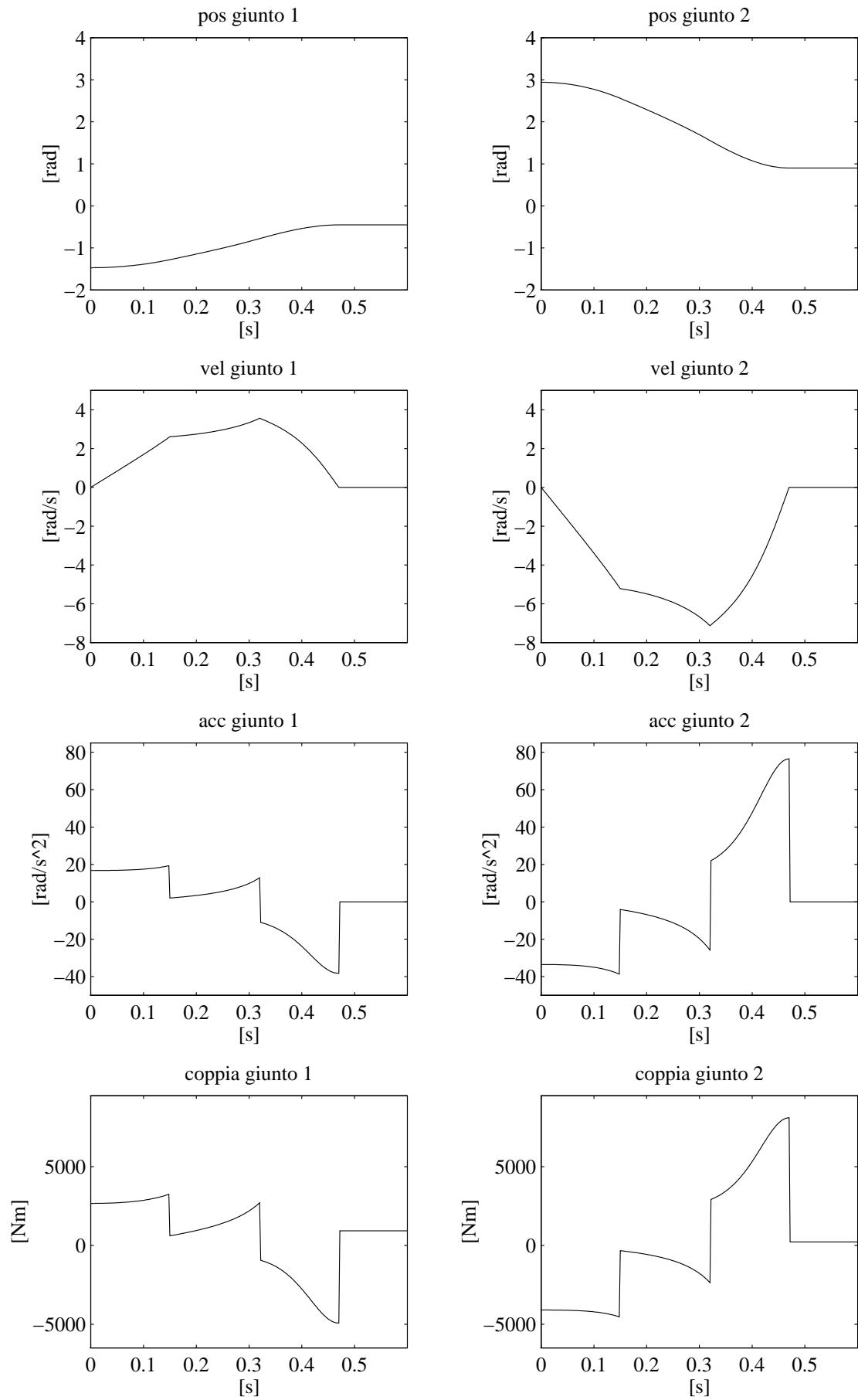


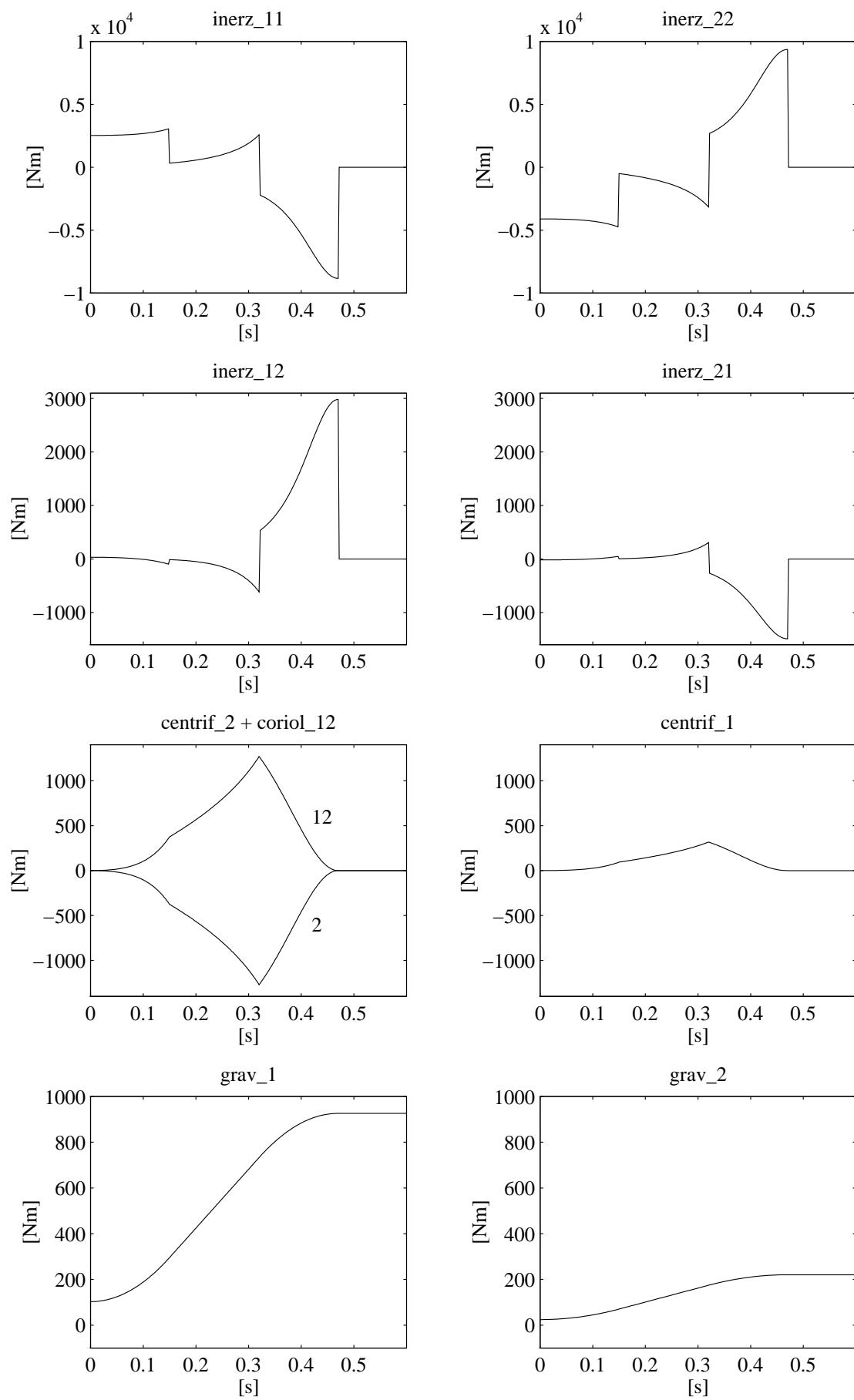




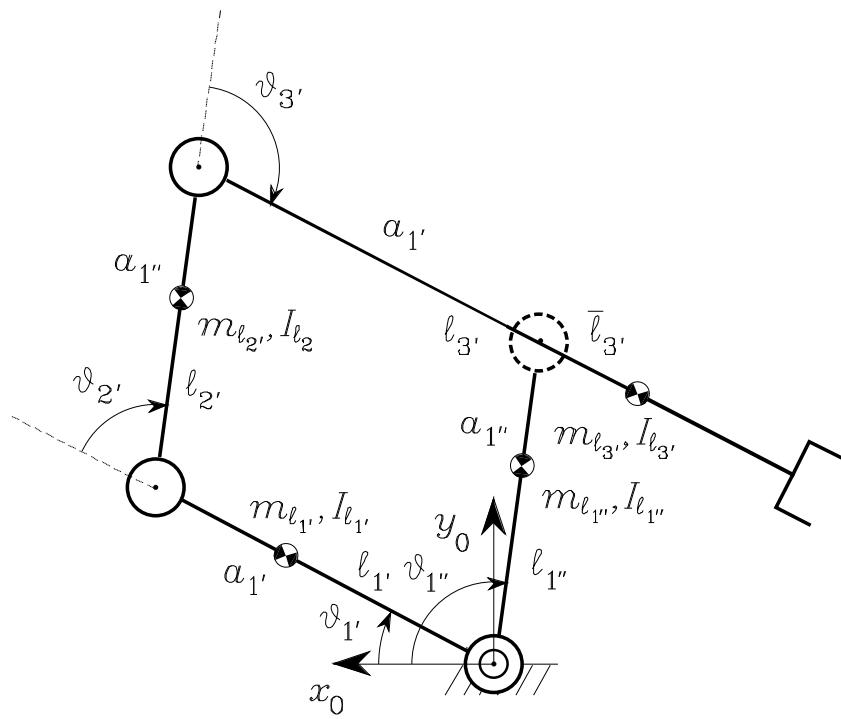








# Manipolatore a parallelogramma



## ★ Jacobiani

$$\mathbf{J}_P^{(\ell_{1'})} = \begin{bmatrix} -\ell_{1'} s_{1'} & 0 & 0 \\ \ell_{1'} c_{1'} & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad \mathbf{J}_P^{(\ell_{2'})} = \begin{bmatrix} -a_{1'} s_{1'} - \ell_2 s_{1'2'} & -\ell_{2'} s_{1'2'} & 0 \\ a_{1'} c_{1'} + \ell_{2'} c_{1'2'} & \ell_2 c_{1'2'} & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$\mathbf{J}_P^{(\ell_{3'})} = \begin{bmatrix} -a_{1'} s_{1'} - a_{2'} s_{1'2'} - \ell_{3'} s_{1'2'3'} & -a_{2'} s_{1'2'} - \ell_{3'} s_{1'2'3'} & -\ell_{3'} s_{1'2'3'} \\ a_{1'} c_{1'} + a_{2'} c_{1'2'} + \ell_{3'} c_{1'2'3'} & a_{2'} c_{1'2'} + \ell_{3'} c_{1'2'3'} & \ell_{3'} c_{1'2'3'} \\ 0 & 0 & 0 \end{bmatrix}$$

$$\mathbf{J}_P^{(\ell_{1''})} = \begin{bmatrix} -\ell_{1''} s_{1''} \\ \ell_{1''} c_{1''} \\ 0 \end{bmatrix}$$

$$\mathbf{J}_O^{(\ell_{1'})} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix} \quad \mathbf{J}_O^{(\ell_{2'})} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 1 & 0 \end{bmatrix} \quad \mathbf{J}_O^{(\ell_{3'})} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 1 & 1 \end{bmatrix}$$

$$\mathbf{J}_O^{(\ell_{1''})} = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

- Matrici di inerzia per i manipolatori virtuali

$$\mathbf{B}'(\mathbf{q}') = \begin{bmatrix} b_{1'1'}(\vartheta_{2'}, \vartheta_{3'}) & b_{1'2'}(\vartheta_{2'}, \vartheta_{3'}) & b_{1'3'}(\vartheta_{2'}, \vartheta_{3'}) \\ b_{2'1'}(\vartheta_{2'}, \vartheta_{3'}) & b_{2'2'}(\vartheta_{3'}) & b_{2'3'}(\vartheta_{3'}) \\ b_{3'1'}(\vartheta_{2'}, \vartheta_{3'}) & b_{3'2'}(\vartheta_{3'}) & b_{3'3'} \end{bmatrix}$$

$$b_{1'1'} = I_{\ell_{1'}} + m_{\ell_{1'}} \ell_{1'}^2 + I_{\ell_{2'}} + m_{\ell_{2'}} (a_{1'}^2 + \ell_{2'}^2 + 2a_{1'} \ell_{2'} c_{2'}) + I_{\ell_{3'}} \\ + m_{\ell_{3'}} (a_{1'}^2 + a_{2'}^2 + \ell_{3'}^2 + 2a_{1'} a_{2'} c_{2'} + 2a_{1'} \ell_{3'} c_{2'3'} + 2a_{2'} \ell_{3'} c_{3'})$$

$$b_{1'2'} = b_{2'1'} = I_{\ell_{2'}} + m_{\ell_{2'}} (\ell_{2'}^2 + a_{1'} \ell_{2'} c_{2'}) + I_{\ell_{3'}} \\ + m_{\ell_{3'}} (a_{2'}^2 + \ell_{3'}^2 + a_{1'} a_{2'} c_{2'} + a_{1'} \ell_{3'} c_{2'3'} + 2a_{2'} \ell_{3'} c_{3'})$$

$$b_{1'3'} = b_{31} = I_{\ell_{3'}} + m_{\ell_{3'}} (\ell_{3'}^2 + a_{1'} \ell_{3'} c_{2'3'} + a_{2'} \ell_{3'} c_{3'})$$

$$b_{2'2'} = I_{\ell_{2'}} + m_{\ell_{2'}} \ell_{2'}^2 + I_{\ell_{3'}} + m_{\ell_{3'}} (a_{2'}^2 + \ell_{3'}^2 + 2a_{2'} \ell_{3'} c_{3'})$$

$$b_{2'3'} = I_{\ell_{3'}} + m_{\ell_{3'}} (\ell_{3'}^2 + a_{2'} \ell_{3'} c_{3'})$$

$$b_{3'3'} = I_{\ell_{3'}} + m_{\ell_{3'}} \ell_{3'}^2$$

$$b_{1''1''} = I_{\ell_{1''}} + m_{\ell_{1''}} \ell_{1''}^2$$

$$\tau_{i'} = \sum_{j'=1'}^{3'} b_{i'j'} \ddot{\vartheta}_{j'} \quad \tau_{1''} = b_{1'1'} \ddot{\vartheta}_{1''}$$

- Manipolatore a catena chiusa

$$\boldsymbol{\tau}_a = \mathbf{B}_a \ddot{\mathbf{q}}_a$$

$$\mathbf{B}_a = \begin{bmatrix} b_{a11} & b_{a12} \\ b_{a21} & b_{a22} \end{bmatrix}$$

$$b_{a11} = I_{\ell_1'} + m_{\ell_1'} \ell_{1'}^2 + m_{\ell_2'} a_{1'}^2,$$

$$+ I_{\ell_3'} + m_{\ell_3'} \ell_{3'}^2 + m_{\ell_3'} a_{1'}^2 - 2a_{1'} m_{\ell_3'} \ell_{3'}$$

$$b_{a12} = b_{a21} = (a_{1'} m_{\ell_2'} \ell_{2'} + a_{1''} m_{\ell_3'} (a_{1'} - \ell_{3'})) \cos(\vartheta_{1''} - \vartheta_{1'})$$

$$b_{a22} = I_{\ell_1'} + m_{\ell_1'} \ell_{1'}^2 + I_{\ell_2'} + m_{\ell_2'} \ell_{2'}^2 + m_{\ell_3'} a_{1''}^2.$$

★ matrice costante e diagonale

$$\frac{m_{\ell_3'} \bar{\ell}_{3'}}{m_{\ell_2'} \ell_{2'}} = \frac{a_{1'}}{a_{1''}}$$

$$b_{a11} = I_{\ell_1'} + m_{\ell_1'} \ell_{1'}^2 + m_{\ell_2'} a_{1'}^2 \left( 1 + \frac{\ell_{2'} \bar{\ell}_{3'}}{a_{1'} a_{1''}} \right) + I_{\ell_3'}$$

$$b_{a22} = I_{\ell_1'} + m_{\ell_1'} \ell_{1'}^2 + I_{\ell_2'} + m_{\ell_2'} \ell_{2'}^2 \left( 1 + \frac{a_{1'} a_{1''}}{\ell_{2'} \bar{\ell}_{3'}} \right)$$

- Forze gravitazionali

$$g_{1'} = (m_{\ell_1} \ell_{1'} + m_{\ell_2} a_{1'} + m_{\ell_3} a_{1'}) g c_{1'} + (m_{\ell_2} \ell_{2'} + m_{\ell_3} a_{2'}) g c_{1'2'}$$
$$+ m_{\ell_3} \ell_{3'} g c_{1'2'3}$$

$$g_{2'} = (m_{\ell_2} \ell_{2'} + m_{\ell_3} a_{2'}) g c_{1'2'} + m_{\ell_3} \ell_{3'} g c_{1'2'3}$$

$$g_{3'} = m_{\ell_3} \ell_{3'} g c_{1'2'3}$$

$$g_{1''} = m_{\ell_1''} \ell_{1''} g c_{1''}$$

$$\mathbf{g}_a = \begin{bmatrix} (m_{\ell_1} \ell_{1'} + m_{\ell_2} a_{1'} - m_{\ell_3} \bar{\ell}_{3'}) g c_{1'} \\ (m_{\ell_1''} \ell_{1''} + m_{\ell_2'} \ell_{2'} + m_{\ell_3} a_{1''}) g c_{1''} \end{bmatrix}$$

# IDENTIFICAZIONE DEI PARAMETRI DINAMICI

- Impiego del modello dinamico per simulazione e controllo
- Stime approssimate dei parametri dinamici
  - ★ semplificazioni di modellazione geometrica
  - ★ effetti dinamici complessi (attriti, etc.)
- Tecniche di identificazione
  - ★ proprietà di linearità

$$\boldsymbol{\tau} = \mathbf{Y}(\boldsymbol{q}, \dot{\boldsymbol{q}}, \ddot{\boldsymbol{q}})\boldsymbol{\pi}$$

- Misure
  - ★  $\boldsymbol{q}$  e  $\dot{\boldsymbol{q}}$  dirette (encoder e tachimetri)
  - ★  $\ddot{\boldsymbol{q}}$  indirette (filtro ricostruttore)
  - ★  $\boldsymbol{\tau}$  dirette (sensore di coppia) o indirette (misure di forza o corrente)

- $N$  insiemi di misure

$$\bar{\boldsymbol{\tau}} = \begin{bmatrix} \boldsymbol{\tau}(t_1) \\ \vdots \\ \boldsymbol{\tau}(t_N) \end{bmatrix} = \begin{bmatrix} \mathbf{Y}(t_1) \\ \vdots \\ \mathbf{Y}(t_N) \end{bmatrix} \boldsymbol{\pi} = \bar{\mathbf{Y}} \boldsymbol{\pi}$$

- Soluzione (a minimi quadrati)

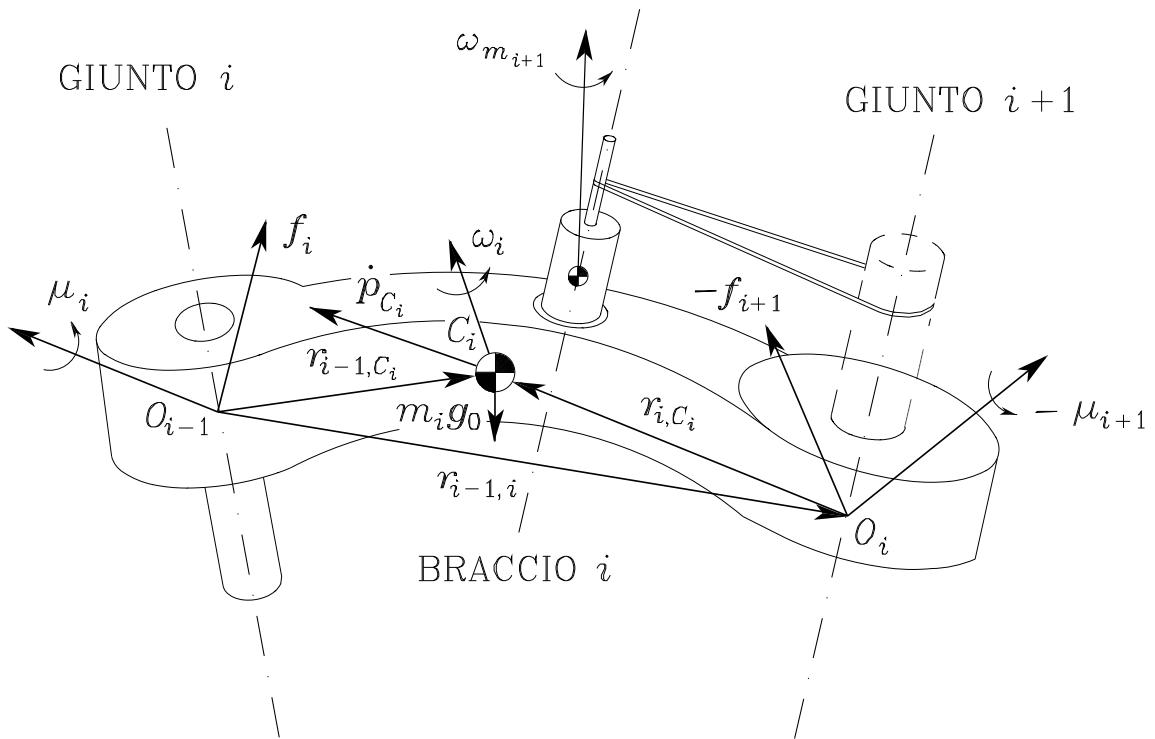
$$\boldsymbol{\pi} = (\bar{\mathbf{Y}}^T \bar{\mathbf{Y}})^{-1} \bar{\mathbf{Y}}^T \bar{\boldsymbol{\tau}}$$

- ★  $\mathbf{Y}$  triangolare  $\implies$  procedimento iterativo
- ★ se  $\varrho(\mathbf{Y}) < p$   $\implies \mathbf{Y}^*$

- Parametri
  - ★ identificabili
  - ★ non identificabili
  - ★ identificabili in combinazione lineare
- Carico
  - ★ modifica parametri dell'ultimo braccio
- Traiettorie
  - ★ di tipo polinomiale (sufficientemente *ricche*)
  - ★ attenzione alle dinamiche non modellate

# FORMULAZIONE DI NEWTON-EULERO

- ★ bilancio di forze e momenti agenti sul generico braccio



- Equazione di *Newton* (moto *traslazionale* del baricentro)

$$\mathbf{f}_i - \mathbf{f}_{i+1} + m_i \mathbf{g}_0 = m_i \ddot{\mathbf{p}}_{C_i}$$

- Equazione di *Eulero* (moto *rotazionale*)

$$\begin{aligned} \boldsymbol{\mu}_i + \mathbf{f}_i \times \mathbf{r}_{i-1,C_i} - \boldsymbol{\mu}_{i+1} - \mathbf{f}_{i+1} \times \mathbf{r}_{i,C_i} \\ = \frac{d}{dt} (\bar{\mathbf{I}}_i \boldsymbol{\omega}_i + k_{r,i+1} \dot{q}_{i+1} I_{m_{i+1}} \mathbf{z}_{m_{i+1}}) \end{aligned}$$

$$\begin{aligned}
\frac{d}{dt}(\bar{\mathbf{I}}_i \boldsymbol{\omega}_i) &= \dot{\mathbf{R}}_i \bar{\mathbf{I}}_i^i \mathbf{R}_i^T \boldsymbol{\omega}_i + \mathbf{R}_i \bar{\mathbf{I}}_i^i \dot{\mathbf{R}}_i^T \boldsymbol{\omega}_i + \mathbf{R}_i \bar{\mathbf{I}}_i^i \mathbf{R}_i^T \dot{\boldsymbol{\omega}}_i \\
&= \mathbf{S}(\boldsymbol{\omega}_i) \mathbf{R}_i \bar{\mathbf{I}}_i^i \mathbf{R}_i^T \boldsymbol{\omega}_i + \mathbf{R}_i \bar{\mathbf{I}}_i^i \mathbf{R}_i^T \mathbf{S}^T(\boldsymbol{\omega}_i) \boldsymbol{\omega}_i + \mathbf{R}_i \bar{\mathbf{I}}_i^i \mathbf{R}_i^T \dot{\boldsymbol{\omega}}_i \\
&= \bar{\mathbf{I}}_i \dot{\boldsymbol{\omega}}_i + \boldsymbol{\omega}_i \times (\bar{\mathbf{I}}_i \boldsymbol{\omega}_i)
\end{aligned}$$

$$\frac{d}{dt}(\dot{q}_{i+1} I_{m_{i+1}} z_{m_{i+1}}) = \ddot{q}_{i+1} I_{m_{i+1}} z_{m_{i+1}} + \dot{q}_{i+1} I_{m_{i+1}} \boldsymbol{\omega}_i \times z_{m_{i+1}}$$

↓

- Equazione di Eulero

$$\begin{aligned}
\boldsymbol{\mu}_i + \mathbf{f}_i \times \mathbf{r}_{i-1,C_i} - \boldsymbol{\mu}_{i+1} - \mathbf{f}_{i+1} \times \mathbf{r}_{i,C_i} &= \bar{\mathbf{I}}_i \dot{\boldsymbol{\omega}}_i + \boldsymbol{\omega}_i \times (\bar{\mathbf{I}}_i \boldsymbol{\omega}_i) \\
&+ k_{r,i+1} \ddot{q}_{i+1} I_{m_{i+1}} z_{m_{i+1}} + k_{r,i+1} \dot{q}_{i+1} I_{m_{i+1}} \boldsymbol{\omega}_i \times z_{m_{i+1}}
\end{aligned}$$

- Forza generalizzata al giunto

$$\tau_i = \begin{cases} \mathbf{f}_i^T \mathbf{z}_{i-1} + k_{ri} I_{mi} \dot{\boldsymbol{\omega}}_{m_i}^T \mathbf{z}_{m_i} & \text{per un giunto prismatico} \\ \boldsymbol{\mu}_i^T \mathbf{z}_{i-1} + k_{ri} I_{mi} \dot{\boldsymbol{\omega}}_{m_i}^T \mathbf{z}_{m_i} & \text{per un giunto rotoidale} \end{cases}$$

# Accelerazioni di un braccio

$$\boldsymbol{\omega}_i = \begin{cases} \boldsymbol{\omega}_{i-1} & \text{prismatico} \\ \boldsymbol{\omega}_{i-1} + \dot{\vartheta}_i \boldsymbol{z}_{i-1} & \text{rotoidale} \end{cases}$$

$$\dot{\boldsymbol{p}}_i = \begin{cases} \dot{\boldsymbol{p}}_{i-1} + \dot{d}_i \boldsymbol{z}_{i-1} + \boldsymbol{\omega}_i \times \boldsymbol{r}_{i-1,i} & \text{prismatico} \\ \dot{\boldsymbol{p}}_{i-1} + \boldsymbol{\omega}_i \times \boldsymbol{r}_{i-1,i} & \text{rotoidale} \end{cases}$$

$$\dot{\boldsymbol{\omega}}_i = \begin{cases} \dot{\boldsymbol{\omega}}_{i-1} & \text{prismatico} \\ \dot{\boldsymbol{\omega}}_{i-1} + \ddot{\vartheta}_i \boldsymbol{z}_{i-1} + \dot{\vartheta}_i \boldsymbol{\omega}_{i-1} \times \boldsymbol{z}_{i-1} & \text{rotoidale} \end{cases}$$

$$\ddot{\boldsymbol{p}}_i = \begin{cases} \ddot{\boldsymbol{p}}_{i-1} + \ddot{d}_i \boldsymbol{z}_{i-1} + 2\dot{d}_i \boldsymbol{\omega}_i \times \boldsymbol{z}_{i-1} \\ \quad + \dot{\boldsymbol{\omega}}_i \times \boldsymbol{r}_{i-1,i} + \boldsymbol{\omega}_i \times (\boldsymbol{\omega}_i \times \boldsymbol{r}_{i-1,i}) & \text{prismatico} \\ \ddot{\boldsymbol{p}}_{i-1} + \dot{\boldsymbol{\omega}}_i \times \boldsymbol{r}_{i-1,i} \\ \quad + \boldsymbol{\omega}_i \times (\boldsymbol{\omega}_i \times \boldsymbol{r}_{i-1,i}) & \text{rotoidale} \end{cases}$$

$$\ddot{\boldsymbol{p}}_{C_i} = \ddot{\boldsymbol{p}}_i + \dot{\boldsymbol{\omega}}_i \times \boldsymbol{r}_{i,C_i} + \boldsymbol{\omega}_i \times (\boldsymbol{\omega}_i \times \boldsymbol{r}_{i,C_i})$$

$$\dot{\boldsymbol{\omega}}_{m_i} = \dot{\boldsymbol{\omega}}_{i-1} + k_{ri} \ddot{q}_i \boldsymbol{z}_{mi} + k_{ri} \dot{q}_i \boldsymbol{\omega}_{i-1} \times \boldsymbol{z}_{mi}$$

## Algoritmo ricorsivo

- Ricorsione *in avanti* di velocità e accelerazioni
  - ★ condizioni iniziali:  $\omega_0, \ddot{\mathbf{p}}_0 - \mathbf{g}_0, \dot{\omega}_0$
  - ★ calcolo di  $\omega_i, \dot{\omega}_i, \ddot{\mathbf{p}}_i, \ddot{\mathbf{p}}_{C_i}, \dot{\omega}_{m_i}$
- Ricorsione *all'indietro* di forze e momenti
  - ★ condizioni terminali:  $\mathbf{h} = [\mathbf{f}_{n+1}^T \quad \boldsymbol{\mu}_{n+1}^T]^T$
  - ★ calcolo di

$$\mathbf{f}_i = \mathbf{f}_{i+1} + m_i \ddot{\mathbf{p}}_{C_i}$$

$$\begin{aligned} \boldsymbol{\mu}_i &= -\mathbf{f}_i \times (\mathbf{r}_{i-1,i} + \mathbf{r}_{i,C_i}) + \boldsymbol{\mu}_{i+1} + \mathbf{f}_{i+1} \times \mathbf{r}_{i,C_i} \\ &\quad + \bar{\mathbf{I}}_i \dot{\omega}_i + \omega_i \times (\bar{\mathbf{I}}_i \omega_i) \\ &\quad + k_{r,i+1} \ddot{q}_{i+1} I_{m_{i+1}} \mathbf{z}_{m_{i+1}} + k_{r,i+1} \dot{q}_{i+1} I_{m_{i+1}} \omega_i \times \mathbf{z}_{m_{i+1}} \end{aligned}$$

$$\tau_i = \begin{cases} \mathbf{f}_i^T \mathbf{z}_{i-1} + k_{ri} I_{mi} \dot{\omega}_{m_i}^T \mathbf{z}_{m_i} + F_{vi} \dot{d}_i + F_{si} \operatorname{sgn}(\dot{d}_i) & \text{prismatico} \\ \boldsymbol{\mu}_i^T \mathbf{z}_{i-1} + k_{ri} I_{mi} \dot{\omega}_{m_i}^T \mathbf{z}_{m_i} + F_{vi} \dot{\vartheta}_i + F_{si} \operatorname{sgn}(\dot{\vartheta}_i) & \text{rotoidale} \end{cases}$$

- Conviene riferire tutti i vettori alla terna corrente sul braccio  $i$

★ le quantità  $\bar{\mathbf{I}}_i^i$ ,  $\mathbf{r}_{i,C_i}^i$ ,  $\mathbf{z}_{m_i}^{i-1}$  sono *costanti*

$$\boldsymbol{\omega}_i^i = \begin{cases} \mathbf{R}_i^{i-1T} \boldsymbol{\omega}_{i-1}^{i-1} & \text{prismatico} \\ \mathbf{R}_i^{i-1T} (\boldsymbol{\omega}_{i-1}^{i-1} + \dot{\vartheta}_i \mathbf{z}_0) & \text{rotoidale} \end{cases} \quad (7.107)$$

$$\dot{\boldsymbol{\omega}}_i^i = \begin{cases} \mathbf{R}_i^{i-1T} \dot{\boldsymbol{\omega}}_{i-1}^{i-1} & \text{prismatico} \\ \mathbf{R}_i^{i-1T} (\dot{\boldsymbol{\omega}}_{i-1}^{i-1} + \ddot{\vartheta}_i \mathbf{z}_0 + \dot{\vartheta}_i \boldsymbol{\omega}_{i-1}^{i-1} \times \mathbf{z}_0) & \text{rotoidale} \end{cases} \quad (7.108)$$

$$\ddot{\mathbf{p}}_i^i = \begin{cases} \mathbf{R}_i^{i-1T} (\ddot{\mathbf{p}}_{i-1}^{i-1} + \ddot{d}_i \mathbf{z}_0) + 2\dot{d}_i \boldsymbol{\omega}_i^i \times \mathbf{R}_i^{i-1T} \mathbf{z}_0 \\ + \dot{\boldsymbol{\omega}}_i^i \times \mathbf{r}_{i-1,i}^i + \boldsymbol{\omega}_i^i \times (\boldsymbol{\omega}_i^i \times \mathbf{r}_{i-1,i}^i) & \text{prismatico} \\ \mathbf{R}_i^{i-1T} \ddot{\mathbf{p}}_{i-1}^{i-1} + \dot{\boldsymbol{\omega}}_i^i \times \mathbf{r}_{i-1,i}^i \\ + \boldsymbol{\omega}_i^i \times (\boldsymbol{\omega}_i^i \times \mathbf{r}_{i-1,i}^i) & \text{rotoidale} \end{cases} \quad (7.109)$$

$$\ddot{\mathbf{p}}_{C_i}^i = \ddot{\mathbf{p}}_i^i + \dot{\boldsymbol{\omega}}_i^i \times \mathbf{r}_{i,C_i}^i + \boldsymbol{\omega}_i^i \times (\boldsymbol{\omega}_i^i \times \mathbf{r}_{i,C_i}^i) \quad (7.110)$$

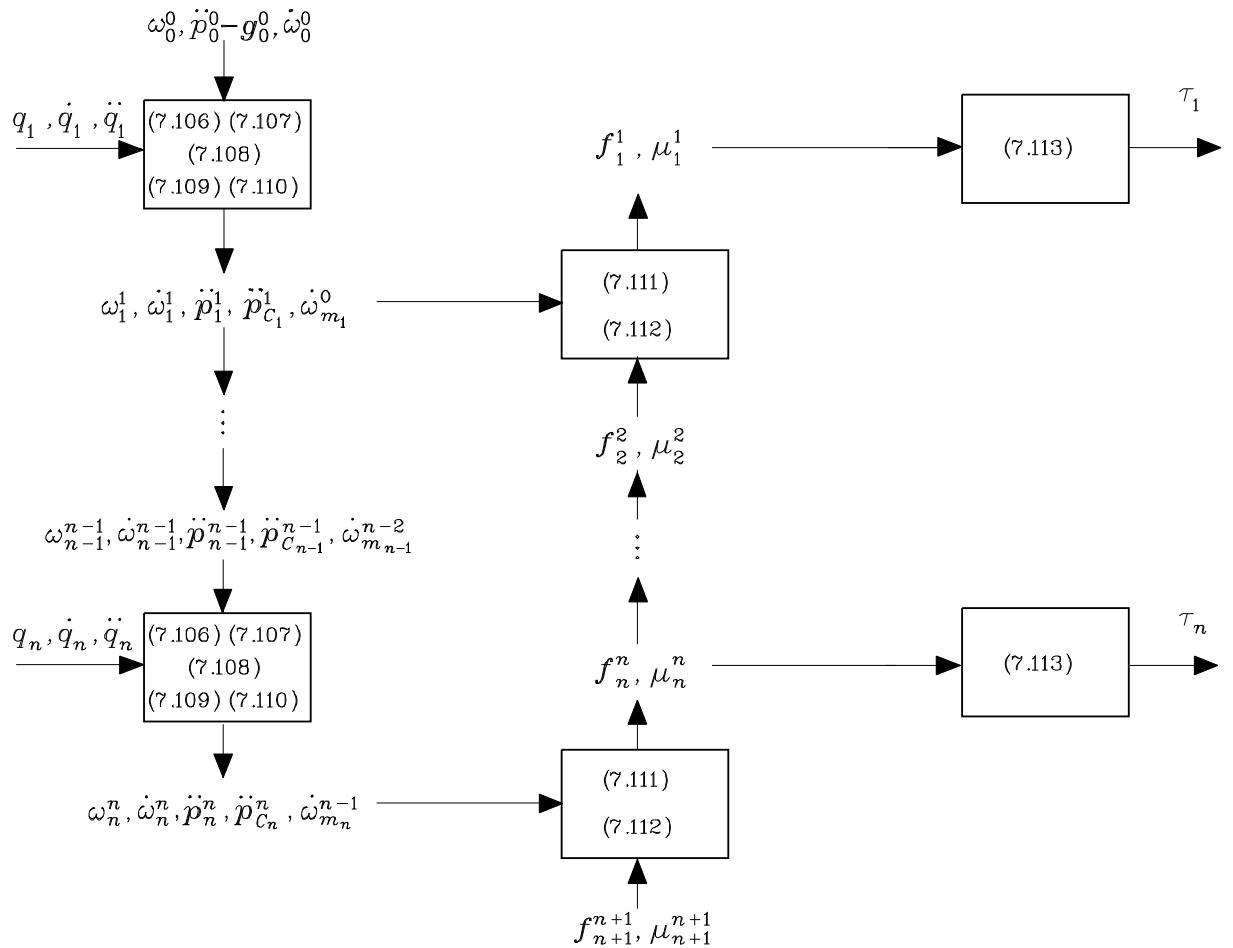
$$\dot{\boldsymbol{\omega}}_{m_i}^{i-1} = \boldsymbol{\omega}_{i-1}^{i-1} + k_{ri} \ddot{q}_i \mathbf{z}_{m_i}^{i-1} + k_{ri} \dot{q}_i \boldsymbol{\omega}_{i-1}^{i-1} \times \mathbf{z}_{m_i}^{i-1} \quad (7.111)$$

$$\mathbf{f}_i^i = \mathbf{R}_{i+1}^i \mathbf{f}_{i+1}^{i+1} + m_i \ddot{\mathbf{p}}_{C_i}^i \quad (7.112)$$

$$\begin{aligned} \boldsymbol{\mu}_i^i = & -\mathbf{f}_i^i \times (\mathbf{r}_{i-1,i}^i + \mathbf{r}_{i,C_i}^i) + \mathbf{R}_{i+1}^i \boldsymbol{\mu}_{i+1}^{i+1} \\ & + \mathbf{R}_{i+1}^i \mathbf{f}_{i+1}^{i+1} \times \mathbf{r}_{i,C_i}^i + \bar{\mathbf{I}}_i^i \dot{\boldsymbol{\omega}}_i^i + \boldsymbol{\omega}_i^i \times (\bar{\mathbf{I}}_i^i \boldsymbol{\omega}_i^i) \\ & + k_{r,i+1} \ddot{q}_{i+1} I_{m_{i+1}} \mathbf{z}_{m_{i+1}}^i + k_{r,i+1} \dot{q}_{i+1} I_{m_{i+1}} \boldsymbol{\omega}_i^i \times \mathbf{z}_{m_{i+1}}^i \end{aligned} \quad (7.113)$$

$$\tau_i = \begin{cases} \mathbf{f}_i^{iT} \mathbf{R}_i^{i-1T} \mathbf{z}_0 + k_{ri} I_{m_i} \dot{\boldsymbol{\omega}}_{m_i}^{i-1T} \mathbf{z}_{m_i}^{i-1} \\ + F_{vi} \dot{d}_i + F_{si} \operatorname{sgn}(\dot{d}_i) & \text{prismatico} \\ \boldsymbol{\mu}_i^{iT} \mathbf{R}_i^{i-1T} \mathbf{z}_0 + k_{ri} I_{m_i} \dot{\boldsymbol{\omega}}_{m_i}^{i-1T} \mathbf{z}_{m_i}^{i-1} \\ + F_{vi} \dot{\vartheta}_i + F_{si} \operatorname{sgn}(\dot{\vartheta}_i) & \text{rotoidale} \end{cases} \quad (7.114)$$

- Struttura computazionale



## Manipolatore planare a due bracci

- Condizioni iniziali

$$\ddot{\mathbf{p}}_0^0 - \mathbf{g}_0^0 = \begin{bmatrix} 0 \\ g \\ 0 \end{bmatrix} \quad \omega_0^0 = \dot{\omega}_0^0 = \mathbf{0}$$

- Condizioni terminali

$$\mathbf{f}_3^3 = \mathbf{0} \quad \mu_3^3 = \mathbf{0}$$

- Quantità riferite alla terna corrente

$$\mathbf{r}_{1,C_1}^1 = \begin{bmatrix} \ell_{C_1} \\ 0 \\ 0 \end{bmatrix} \quad \mathbf{r}_{0,1}^1 = \begin{bmatrix} a_1 \\ 0 \\ 0 \end{bmatrix}$$

$$\mathbf{r}_{2,C_2}^2 = \begin{bmatrix} \ell_{C_2} \\ 0 \\ 0 \end{bmatrix} \quad \mathbf{r}_{1,2}^2 = \begin{bmatrix} a_2 \\ 0 \\ 0 \end{bmatrix}$$

$$\mathbf{z}_{m_1}^1 = \mathbf{z}_0^1 = \mathbf{z}_{m_2}^1 = \mathbf{z}_{m_2}^2 = \mathbf{z}_1^2 = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

★ matrici di rotazione

$$\mathbf{R}_i^{i-1} = \begin{bmatrix} c_i & -s_i & 0 \\ s_i & c_i & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad i = 1, 2 \quad \mathbf{R}_3^2 = \mathbf{I}$$

- Ricorsione in avanti: braccio 1

$$\boldsymbol{\omega}_1^1 = \begin{bmatrix} 0 \\ 0 \\ \dot{\vartheta}_1 \end{bmatrix}$$

$$\dot{\boldsymbol{\omega}}_1^1 = \begin{bmatrix} 0 \\ 0 \\ \ddot{\vartheta}_1 \end{bmatrix}$$

$$\ddot{\boldsymbol{p}}_1^1 = \begin{bmatrix} -a_1 \dot{\vartheta}_1^2 + gs_1 \\ a_1 \ddot{\vartheta}_1 + gc_1 \\ 0 \end{bmatrix}$$

$$\ddot{\boldsymbol{p}}_{C_1}^1 = \begin{bmatrix} -(\ell_{C_1} + a_1) \dot{\vartheta}_1^2 + gs_1 \\ (\ell_{C_1} + a_1) \ddot{\vartheta}_1 + gc_1 \\ 0 \end{bmatrix}$$

$$\dot{\boldsymbol{\omega}}_{m_1}^0 = \begin{bmatrix} 0 \\ 0 \\ k_{r1} \ddot{\vartheta}_1 \end{bmatrix}$$

- Ricorsione in avanti: braccio 2

$$\boldsymbol{\omega}_2^2 = \begin{bmatrix} 0 \\ 0 \\ \dot{\vartheta}_1 + \dot{\vartheta}_2 \end{bmatrix}$$

$$\dot{\boldsymbol{\omega}}_2^2 = \begin{bmatrix} 0 \\ 0 \\ \ddot{\vartheta}_1 + \ddot{\vartheta}_2 \end{bmatrix}$$

$$\ddot{\boldsymbol{p}}_2^2 = \begin{bmatrix} a_1 s_2 \ddot{\vartheta}_1 - a_1 c_2 \dot{\vartheta}_1^2 - a_2 (\dot{\vartheta}_1 + \dot{\vartheta}_2)^2 + g s_{12} \\ a_1 c_2 \ddot{\vartheta}_1 + a_2 (\ddot{\vartheta}_1 + \ddot{\vartheta}_2) + a_1 s_2 \dot{\vartheta}_1^2 + g c_{12} \\ 0 \end{bmatrix}$$

$$\ddot{\boldsymbol{p}}_{C_2}^2 = \begin{bmatrix} a_1 s_2 \ddot{\vartheta}_1 - a_1 c_2 \dot{\vartheta}_1^2 - (\ell_{C_2} + a_2)(\dot{\vartheta}_1 + \dot{\vartheta}_2)^2 + g s_{12} \\ a_1 c_2 \ddot{\vartheta}_1 + (\ell_{C_2} + a_2)(\ddot{\vartheta}_1 + \ddot{\vartheta}_2) + a_1 s_2 \dot{\vartheta}_1^2 + g c_{12} \\ 0 \end{bmatrix}$$

$$\dot{\boldsymbol{\omega}}_{m_2}^1 = \begin{bmatrix} 0 \\ 0 \\ \ddot{\vartheta}_1 + k_{r2} \ddot{\vartheta}_2 \end{bmatrix}$$

- Ricorsione all'indietro: braccio 2

$$\mathbf{f}_2^2 = m_2 \ddot{\mathbf{p}}_{C_2}^2$$

$$\boldsymbol{\mu}_2^2 = \begin{bmatrix} * \\ * \\ \bar{I}_{2zz}(\ddot{\vartheta}_1 + \ddot{\vartheta}_2) + m_2(\ell_{C_2} + a_2)^2(\ddot{\vartheta}_1 + \ddot{\vartheta}_2) \\ + m_2 a_1 (\ell_{C_2} + a_2) c_2 \ddot{\vartheta}_1 \\ + m_2 a_1 (\ell_{C_2} + a_2) s_2 \dot{\vartheta}_1^2 + m_2 (\ell_{C_2} + a_2) g c_{12} \end{bmatrix}$$

$$\begin{aligned} \tau_2 = & \left( \bar{I}_{2zz} + m_2 \left( (\ell_{C_2} + a_2)^2 + a_1 (\ell_{C_2} + a_2) c_2 \right) + k_{r2} I_{m_2} \right) \ddot{\vartheta}_1 \\ & + \left( \bar{I}_{2zz} + m_2 (\ell_{C_2} + a_2)^2 + k_{r2}^2 I_{m_2} \right) \ddot{\vartheta}_2 \\ & + m_2 a_1 (\ell_{C_2} + a_2) s_2 \dot{\vartheta}_1^2 + m_2 (\ell_{C_2} + a_2) g c_{12} \end{aligned}$$

- Ricorsione all'indietro: braccio 1

$$\boldsymbol{f}_1^1 = \begin{bmatrix} -m_2(\ell_{C_2} + a_2)s_2(\ddot{\vartheta}_1 + \ddot{\vartheta}_2) - m_1(\ell_{C_1} + a_1)\dot{\vartheta}_1^2 - m_2a_1\dot{\vartheta}_1^2 \\ -m_2(\ell_{C_2} + a_2)c_2(\dot{\vartheta}_1 + \dot{\vartheta}_2)^2 + (m_1 + m_2)gs_1 \\ m_1(\ell_{C_1} + a_1)\ddot{\vartheta}_1 + m_2a_1\ddot{\vartheta}_1 + m_2(\ell_{C_2} + a_2)c_2(\ddot{\vartheta}_1 + \ddot{\vartheta}_2) \\ -m_2(\ell_{C_2} + a_2)s_2(\dot{\vartheta}_1 + \dot{\vartheta}_2)^2 + (m_1 + m_2)gc_1 \\ 0 \end{bmatrix}$$

$$\boldsymbol{\mu}_1^1 = \begin{bmatrix} * \\ * \\ \bar{I}_{1zz}\ddot{\vartheta}_1 + m_2a_1^2\ddot{\vartheta}_1 + m_1(\ell_{C_1} + a_1)^2\ddot{\vartheta}_1 + m_2a_1(\ell_{C_2} + a_2)c_2\ddot{\vartheta}_1 \\ +\bar{I}_{2zz}(\ddot{\vartheta}_1 + \ddot{\vartheta}_2) + m_2a_1(\ell_{C_2} + a_2)c_2(\ddot{\vartheta}_1 + \ddot{\vartheta}_2) \\ +m_2(\ell_{C_2} + a_2)^2(\ddot{\vartheta}_1 + \ddot{\vartheta}_2) + k_{r2}I_{m_2}\ddot{\vartheta}_2 \\ +m_2a_1(\ell_{C_2} + a_2)s_2\dot{\vartheta}_1^2 - m_2a_1(\ell_{C_2} + a_2)s_2(\dot{\vartheta}_1 + \dot{\vartheta}_2)^2 \\ +m_1(\ell_{C_1} + a_1)gc_1 + m_2a_1gc_1 + m_2(\ell_{C_2} + a_2)gc_{12} \end{bmatrix}$$

$$\begin{aligned} \tau_1 = & \left( \bar{I}_{1zz} + m_1(\ell_{C_1} + a_1)^2 + k_{r1}^2 I_{m_1} + \bar{I}_{2zz} \right. \\ & \left. + m_2 \left( a_1^2 + (\ell_{C_2} + a_2)^2 + 2a_1(\ell_{C_2} + a_2)c_2 \right) \right) \ddot{\vartheta}_1 \\ & + \left( \bar{I}_{2zz} + m_2 \left( (\ell_{C_2} + a_2)^2 + a_1(\ell_{C_2} + a_2)c_2 \right) + k_{r2}I_{m_2} \right) \ddot{\vartheta}_2 \\ & - 2m_2a_1(\ell_{C_2} + a_2)s_2\dot{\vartheta}_1\dot{\vartheta}_2 - m_2a_1(\ell_{C_2} + a_2)s_2\dot{\vartheta}_2^2 \\ & + (m_1(\ell_{C_1} + a_1) + m_2a_1)gc_1 + m_2(\ell_{C_2} + a_2)gc_{12} \end{aligned}$$

- Equivalenza col modello di Lagrange

- ★ parametri di bracci e rotorì

$$m_1 = m_{\ell_1} + m_{m_2}$$

$$m_1 \ell_{C_1} = m_{\ell_1} (\ell_1 - a_1)$$

$$\bar{I}_{1zz} + m_1 \ell_{C_1}^2 = \hat{I}_1 = I_{\ell_1} + m_{\ell_1} (\ell_1 - a_1)^2 + I_{m_2}$$

$$m_2 = m_{\ell_2}$$

$$m_2 \ell_{C_2} = m_{\ell_2} (\ell_2 - a_2)$$

$$\bar{I}_{2zz} + m_2 \ell_{C_2}^2 = \hat{I}_2 = I_{\ell_2} + m_{\ell_2} (\ell_2 - a_2)^2$$

# DINAMICA DIRETTA E DINAMICA INVERSA

- Formulazione di *Lagrange*
  - ★ è *sistematica* e di facile comprensione
  - ★ fornisce le equazioni del moto in una *forma analitica compatta* che evidenzia la matrice di inerzia, la matrice a fattore delle forze centrifughe e di Coriolis e il vettore delle forze gravitazionali (utile ai fini della *sintesi del controllo*)
  - ★ è efficace se si vogliono portare in conto effetti meccanici più complessi quali ad esempio le deformazioni elastiche dei bracci
- Formulazione di *Newton-Eulero*
  - ★ è intrinsecamente un metodo *ricorsivo* che risulta efficiente da un punto di vista computazionale

- Dinamica *diretta*
  - ★ note  $\mathbf{q}(t_0)$ ,  $\dot{\mathbf{q}}(t_0)$ ,  $\boldsymbol{\tau}(t)$  (e  $\mathbf{h}_e(t)$ ), determinare  $\ddot{\mathbf{q}}(t)$ ,  $\dot{\mathbf{q}}(t)$ ,  $\mathbf{q}(t)$  per  $t > t_0$
  - ★ utile in simulazione
- Dinamica *inversa*
  - ★ note  $\ddot{\mathbf{q}}(t)$ ,  $\dot{\mathbf{q}}(t)$ ,  $\mathbf{q}(t)$  (e  $\mathbf{h}_e(t)$ ), determinare  $\boldsymbol{\tau}(t)$
  - ★ utile per la pianificazione e il controllo
- Per un manipolatore ad  $n$  giunti, il numero di operazioni per il calcolo della dinamica è:
  - ★  $O(n^2)$  per la *dinamica diretta*
  - ★  $O(n)$  per la *dinamica inversa*

- Dinamica diretta con Lagrange

$$\ddot{\mathbf{q}} = \mathbf{B}^{-1}(\mathbf{q})(\boldsymbol{\tau} - \boldsymbol{\tau}')$$

$$\boldsymbol{\tau}'(\mathbf{q}, \dot{\mathbf{q}}) = \mathbf{C}(\mathbf{q}, \dot{\mathbf{q}})\dot{\mathbf{q}} + \mathbf{F}_v\dot{\mathbf{q}} + \mathbf{f}_s(\dot{\mathbf{q}}) + \mathbf{g}(\mathbf{q}) + \mathbf{J}^T(\mathbf{q})\mathbf{h}_e$$

- ★ note  $\mathbf{q}(t_k)$ ,  $\dot{\mathbf{q}}(t_k)$ ,  $\boldsymbol{\tau}(t_k)$ , si calcola  $\ddot{\mathbf{q}}(t_k)$
- ★ si integra numericamente con passo  $\Delta t$ :  $\dot{\mathbf{q}}(t_{k+1})$ ,  $\mathbf{q}(t_{k+1})$

- Dinamica diretta con Newton-Eulero

- ★ note  $\mathbf{q}$ ,  $\dot{\mathbf{q}}$ , si calcola:

$\boldsymbol{\tau}'$  come  $\boldsymbol{\tau}$  per  $\ddot{\mathbf{q}} = \mathbf{0}$

colonna  $\mathbf{b}_i$  di  $\mathbf{B}$  come  $\boldsymbol{\tau}$  per  $\mathbf{g}_0 = \mathbf{0}$ ,  $\dot{\mathbf{q}} = \mathbf{0}$ ,  $\ddot{q}_i = 1$  e  $\ddot{q}_j = 0$  per  $j \neq i$

- ★ nota  $\boldsymbol{\tau}$ , si calcola  $\ddot{\mathbf{q}}$  e si integra numericamente

# MODELLO DINAMICO NELLO SPAZIO OPERATIVO

- Ridondanza
  - ★ coordinate generalizzate (moti interni?)
- Singolarità
- Modello nello spazio dei giunti

$$\ddot{\mathbf{q}} = -\mathbf{B}^{-1}(\mathbf{q})\mathbf{C}(\mathbf{q}, \dot{\mathbf{q}})\dot{\mathbf{q}} - \mathbf{B}^{-1}(\mathbf{q})\mathbf{g}(\mathbf{q}) + \mathbf{B}^{-1}(\mathbf{q})\mathbf{J}^T(\mathbf{q})(\boldsymbol{\gamma} - \mathbf{h}_e)$$

- Cinematica differenziale del secondo ordine

$$\begin{aligned}\ddot{\mathbf{x}}_e &= \mathbf{J}_A(\mathbf{q})\ddot{\mathbf{q}} + \dot{\mathbf{J}}_A(\mathbf{q}, \dot{\mathbf{q}})\dot{\mathbf{q}} \\ &= -\mathbf{J}_A\mathbf{B}^{-1}\mathbf{C}\dot{\mathbf{q}} - \mathbf{J}_A\mathbf{B}^{-1}\mathbf{g} + \dot{\mathbf{J}}_A\dot{\mathbf{q}} + \mathbf{J}_A\mathbf{B}^{-1}\mathbf{J}_A^T(\boldsymbol{\gamma}_A - \mathbf{h}_A)\end{aligned}$$

- Trasformazioni

$$\begin{aligned}\boldsymbol{B}_A &= (\boldsymbol{J}_A \boldsymbol{B}^{-1} \boldsymbol{J}_A^T)^{-1} \\ \boldsymbol{C}_A \dot{\boldsymbol{x}}_e &= \boldsymbol{B}_A \boldsymbol{J}_A \boldsymbol{B}^{-1} \boldsymbol{C} \dot{\boldsymbol{q}} - \boldsymbol{B}_A \dot{\boldsymbol{J}}_A \dot{\boldsymbol{q}} \\ \boldsymbol{g}_A &= \boldsymbol{B}_A \boldsymbol{J}_A \boldsymbol{B}^{-1} \boldsymbol{g}\end{aligned}$$

★ manipolatore non ridondante

$$\begin{aligned}\boldsymbol{B}_A &= \boldsymbol{J}_A^{-T} \boldsymbol{B} \boldsymbol{J}_A^{-1} \\ \boldsymbol{C}_A \dot{\boldsymbol{x}}_e &= \boldsymbol{J}_A^{-T} \boldsymbol{C} \dot{\boldsymbol{q}} - \boldsymbol{B}_A \dot{\boldsymbol{J}}_A \dot{\boldsymbol{q}} \\ \boldsymbol{g}_A &= \boldsymbol{J}_A^{-T} \boldsymbol{g}\end{aligned}$$

- Equazioni del moto

$$\boldsymbol{B}_A(\boldsymbol{x}_e) \ddot{\boldsymbol{x}}_e + \boldsymbol{C}_A(\boldsymbol{x}_e, \dot{\boldsymbol{x}}_e) \dot{\boldsymbol{x}}_e + \boldsymbol{g}_A(\boldsymbol{x}_e) = \boldsymbol{\gamma}_A - \boldsymbol{h}_A$$

- Dinamica diretta
  - ★ note  $\tau(t)$  (e  $h_e(t)$ ), determinare  $\ddot{x}_e(t)$ ,  $\dot{x}_e(t)$ ,  $x_e(t)$
- Soluzione
  - ★ dinamica diretta ai giunti
  - ★ cinematica diretta
- Dinamica inversa
  - ★ note  $x_e(t)$ ,  $\dot{x}_e(t)$ ,  $\ddot{x}_e(t)$  (e  $h_e(t)$ ), determinare  $\tau(t)$
- Soluzione (ridondanza a livello cinematico)
  - ★ cinematica inversa
  - ★ dinamica inversa ai giunti
- Soluzione (ridondanza a livello dinamico)
  - ★ modello dinamico nello spazio operativo
  - ★ coppie con gestione di moti interni

- Dinamica inversa

$$\begin{aligned}
 \boldsymbol{\gamma}_A - \boldsymbol{h}_A &= \boldsymbol{B}_A(\boldsymbol{x}_e)\ddot{\boldsymbol{x}}_e + \boldsymbol{C}_A(\boldsymbol{x}_e, \dot{\boldsymbol{x}}_e)\dot{\boldsymbol{x}}_e + \boldsymbol{g}_A(\boldsymbol{x}_e) \\
 &= \boldsymbol{B}_A(\ddot{\boldsymbol{x}}_e - \dot{\boldsymbol{J}}_A \dot{\boldsymbol{q}}) + \boldsymbol{B}_A \boldsymbol{J}_A \boldsymbol{B}^{-1} \boldsymbol{C} \dot{\boldsymbol{q}} + \boldsymbol{B}_A \boldsymbol{J}_A \boldsymbol{B}^{-1} \boldsymbol{g} \\
 &= \boldsymbol{B}_A \boldsymbol{J}_A \ddot{\boldsymbol{q}} + \boldsymbol{B}_A \boldsymbol{J}_A \boldsymbol{B}^{-1} \boldsymbol{C} \dot{\boldsymbol{q}} + \boldsymbol{B}_A \boldsymbol{J}_A \boldsymbol{B}^{-1} \boldsymbol{g}
 \end{aligned}$$

★  $\boldsymbol{J}_A^T$ : pseudo-inversa destra pesata secondo  $\boldsymbol{B}^{-1}$  di

$$\bar{\boldsymbol{J}}_A^T(\boldsymbol{q}) = \boldsymbol{B}_A(\boldsymbol{q}) \boldsymbol{J}_A(\boldsymbol{q}) \boldsymbol{B}^{-1}(\boldsymbol{q})$$



$$\begin{aligned}
 \boldsymbol{\gamma}_A - \boldsymbol{h}_A &= \bar{\boldsymbol{J}}_A^T(\boldsymbol{B} \ddot{\boldsymbol{q}} + \boldsymbol{C} \dot{\boldsymbol{q}} + \boldsymbol{g}) \\
 &= \bar{\boldsymbol{J}}_A^T(\boldsymbol{\tau} - \boldsymbol{J}_A^T \boldsymbol{h}_A)
 \end{aligned}$$

$$\boldsymbol{\tau} = \boldsymbol{J}_A^T(\boldsymbol{q}) \boldsymbol{\gamma}_A + (\boldsymbol{I} - \boldsymbol{J}_A^T(\boldsymbol{q}) \bar{\boldsymbol{J}}_A^T(\boldsymbol{q})) \boldsymbol{\tau}_0$$

# ELLISSOIDE DI MANIPOLABILITÀ DINAMICA

- Sfera nello spazio dei giunti

$$\boldsymbol{\tau}^T \boldsymbol{\tau} = 1$$

★ modello dinamico con  $\dot{\boldsymbol{q}} = \mathbf{0}$ ,  $\boldsymbol{h} = \mathbf{0}$

$$\boldsymbol{B}(\boldsymbol{q})\ddot{\boldsymbol{q}} + \boldsymbol{g}(\boldsymbol{q}) = \boldsymbol{\tau}$$

★ accelerazioni ai giunti ( $\dot{\boldsymbol{q}} = \mathbf{0}$ )

$$\dot{\boldsymbol{v}}_e = \boldsymbol{J}(\boldsymbol{q})\ddot{\boldsymbol{q}}$$

$$\Downarrow$$

$$\boldsymbol{\tau} = \boldsymbol{B}(\boldsymbol{q})\boldsymbol{J}^\dagger(\boldsymbol{q})\dot{\boldsymbol{v}}_e + \boldsymbol{g}(\boldsymbol{q})$$

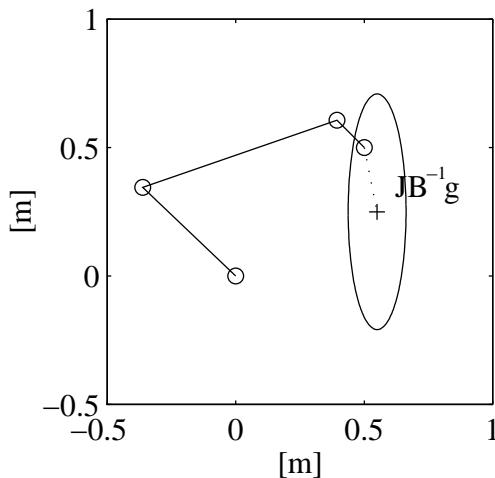
- Ellissoide nello spazio operativo

$$(\boldsymbol{B}(\boldsymbol{q})\boldsymbol{J}^\dagger(\boldsymbol{q})\dot{\boldsymbol{v}}_e + \boldsymbol{g}(\boldsymbol{q}))^T (\boldsymbol{B}(\boldsymbol{q})\boldsymbol{J}^\dagger(\boldsymbol{q})\dot{\boldsymbol{v}}_e + \boldsymbol{g}(\boldsymbol{q})) = 1$$

$$\begin{aligned}
\mathbf{B} \mathbf{J}^\dagger \dot{\mathbf{v}}_e + \mathbf{g} &= \mathbf{B}(\mathbf{J}^\dagger \dot{\mathbf{v}}_e + \mathbf{B}^{-1} \mathbf{g}) \\
&= \mathbf{B}(\mathbf{J}^\dagger \dot{\mathbf{v}}_e + \mathbf{B}^{-1} \mathbf{g} + \mathbf{J}^\dagger \mathbf{J} \mathbf{B}^{-1} \mathbf{g} - \mathbf{J}^\dagger \mathbf{J} \mathbf{B}^{-1} \mathbf{g}) \\
&= \mathbf{B}(\mathbf{J}^\dagger \dot{\mathbf{v}}_e + \mathbf{J}^\dagger \mathbf{J} \mathbf{B}^{-1} \mathbf{g} + (\mathbf{I} - \mathbf{J}^\dagger \mathbf{J}) \mathbf{B}^{-1} \mathbf{g}) \\
&= \mathbf{B} \mathbf{J}^\dagger (\dot{\mathbf{v}}_e + \mathbf{J} \mathbf{B}^{-1} \mathbf{g})
\end{aligned}$$

↓

$$(\dot{\mathbf{v}}_e + \mathbf{J} \mathbf{B}^{-1} \mathbf{g})^T \mathbf{J}^{\dagger T} \mathbf{B}^T \mathbf{B} \mathbf{J}^\dagger (\dot{\mathbf{v}}_e + \mathbf{J} \mathbf{B}^{-1} \mathbf{g}) = 1$$



★ manipolatore non ridondante

$$(\dot{\mathbf{v}}_e + \mathbf{J} \mathbf{B}^{-1} \mathbf{g})^T \mathbf{J}^{-T} \mathbf{B}^T \mathbf{B} \mathbf{J}^{-1} (\dot{\mathbf{v}}_e + \mathbf{J} \mathbf{B}^{-1} \mathbf{g}) = 1$$

## SCALATURA DINAMICA DI TRAIETTORIE

- ★ vincoli dinamici (coppie troppo elevate)

$$\begin{aligned}\boldsymbol{\tau}(t) &= \mathbf{B}(\mathbf{q}(t))\ddot{\mathbf{q}}(t) + \mathbf{C}(\mathbf{q}(t), \dot{\mathbf{q}}(t))\dot{\mathbf{q}}(t) + \mathbf{g}(\mathbf{q}(t)) \\ &= \mathbf{B}(\mathbf{q}(t))\ddot{\mathbf{q}}(t) + \boldsymbol{\Gamma}(\mathbf{q}(t))[\dot{\mathbf{q}}(t)\dot{\mathbf{q}}(t)] + \mathbf{g}(\mathbf{q}(t)) \\ &= \boldsymbol{\tau}_s(t) + \mathbf{g}(\mathbf{q}(t))\end{aligned}$$

$$\mathbf{C}(\mathbf{q}, \dot{\mathbf{q}})\dot{\mathbf{q}} = \boldsymbol{\Gamma}(\mathbf{q})[\dot{\mathbf{q}}\dot{\mathbf{q}}]$$

$$[\dot{\mathbf{q}}\dot{\mathbf{q}}] = [\dot{q}_1^2 \quad \dot{q}_1\dot{q}_2 \quad \dots \quad \dot{q}_{n-1}\dot{q}_n \quad \dot{q}_n^2]^T$$

★ scalatura temporale  $r(t)$ :  $r(0) = 0$   $r(t_f) = \bar{t}_f$

$$\mathbf{q}(t) = \bar{\mathbf{q}}(r(t))$$

$$\dot{\mathbf{q}} = \dot{r} \bar{\mathbf{q}}'(r)$$

$$\ddot{\mathbf{q}} = \dot{r}^2 \bar{\mathbf{q}}''(r) + \ddot{r} \bar{\mathbf{q}}'(r)$$

$\Downarrow$

$$\begin{aligned} \boldsymbol{\tau} &= \dot{r}^2 \left( \mathbf{B}(\bar{\mathbf{q}}(r)) \bar{\mathbf{q}}''(r) + \boldsymbol{\Gamma}(\bar{\mathbf{q}}(r)) [\bar{\mathbf{q}}'(r) \bar{\mathbf{q}}'(r)] \right) + \ddot{r} \mathbf{B}(\bar{\mathbf{q}}(r)) \bar{\mathbf{q}}'(r) \\ &\quad + \mathbf{g}(\bar{\mathbf{q}}(r)) \end{aligned}$$

$$= \boldsymbol{\tau}_s(t) + \mathbf{g}(\bar{\mathbf{q}}(r))$$

$$\bar{\boldsymbol{\tau}}_s(r) = \mathbf{B}(\bar{\mathbf{q}}(r)) \bar{\mathbf{q}}''(r) + \boldsymbol{\Gamma}(\bar{\mathbf{q}}(r)) [\bar{\mathbf{q}}'(r) \bar{\mathbf{q}}'(r)]$$

$\Downarrow$

$$\boldsymbol{\tau}_s(t) = \dot{r}^2 \bar{\boldsymbol{\tau}}_s(r) + \ddot{r} \mathbf{B}(\bar{\mathbf{q}}(r)) \bar{\mathbf{q}}'(r)$$

★ scelta semplice  $r(t) = ct$

$$\boldsymbol{\tau}_s(t) = c^2 \bar{\boldsymbol{\tau}}_s(ct)$$

giunto  $q_i$  in corrispondenza della violazione maggiore

$$\frac{|\boldsymbol{\tau}_s|}{|\bar{\boldsymbol{\tau}}_i - \mathbf{g}(q_i)|} = c^2$$